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Orthogonality for Quantum Latin Isometry Squares

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Goyeneche et al recently proposed a notion of orthogonality for quantum Latin squares, and showed that orthogonal quantum Latin squares yield quantum codes. We give a simplified characterization of orthogonality for quantum Latin squares, which we show is equivalent to the existing notion. We use this simplified characterization to give an upper bound for the number of mutually orthogonal quantum Latin squares of a given size, and to give the first examples of orthogonal quantum Latin squares that do not arise from ordinary Latin squares. We then discuss quantum Latin isometry squares, generalizations of quantum Latin squares recently introduced by Benoist and Nechita, and define a new orthogonality property for these objects, showing that it also allows the construction of quantum codes. We give a new characterization of unitary error bases using these structures.

1 Introduction

1.1 Summary

At QPL 2016 the present authors introduced *quantum Latin squares* [12, 15], as quantum structures generalizing the well-known Latin squares from classical combinatorics [6]. Since then this work has been built on separately by a number of researchers: in particular, by Goyeneche, Raissi, Di Martino and Życzkowski [7], who propose a notion of *orthogonality* for quantum Latin squares which allows the construction of quantum codes; and also by Benoist and Nechita [4], who introduce *matrices of partial isometries of type $(C1, C2, C3, C4)$* , generalizations of quantum Latin squares which characterize system-environment observables preserving a certain set of pointer states.

In this paper we give a new formulation of orthogonality for quantum Latin squares, and use it to relate and generalize the works just cited, and extend them in certain ways. In particular, we highlight the following key contributions.

- We give a new, simplified definition of orthogonality for quantum Latin squares, and show that it is equivalent to the existing definition of Goyeneche et al [7, Definition 3]. (Definition 4, Theorem 10.)
- We give the first example of a pair of orthogonal quantum Latin squares which are not equivalent to a pair of classical Latin squares. (Example 5 and Proposition 13.)
- We show that there can be at most $n - 1$ mutually orthogonal quantum Latin squares of dimension n (Theorem 18.)
- We introduce *quantum Latin isometry squares* based on the *matrices of partial isometries of type $(C1, C2, C3, C4)$* defined by Benoist and Nechita [4, Definition 3.2], and define a new notion of orthogonality for these objects. (Definitions 19 and 26.)

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- We show how orthogonal quantum Latin isometry squares can be used to build quantum codes. (Theorem 33.)
- We show that unitary error bases give rise to orthogonal pairs of quantum Latin isometry squares, and in fact can be characterized in terms of them. (Theorem 37.)

1.2 Related work

Since the introduction of quantum Latin squares by the present authors [15], two notions of orthogonality for quantum Latin squares have been introduced, both of which extend the standard notion for classical Latin squares.

The first such notion, to which we refer here as *left orthogonality*, was introduced by the first author [12], who showed it could be used to construct maximally entangled mutually unbiased bases. Given a pair of classical Latin squares which are left orthogonal by this definition, the left conjugates of each square are orthogonal Latin squares in the traditional sense [11]. This notion of orthogonality between QLS is not comparable to that which we study in this paper.

More recently, Goyeneche et al [7] introduced another notion of orthogonality for quantum Latin squares, which also extends the traditional definition for classical Latin squares, and the definition which we study here is equivalent. They extended their notion to quantum orthogonal arrays, more general objects which we do not consider here.

Remark 1. Note that the definition of mutually orthogonal quantum Latin squares introduced by Goyeneche et al includes not only families of quantum Latin squares satisfying an orthogonality condition but also entangled multi-partite states known as *essentially quantum Latin squares*. This more general definition is exactly equivalent to that of *perfect tensors* which are already well studied. We only consider orthogonality between two or more quantum Latin squares in this paper, and only use the terms orthogonal quantum Latin squares and mutually orthogonal quantum Latin squares to refer to such objects.

1.3 Outline

This paper has the following structure. In Section 2, we give background on quantum Latin squares, introduce our new definition of orthogonality, and explore its consequences, especially in relation to the work of Goyeneche et al [7]. In Section 3, we define quantum Latin isometry squares based on the work of Benoist and Nechita [4], and investigate a new notion of orthogonality for these objects as well as a connection to unitary error bases.

2 Quantum Latin squares and orthogonality

In this section we prove our main results concerning orthogonal quantum Latin squares. In Section 2.1 we recall the definition of quantum Latin squares, give our new definition of orthogonality, and give a nontrivial example. In Section 2.2 we show that our notion of orthogonality is equivalent to a previous, more complicated definition due to Goyeneche et al [7]. In Section 2.3 we explore the connection between equivalence and orthogonality of quantum Latin squares, and show that our example of orthogonal quantum Latin squares is not equivalent to a pair of orthogonal classical Latin squares. In Section 2.4, we give a simpler definition of orthogonality for families of quantum Latin squares, and show it agrees with that due to Goyeneche et al. In Section 2.5, we prove an upper bound on the number of mutually orthogonal quantum Latin squares that can exist in any dimension.

Then the following arrays are a pair of orthogonal quantum Latin squares of dimension 9:

0⟩	2⟩	1⟩	3⟩	5⟩	4⟩	6⟩	8⟩	7⟩	0⟩	2⟩	1⟩	3⟩	5⟩	4⟩	6⟩	8⟩	7⟩
2⟩	1⟩	0⟩	5⟩	4⟩	3⟩	8⟩	7⟩	6⟩	1⟩	0⟩	2⟩	4⟩	3⟩	5⟩	7⟩	6⟩	8⟩
1⟩	0⟩	2⟩	4⟩	3⟩	5⟩	7⟩	6⟩	8⟩	2⟩	1⟩	0⟩	5⟩	4⟩	3⟩	8⟩	7⟩	6⟩
6⟩	8⟩	7⟩	0⟩	2⟩	1⟩	3⟩	5⟩	4⟩	3⟩	5⟩	4⟩	6⟩	8⟩	7⟩	$U 0\rangle$	$U 2\rangle$	$U 1\rangle$
8⟩	7⟩	6⟩	2⟩	1⟩	0⟩	5⟩	4⟩	3⟩	4⟩	3⟩	5⟩	7⟩	6⟩	8⟩	$U 1\rangle$	$U 0\rangle$	$U 2\rangle$
7⟩	6⟩	8⟩	1⟩	0⟩	2⟩	4⟩	3⟩	5⟩	5⟩	4⟩	3⟩	8⟩	7⟩	6⟩	$U 2\rangle$	$U 1\rangle$	$U 0\rangle$
$U 3\rangle$	$U 5\rangle$	$U 4\rangle$	6⟩	8⟩	7⟩	$U 0\rangle$	$U 2\rangle$	$U 1\rangle$	$U 6\rangle$	$U 8\rangle$	$U 7\rangle$	0⟩	2⟩	1⟩	3⟩	5⟩	4⟩
$U 5\rangle$	$U 4\rangle$	$U 3\rangle$	8⟩	7⟩	6⟩	$U 2\rangle$	$U 1\rangle$	$U 0\rangle$	$U 7\rangle$	$U 6\rangle$	$U 8\rangle$	1⟩	0⟩	2⟩	4⟩	3⟩	5⟩
$U 4\rangle$	$U 3\rangle$	$U 5\rangle$	7⟩	6⟩	8⟩	$U 1\rangle$	$U 0\rangle$	$U 2\rangle$	$U 8\rangle$	$U 7\rangle$	$U 6\rangle$	2⟩	1⟩	0⟩	5⟩	4⟩	3⟩

(1)

We now consider some equivalent characterizations of orthogonality, which will be useful later.

Lemma 6. *Two quantum Latin squares Φ, Ψ are orthogonal if and only if one, and hence both, of the following equivalent conditions hold:*

$$\sum_{i,j=0}^{n-1} |\Phi_{ij}\rangle\langle\Phi_{ij}| \otimes |\Psi_{ij}\rangle\langle\Psi_{ij}| = \mathbb{I}_{n^2} \tag{2}$$

$$\sum_{i,j,p,q=0}^{n-1} \langle\Phi_{ij}|\Phi_{pq}\rangle\langle\Psi_{ij}|\Psi_{pq}\rangle|i\rangle\langle p| = \mathbb{I}_{n^2} \tag{3}$$

Proof. For the first condition, equation (2) says that if we sum up outer products of each element of the family $\{|\Phi_{ij}\rangle \otimes |\Psi_{ij}\rangle | i, j \in [n]\}$, we get the identity; clearly this is equivalent to the statement that the family yields an orthonormal basis. For the second condition, consider the linear map $S = \sum_{i,j} |ij\rangle\langle\Phi_{ij}| \langle\Psi_{ij}|$, an operator on $\mathbb{C}^n \otimes \mathbb{C}^n$. The quantum Latin squares Φ, Ψ are orthogonal if and only if this map is unitary, since it transports the orthonormal basis $\{|\Phi_{ij}\rangle \otimes |\Psi_{ij}\rangle | i, j \in [n]\}$ to the computational basis. Since it is an operator on a finite-dimensional Hilbert space, S is unitary if and only if it is an isometry, and equation (3) is the isometry condition. \square

Orthogonality of quantum Latin squares is unaffected by conjugation of one of the squares.

Definition 7. Given a quantum Latin square Ψ , its *conjugate* Ψ^* is the quantum Latin square with entries $(\Psi^*)_{ij} = (\Psi_{ij})^*$.

Lemma 8. *Two quantum Latin squares Φ, Ψ are orthogonal just when Φ^*, Ψ are orthogonal.*

Proof. Suppose Φ, Ψ are orthogonal quantum Latin squares. Then by equation (3), it follows that $\sum_{i,j,p,q=0}^{n-1} \langle\Phi_{ij}|\Phi_{pq}\rangle\langle\Psi_{ij}|\Psi_{pq}\rangle = \delta_{ip}\delta_{jq}$. So for all $(i, j) \neq (p, q)$ either $\langle\Phi_{ij}|\Phi_{pq}\rangle = 0$ or $\langle\Psi_{ij}|\Psi_{pq}\rangle = 0$, and we know that $\langle\Phi_{ij}|\Phi_{ij}\rangle = \langle\Psi_{ij}|\Psi_{ij}\rangle = 1$. Since $0, 1 \in \mathbb{R}$, we conclude that $\sum_{i,j,p,q=0}^{n-1} \langle\Phi_{ij}^*|\Phi_{pq}^*\rangle\langle\Psi_{ij}|\Psi_{pq}\rangle = \delta_{ip}\delta_{jq}$, and hence by equation (2) it follows that Φ^*, Ψ are orthogonal. The converse then follows since $(\Phi^*)^* = \Phi$. \square

2.2 Relationship to previous notion of orthogonality

The following definition of orthogonality for quantum Latin squares has recently been proposed. It is less conceptual than our Definition 4, and more complex to work with.

Definition 9 ([7], Definition 3). Two quantum Latin squares Φ, Ψ are *GRMZ-orthogonal* when for each tensor factor $X \in \{A, B, C\}$, the following holds:

$$\text{Tr}_X \left(\sum_{i,p,j=0}^{n-1} |\Phi_{ij}\rangle\langle\Phi_{ij}|_A \otimes |\Psi_{pj}\rangle\langle\Psi_{pj}|_B \otimes |i\rangle\langle p|_C \right) = \mathbb{I}_{n^2} \quad (4)$$

Note that in the presentation of this definition we have restricted the original definition to orthogonality between pairs of quantum Latin squares (please refer to Remark 1). We now show that Definition 9 is equivalent to our Definition 4.

Theorem 10. *Two quantum Latin squares Φ, Ψ are orthogonal if and only if they are GRMZ-orthogonal.*

Proof. We first consider equation (4) for the case $X = C$, which yields the following equation:

$$\sum_{i,j,p=0}^{n-1} |\Phi_{ij}\rangle\langle\Phi_{pj}| \otimes |\Psi_{ij}\rangle\langle\Psi_{pj}| \langle p|i\rangle = \sum_{i,j=0}^{n-1} |\Phi_{ij}\rangle\langle\Phi_{ij}| \otimes |\Psi_{ij}\rangle\langle\Psi_{ij}| \langle ij|ij\rangle = \mathbb{I}_{n^2} \quad (5)$$

This corresponds to our equation (2). By Lemma 6, it will hold if and only if Φ, Ψ are orthogonal.

We now show that the trace conditions over $X = A$ and $X = B$ in the GRMZ-orthogonality definition are redundant, in the sense that they hold automatically for all pairs of quantum Latin squares Φ, Ψ , regardless of orthogonality. We analyze the case that $X = B$; the case $X = A$ is similar. The trace condition yields the following equation:

$$\sum_{i,j,p=0}^{n-1} \langle\Phi_{ij}|\Phi_{pj}\rangle|\Psi_{ij}\rangle\langle\Psi_{pj}| \otimes |i\rangle\langle p| = \mathbb{I}_{n^2}$$

But this equation holds for any pair of quantum Latin squares Φ, Ψ , as follows:

$$\begin{aligned} \sum_{i,j,p=0}^{n-1} \langle\Phi_{ij}|\Phi_{pj}\rangle|\Psi_{ij}\rangle\langle\Psi_{pj}| \otimes |i\rangle\langle p| &= \sum_{i,j,p=0}^{n-1} \delta_{ip}|\Psi_{ij}\rangle\langle\Psi_{pj}| \otimes |i\rangle\langle p| = \sum_{i,j=0}^{n-1} |\Psi_{ij}\rangle\langle\Psi_{ij}| \otimes |i\rangle\langle i| \\ &= \sum_{i=0}^{n-1} \left(\sum_{j=0}^{n-1} |\Psi_{ij}\rangle\langle\Psi_{ij}| \right) \otimes |i\rangle\langle i| = \sum_{i=0}^{n-1} \mathbb{I}_n \otimes |i\rangle\langle i| = \mathbb{I}_n \otimes \left(\sum_{i=0}^{n-1} |i\rangle\langle i| \right) = \mathbb{I}_n \otimes \mathbb{I}_n \end{aligned}$$

Here the first equality uses the fact that Φ is a QLS, the second equality uses the definition of the Kronecker delta function, the third equality rearranges the sum, the fourth equality uses the fact that Ψ is a QLS, and the final equalities are trivial algebraic manipulations. This completes the proof. \square

2.3 Equivalence and orthogonality

Two classical Latin squares are said to be equivalent if one can be transformed into the other by permutations of the rows, columns or computational basis state labels. Similarly, there is a notion of equivalence between quantum Latin squares [15], which we now recall.

Definition 11. Two quantum Latin squares Φ, Ψ of dimension n are *equivalent* if there exists some unitary operator U on \mathbb{C}^n , family of modulus-1 complex numbers c_{ij} , and permutations $\sigma, \tau \in S_n$, such that the following holds for all $i, j \in [n]$:

$$c_{ij}U|\Phi_{\sigma(i),\tau(j)}\rangle = |\Psi_{ij}\rangle \quad (6)$$

Orthogonality is preserved by taking potentially different equivalences of each QLS, as long as the same pair of permutations are used.

Lemma 12. *Given quantum Latin squares Φ, Ψ, Φ', Ψ' of dimension n , unitary operators U, V on \mathbb{C}^n , families of modulus-1 complex numbers c_{ij}, d_{ij} , and permutations $\sigma, \tau \in S_n$ such that*

$$|\Phi'_{ij}\rangle := c_{ij}U|\Phi_{\sigma(i),\tau(j)}\rangle \quad |\Psi'_{ij}\rangle := d_{ij}V|\Psi_{\sigma(i),\tau(j)}\rangle \quad (7)$$

then Φ, Ψ are orthogonal if and only if Φ', Ψ' are orthogonal.

Proof. By Lemma 6 we have the following for all i, j, m, n :

$$\begin{aligned} \langle \Phi_{mn} | \Phi_{ij} \rangle \langle \Psi_{mn} | \Psi_{ij} \rangle &= \delta_{im} \delta_{jn} \\ \Leftrightarrow \langle \Phi_{m'n'} | U^\dagger \circ U | \Phi_{i'j'} \rangle \langle \Psi_{m'n'} | V^\dagger \circ V | \Psi_{i'j'} \rangle &= \delta_{i'm'} \delta_{j'n'} c_{i'j'}^* c_{i'j'}^* d_{i'j'}^* d_{i'j'}^* = \delta_{i'm'} \delta_{j'n'} c_{m'n'}^* c_{i'j'}^* d_{m'n'}^* d_{i'j'}^* \\ \Leftrightarrow \langle \Phi_{m'n'} | U^\dagger c_{m'n'}^* c_{i'j'}^* U | \Phi_{i'j'} \rangle \langle \Psi_{m'n'} | V^\dagger d_{m'n'}^* d_{i'j'}^* V | \Psi_{i'j'} \rangle &= \delta_{i'm'} \delta_{j'n'} \\ \Leftrightarrow \langle \Phi'_{mn} | \Phi'_{ij} \rangle \langle \Psi'_{mn} | \Psi'_{ij} \rangle &= \delta_{im} \delta_{jn} \end{aligned}$$

Where $\sigma(i) = i', \tau(j) = j', \sigma(m) = m'$ and $\tau(n) = n'$. \square

We now show that the pair of orthogonal quantum Latin squares illustrated in Example 5 are not equivalent to any pair of orthogonal Latin squares.

Proposition 13. *The orthogonal quantum Latin squares of Example 5 are not equivalent to a pair of orthogonal classical Latin squares.*

Proof. It is enough to show that the left-hand quantum Latin square of Example 5, which we call Φ , is not equivalent to a classical Latin square. Clearly no permutation of the rows or columns could transform Φ into a classical Latin square. Suppose for a contradiction that there exists a unitary operator V and a set of phases c_{ij} such that $|\eta_{ij}\rangle := c_{ij}V|\Phi_{ij}\rangle$ are all computational basis elements, and therefore yield a classical Latin square. Then for all i, j, m, n , we must have $\langle \eta_{mn} | \eta_{ij} \rangle = 0$ or 1 . We choose $m = 0, n = 3, i = 6$ and $j = 2$ to obtain $\langle \eta_{03} | \eta_{62} \rangle = c_{03}^* c_{62} \langle \Phi_{03} | V^\dagger V | \Phi_{62} \rangle = c_{03}^* c_{62} \langle \Phi_{03} | \Phi_{62} \rangle = c_{03}^* c_{62} \langle 3 | U | 4 \rangle = c_{03}^* c_{62} (1+i)/\sqrt{3}$. But since the c_{ij} have modulus 1, this can never equal 0 or 1, and the contradiction is established. \square

2.4 Generalization to multiple systems

We now extend this definition to families of quantum Latin squares, generalising mutually orthogonal Latin squares. In particular, we show that no essentially new concept is introduced, with the existing pairwise orthogonality property being sufficient.

Definition 14 (MOQLS). A family of m quantum Latin squares $\{\Phi^k | k \in [m]\}$ are *mutually orthogonal* if they are pairwise orthogonal.

We now present the definition of mutually orthogonal quantum Latin squares due to Goyeneche et al. As usual we only consider the definition with respect to families of quantum Latin squares (see Remark 1).

Definition 15 (GRMZ-MOQLS). A family of m quantum Latin squares $\{\Phi^k | k \in [m]\}$ are *GRMZ-mutually orthogonal* when the following equations hold, where X indicates a partial trace over any of the $m+1$ subsystems:

$$\text{Tr}_X \left(\sum_{i,j,p,q=0}^{n-1} |\Phi_{ij}^0\rangle \langle \Phi_{pq}^0| \otimes |\Phi_{ij}^1\rangle \langle \Phi_{pq}^1| \otimes \dots \otimes |\Phi_{ij}^{m-1}\rangle \langle \Phi_{pq}^{m-1}| \otimes |ij\rangle \langle pq| \right) = \mathbb{I}_{n^2} \quad (8)$$

This definition is equivalent to Definition 14.

Proposition 16 (MOQLS = GRMZ-MOQLS). *A family of m quantum Latin squares $\{\Phi^k | k \in [m]\}$ are mutually orthogonal just when they are GRMZ-mutually orthogonal.*

Proof. We label the k th QLS system by A_k and the two other systems in equation (8) as α and β . So X can range over m element subsets of $\{A_0, A_1, \dots, A_{m-1}, \alpha, \beta\}$. We will label such sets by the two elements that are NOT included so for example $(A_g, A_h) = \{A_0, \dots, A_{g-1}, A_{g+1}, \dots, A_{h-1}, A_{h+1}, \dots, A_{m-1}, \alpha, \beta\}$.

First we show that for all g and h , substituting $X = (A_g, A_h)$ into equation (8) reduces to equation (2) and so by varying g and h we obtain Definition 14. Let $X = (A_g, A_h)$ then we have:

$$\sum_{i,j,p,q=0}^{n-1} \langle \Phi_{pq}^0 | \Phi_{ij}^0 \rangle \dots \langle \Phi_{pq}^{g-1} | \Phi_{ij}^{g-1} \rangle | \Phi_{pq}^g \rangle \langle \Phi_{ij}^g | \langle \Phi_{pq}^{g+1} | \Phi_{ij}^{g+1} \rangle \dots \langle \Phi_{pq}^{h-1} | \Phi_{ij}^{h-1} \rangle | \Phi_{pq}^h \rangle \langle \Phi_{ij}^h | \langle \Phi_{pq}^{h+1} | \Phi_{ij}^{h+1} \rangle \dots \langle \Phi_{pq}^{m-1} | \Phi_{ij}^{m-1} \rangle \langle pq | ij \rangle = \mathbb{I}_{n^2} \quad \Leftrightarrow \quad \sum_{i,j=0}^{n-1} | \Phi_{ij}^g \rangle \langle \Phi_{ij}^g | \otimes | \Phi_{ij}^h \rangle \langle \Phi_{ij}^h | = \mathbb{I}_{n^2}$$

Thus by Lemma 6, equation (8) with $X = (A_g, A_h)$ holds if and only if Φ^g and Φ^h are orthogonal. We now show that Definition 14 implies equation (8) for all other possible values of X .

Since $\sum_{j=0}^{n-1} \langle \Phi_{ij} | \Phi_{pj} \rangle = \delta_{ip}$ by the quantum Latin square property, we have that for all k , substituting $X = (A_k, \alpha)$ into equation (8) reduces to $\sum_{i,j=0}^{n-1} | \Phi_{ij}^k \rangle \langle \Phi_{ij}^k | \otimes | j \rangle \langle j | = \mathbb{I}_{n^2}$, which holds for all QLS. Similarly by setting $X = (A_k, \beta)$ we obtain $\sum_{i,j=0}^{n-1} | \Phi_{ij}^k \rangle \langle \Phi_{ij}^k | \otimes | j \rangle \langle j | = \mathbb{I}_{n^2}$ which again holds for all QLS. Finally we are left with $X = (\alpha, \beta)$, which gives the following:

$$\sum_{i,j,p,q=0}^{n-1} \langle \Phi_{ij}^0 | \Phi_{pq}^0 \rangle \dots \langle \Phi_{ij}^{m-1} | \Phi_{pq}^{m-1} \rangle | ij \rangle \langle pq | = \mathbb{I}_{n^2} \tag{9}$$

Split the m QLS into pairs. Equation (3) is equivalent to $\sum_{i,j,p,q=0}^{n-1} \langle \Phi_{ij} | \Phi_{pq} \rangle \langle \Psi_{ij} | \Psi_{pq} \rangle = \delta_{ip} \delta_{qj}$. If m is even then the LHS of equation (9) becomes $\sum_{i,j,p,q=0}^{n-1} \delta_{ip} \delta_{pq} | ij \rangle \langle pq |$ which is a resolution of the identity. For m odd we have $\sum_{i,j=0}^{n-1} \langle \Phi_{ij}^{m-1} | \Phi_{ij}^{m-1} \rangle | ij \rangle \langle ij |$ which again is a resolution of the identity since all entries of a QLS are unit vectors. \square

It follows as a corollary of Lemma 12 that MOQLS are preserved by equivalences in the same way as pairs of orthogonal QLS.

Corollary 17. *Given a set of MOQLS Φ^k , the set of quantum Latin squares with entries $c_{ij} U_k | \Phi_{\sigma(i), \tau(j)}^k \rangle$ are also mutually orthogonal, for any set of unitary operators U_k , complex phases c_{ij} and permutations σ, τ .*

2.5 Upper bounds on the number of mutually orthogonal quantum Latin squares

We now show that the upper bound for the number of MOQLS of a given size is equal to the upper bound for MOLS.

Theorem 18. *Any family of MOQLS of dimension n has size at most $n - 1$.*

Proof. Suppose that we have a set of m -MOQLS $|\Phi_{ij}^0\rangle, \dots, |\Phi_{ij}^{m-1}\rangle$ of size $n \times n$. By Corollary 17 we can apply unitaries to each QLS such that the first row of every QLS is the ordered computational basis

$|i\rangle, i \in [n]$ so $|\Phi_{0i}^k\rangle = |i\rangle$ for all $k \in [m], i \in [n]$. Consider $\langle \Phi_{10}^k | \Phi_{10}^l \rangle$ for some $k, l \in [m]$ such that $k \neq l$. We have that:

$$\begin{aligned} \langle \Phi_{10}^k | \Phi_{10}^l \rangle &= \sum_{i=0}^{n-1} \langle \Phi_{10}^k | i \rangle \langle i | \Phi_{10}^l \rangle = \sum_{i=0}^{n-1} \langle \Phi_{10}^k | \Phi_{0i}^k \rangle \langle \Phi_{0i}^l | \Phi_{10}^l \rangle = \sum_{i=0}^{n-1} \langle \Phi_{10}^k | \Phi_{0i}^k \rangle \langle \Phi_{10}^{l*} | \Phi_{0i}^{l*} \rangle \\ &= \sum_{i,m,n,p,q=0}^{n-1} \langle \Phi_{10}^k | \Phi_{0i}^k \rangle \langle \Phi_{10}^{l*} | \Phi_{0i}^{l*} \rangle \langle 10 | mn \rangle \langle pq | 0i \rangle = \sum_i^{n-1} \langle 10 | 0i \rangle = \langle 1 | 0 \rangle = 0 \end{aligned}$$

The first equality is a resolution of the identity, the second holds since $|\Phi_{0i}^k\rangle = |\Phi_{0i}^l\rangle = |i\rangle$, the third is a straightforward property of inner products, the fourth equality is simple algebraic rearrangement and the fifth equality is due to Lemma 8 and equation 3. So the m unit vectors $|\Phi_{10}^i\rangle$ together with $|0\rangle$ are $m + 1$ linearly independent vectors. Thus m can be at most $n - 1$. \square

3 Quantum Latin isometry squares and quantum error detecting codes

In this section we introduce quantum Latin isometry squares a generalization of quantum Latin squares and use them to construct quantum error detecting codes. In Section 3.1 we give the definition of quantum isometry Latin squares and give a simple example. In Section 3.2 we compose pairs of compatible quantum isometry Latin squares and thereby recover both *matrices of partial isometries* and *projective permutation matrices*. In Section 3.3 we give orthogonality criteria for pairs and families of quantum isometry Latin squares. In Section 3.4 we show how quantum error detecting codes can be constructed from orthogonal pairs of quantum isometry Latin squares. Finally in Section 3.5 we show that *unitary error bases* can be characterized as quantum isometry Latin squares that are orthogonal to the *identity square*.

3.1 Quantum isometry Latin squares

A normalized vector $|\Psi\rangle$ of dimension n is a trivial example of an isometry $|\Psi\rangle : \mathbb{C} \rightarrow \mathbb{C}^n$. We can thus consider n dimensional QLS as arrays of isometries of this type. This perspective leads to the following definition which generalises QLS.

Definition 19 (Quantum isometry Latin square). An n -by- n array of isometries $k_{ij} : \mathbb{C}^{a_{ij}} \rightarrow \mathbb{C}^d$ is a *quantum isometry Latin square*, denoted (k_{ij}, a_{ij}, d) if the following hold for all $i, j, p, q \in \{0, \dots, n-1\}$:

$$k_{ip}^\dagger \circ k_{iq} = \delta_{pq} \mathbb{I}_{a_{ip}} \quad (10)$$

$$k_{pj}^\dagger \circ k_{qj} = \delta_{pq} \mathbb{I}_{a_{mj}} \quad (11)$$

$$\sum_{i=0}^{n-1} k_{ij} \circ k_{ij}^\dagger = \sum_{j=0}^{n-1} k_{ij} \circ k_{ij}^\dagger = \mathbb{I}_d \quad (12)$$

Remark 20. A quantum Latin square Φ of size $n \times n$ is a quantum isometry Latin squares such that $a_{ij} = 1$ for all i and j , and is therefore of the form $(|\Phi_{ij}\rangle, 1, n)$.

As a first example, we show how to construct quantum Latin isometry squares from arbitrary families of unitaries.

Example 21. For a Hilbert space \mathbb{C}^n equipped with a family of m unitaries $\mathcal{U} := \{U_i : \mathbb{C}^n \rightarrow \mathbb{C}^n | i \in [m]\}$, we can build a quantum Latin isometry square, denoted $L(\mathcal{U})$, of size m :

$$L(\mathcal{U}) := (U_i \delta_{ij}, n \delta_{ij}, n)$$

Such a quantum Latin isometry square $L(\mathcal{U})$ is diagonal, with nonzero isometries only on the leading diagonal. It is straightforward to see that equations (10),(11) and (12) are satisfied.

3.2 Skew projective permutation matrices

Pairs of quantum isometry Latin squares that share the same multiset of values a_{ij} can be composed to form a new structure.

Definition 22 (Skew projective permutation matrix). Given a pair of n -by- n quantum isometry Latin squares (k_{ij}, a_{ij}, b) and (q_{ij}, a_{ij}, b) , let $T_{ij} := q_{ij} \circ k_{ij}^\dagger$. We define the n -by- n array of linear operators $T_{ij} : \mathbb{C}^b \rightarrow \mathbb{C}^b$ to be a *skew projective permutation matrix* (skew PPM).

We now show that skew projective permutation matrices are precisely the *matrices of partial isometries of type (C1,C2,C3,C4)* introduced in a recent paper by Benoist and Nechita [4]. These structures were shown to characterise quantum channels preserving pointer states. We first require the following definition.

Definition 23 (Partial isometry [8]). A *partial isometry* is a linear map such that the restriction to the orthogonal complement of its kernel is an isometry. Alternatively, a partial isometry A is a linear map such that $A \circ A^\dagger \circ A = A$. The *initial space* of a partial isometry is the orthogonal complement of its kernel. The *final space* is its range.

Two examples of partial isometries are orthogonal projectors and unitaries. We now give Benoist and Nechita’s definition.

Definition 24 ([4], Definition 3.2 conditions (C1) to (C4)). An n -by- n *matrix of partial isometries of type (C1,C2,C3,C4)* and dimension b is an n -by- n array of partial isometries $T_{ij} : \mathbb{C}^b \rightarrow \mathbb{C}^b$ such that along each row and column the initial and final spaces of the T_{ij} partition \mathbb{C}^b .

Skew PPMs are matrices of partial isometries of type (C1,C2,C3,C4).

Lemma 25. Given a pair n -by- n quantum isometry Latin squares (k_{ij}, a_{ij}, b) and (q_{ij}, a_{ij}, b) , the corresponding skew PPM is a b dimensional n -by- n matrix of partial isometries of type (C1,C2,C3,C4).

Proof. Given the pair of isometries $k_{ij}, q_{ij} : \mathbb{C}^{a_{ij}} \rightarrow \mathbb{C}^b$ for some $i, j \in [n]$, we form the following composite linear maps:

$$K_{ij} := k_{ij} \circ k_{ij}^\dagger \qquad Q_{ij} := q_{ij} \circ q_{ij}^\dagger \qquad T_{ij} := q_{ij} \circ k_{ij}^\dagger$$

It is easy to see that K_{ij} and Q_{ij} are orthogonal projectors. The linear map T_{ij} is a partial isometry since $T_{ij} \circ T_{ij}^\dagger \circ T_{ij} = q_{ij} \circ k_{ij}^\dagger \circ k_{ij} \circ q_{ij}^\dagger \circ q_{ij} \circ k_{ij}^\dagger = q_{ij} \circ k_{ij}^\dagger = T_{ij}$ with each equality holding either by definition or using the properties of an isometry. It can easily be checked that the initial and final spaces of each T_{ij} are the spaces projected onto by K_{ij} and Q_{ij} respectively. By equations (10),(11) and (12) the initial and final spaces of the T_{ij} partition \mathbb{C}^b along the rows and columns as required. \square

Skew PPMs also appear in another, currently very active area of research. Skew PPMs are a generalisation of *projective permutation matrices* (PPMs). PPMs are square arrays of orthogonal projectors that form a projective POVM on every row and column. PPMs are skew PPMs coming from pairs of identical quantum isometry Latin squares. In this case the partial isometries T_{ij} are all orthogonal projectors, having the same initial and final spaces. Clearly these projectors form POVMs on every row and column since the spaces they project onto partition the whole Hilbert space along every row and column by Lemma 25.

PPMs, also known as magic unitaries and quantum bijections between classical sets have recently appeared in the context of quantum non-local games [1, 2, 13, 14] and the study of compact quantum groups [3, 5, 17].

3.3 Orthogonal quantum isometry Latin squares

We now extend the definition of orthogonal quantum Latin squares to quantum isometry Latin squares.

Definition 26 (Orthogonal quantum isometry Latin squares). A pair of quantum isometry Latin squares (k_{ij}, a_{ij}, d) and (q_{ij}, a_{ij}, d) are *orthogonal* if the operators $T_{ij} = q_{ij} \circ k_{ij}^\dagger$ span the space of operators and for all non-zero T_{ij} we have that $\text{Tr}(T_{ij}^\dagger \circ T_{ij}) = a$ for some $a \in \mathbb{C}$. We say that the T_{ij} form an orthogonal skew PPM.

We now give a more algebraic characterisation of orthogonal quantum isometry Latin squares.

Lemma 27. Given a pair of $n \times n$ quantum isometry Latin squares, $K = (k_{ij}, a_{ij}, d)$ and $Q = (q_{ij}, a_{ij}, d)$ define the following linear map $S : \mathbb{C}^n \otimes \mathbb{C}^n \rightarrow \mathbb{C}^d \otimes \mathbb{C}^d$:

$$S := \sum_{i,j,p,q=0}^{n-1} \sum_{x=0}^{d-1} |ij\rangle \langle x| q_{ij} \circ k_{ij}^\dagger \otimes \langle x| \quad (13)$$

K and Q are orthogonal just when S is an isometry.

Proof. Assuming K and Q are orthogonal, the partial isometries $q_{ij} \circ k_{ij}^\dagger$ span the operator space. We therefore have:

$$\mathbb{I}_{d^2} = \sum_{i,j=0}^{n-1} \sum_{x,y=0}^{d-1} |y\rangle \otimes k_{ij} \circ q_{ij}^\dagger |y\rangle \langle x| q_{ij} \circ k_{ij}^\dagger \otimes \langle x| \quad (14)$$

$$= \sum_{i,j,p,q=0}^{n-1} \sum_{x,y=0}^{d-1} |y\rangle \otimes k_{pq} \circ q_{pq}^\dagger |y\rangle \langle pq|ij\rangle \langle x| q_{ij} \circ k_{ij}^\dagger \otimes \langle x| \quad (15)$$

$$= S^\dagger \circ S \quad (16)$$

The other direction follows straightforwardly. \square

Remark 28. Note that $\text{Tr}(T_{ij}^\dagger \circ T_{pq}) = \delta_{ip} \delta_{jq} \text{Tr}(T_{ij}^\dagger \circ T_{ij})$.

Remark 29. A PPM can never be orthogonal since the projectors T_{ij} of a PPM span the operator space for every row and column.

Orthogonal quantum isometry Latin squares generalise orthogonal QLS (and therefore orthogonal Latin squares).

Lemma 30. Pairs of QLS are orthogonal quantum isometry Latin squares if and only if they are orthogonal quantum Latin squares.

Proof. Consider a pair of n -by- n QLS $(|k_{ij}\rangle, 1, n)$ and $(|q_{ij}\rangle, 1, n)$ such that they are orthogonal quantum isometry Latin squares by Definition 26. By Lemma 27, S is an isometry. Since S is a linear operator on a finite dimensional Hilbert space, S is unitary. This yields the following equation:

$$\sum_{i,j,x,y=0}^{n-1} |k_{ij}\rangle\langle q_{ij}|x\rangle\langle y|q_{ij}\rangle\langle k_{ij}| \otimes |x\rangle\langle y| = \sum_{i,j=0}^{n-1} |k_{ij}\rangle\langle k_{ij}| \otimes |q_{ij}^*\rangle\langle q_{ij}^*| = \mathbb{I}_{n^2}$$

By Lemmas 6 and 8 this holds if and only if $|q_{ij}\rangle$ and $|k_{ij}\rangle$ are orthogonal QLS. □

We define *mutually orthogonal quantum isometry Latin squares*, to be sets of pairwise orthogonal quantum isometry Latin squares thus generalising MOLS and MOQLS (see Definition 14).

We now present an example pair of orthogonal quantum isometry Latin squares which are not QLS.

Example 31. We present a pair of orthogonal quantum Latin isometry squares Q and K and associated orthogonal skew PPM, T . We have $n = 8$, $d = 4$ and $a_{ij} = 2$ or 0 for all $i, j \in \{0, \dots, 7\}$. There are $d^2 = 16$ non-zero T_{ij} as required to span the operator space.

We fix the computational basis $|a\rangle, |b\rangle$ for \mathbb{C}^2 and $|0\rangle, |1\rangle, |2\rangle, |3\rangle$ for \mathbb{C}^4 .

We present the first quantum Latin isometry square Q :

$ 0\rangle\langle a + 1\rangle\langle b $	$ 2\rangle\langle a + 3\rangle\langle b $	0	0	0	0	0	0
$ 2\rangle\langle a + 3\rangle\langle b $	$ 1\rangle\langle a + 0\rangle\langle b $	0	0	0	0	0	0
0	0	$ 0\rangle\langle a - 1\rangle\langle b $	$ 3\rangle\langle b - 2\rangle\langle a $	0	0	0	0
0	0	$ 2\rangle\langle a - 3\rangle\langle b $	$ 0\rangle\langle a - 1\rangle\langle b $	0	0	0	0
0	0	0	0	$ 1\rangle\langle a + 2\rangle\langle b $	$ 0\rangle\langle a + 3\rangle\langle b $	0	0
0	0	0	0	$ 3\rangle\langle b - 0\rangle\langle a $	$ 2\rangle\langle b - 1\rangle\langle a $	0	0
0	0	0	0	0	0	$ 0\rangle\langle a + 2\rangle\langle b $	$ 3\rangle\langle b - 1\rangle\langle a $
0	0	0	0	0	0	$ 1\rangle\langle a + 3\rangle\langle b $	$ 2\rangle\langle b - 0\rangle\langle a $

Now we present the quantum Latin isometry square K :

$ 0\rangle\langle a + 1\rangle\langle b $	$ 2\rangle\langle a + 3\rangle\langle b $	0	0	0	0	0	0
$ 2\rangle\langle a + 3\rangle\langle b $	$ 0\rangle\langle a + 1\rangle\langle b $	0	0	0	0	0	0
0	0	$ 0\rangle\langle a + 1\rangle\langle b $	$ 3\rangle\langle a + 2\rangle\langle b $	0	0	0	0
0	0	$ 2\rangle\langle a + 3\rangle\langle b $	$ 1\rangle\langle a + 0\rangle\langle b $	0	0	0	0
0	0	0	0	$ 2\rangle\langle a + 1\rangle\langle b $	$ 3\rangle\langle a + 0\rangle\langle b $	0	0
0	0	0	0	$ 3\rangle\langle a + 0\rangle\langle b $	$ 2\rangle\langle a + 1\rangle\langle b $	0	0
0	0	0	0	0	0	$ 2\rangle\langle a + 0\rangle\langle b $	$ 3\rangle\langle a + 1\rangle\langle b $
0	0	0	0	0	0	$ 3\rangle\langle a + 1\rangle\langle b $	$ 2\rangle\langle a + 0\rangle\langle b $

Finally we present the associated skew projective permutation matrix T with entries T_{ij} :

$ 0\rangle\langle 0 + 1\rangle\langle 1 $	$ 3\rangle\langle 2 + 2\rangle\langle 3 $	0	0	0	0	0	0
$ 2\rangle\langle 2 + 3\rangle\langle 3 $	$ 1\rangle\langle 0 + 0\rangle\langle 1 $	0	0	0	0	0	0
0	0	$ 0\rangle\langle 0 - 1\rangle\langle 1 $	$ 3\rangle\langle 2 - 2\rangle\langle 3 $	0	0	0	0
0	0	$ 2\rangle\langle 2 - 3\rangle\langle 3 $	$ 0\rangle\langle 1 - 1\rangle\langle 0 $	0	0	0	0
0	0	0	0	$ 2\rangle\langle 1 + 1\rangle\langle 2 $	$ 3\rangle\langle 0 + 0\rangle\langle 3 $	0	0
0	0	0	0	$ 3\rangle\langle 0 - 0\rangle\langle 3 $	$ 2\rangle\langle 1 - 1\rangle\langle 2 $	0	0
0	0	0	0	0	0	$ 2\rangle\langle 0 + 0\rangle\langle 2 $	$ 3\rangle\langle 1 - 1\rangle\langle 3 $
0	0	0	0	0	0	$ 3\rangle\langle 1 + 1\rangle\langle 3 $	$ 2\rangle\langle 0 - 0\rangle\langle 2 $

3.4 Quantum error detecting codes

We now prove the main result of this section, a construction of quantum error detecting codes from orthogonal skew PPMs. In their famous 1997 paper Knill and Laflamme proved that certain one-to-three 4-valent tensors can be used as encoding maps for quantum codes that can detect any single local error.

Theorem 32. [10] *Given a three-to-one tensor $\langle E_{ijk} : \mathbb{C}^a \rightarrow \mathbb{C}^b \otimes \mathbb{C}^c \otimes \mathbb{C}^d$, it is an encoding map that detects a single error if the following hold:*

$$\sum_{i=0}^{b-1} \sum_{j=0}^{c-1} \sum_{k=0}^{d-1} |E_{ijl}\rangle \langle E_{ijk}| \otimes |l\rangle \langle k| = \mathbb{I}_a \otimes \mathbb{I}_d \quad (17)$$

$$\sum_{i=0}^{b-1} \sum_{l=0}^{c-1} \sum_{k=0}^{d-1} |E_{ilk}\rangle \langle E_{ijk}| \otimes |l\rangle \langle j| = \mathbb{I}_a \otimes \mathbb{I}_c \quad (18)$$

$$\sum_{l=0}^{b-1} \sum_{j=0}^{c-1} \sum_{k=0}^{d-1} |E_{ljk}\rangle \langle E_{ijk}| \otimes |l\rangle \langle i| = \mathbb{I}_a \otimes \mathbb{I}_b \quad (19)$$

Theorem 33. *Given an n -by- n pair of orthogonal quantum isometry Latin squares (k_{ij}, a_{ij}, d) and (q_{ij}, a_{ij}, d) ; the following one-to-three tensor is an encoding map that detects a single error:*

$$\langle T| := \sum_{i,j=0}^{n-1} |i\rangle \otimes q_{ij} \circ k_{ij}^\dagger \otimes |j\rangle \quad (20)$$

Proof. First we show equation (17), we have:

$$\begin{aligned} \sum_{x=0}^{d-1} \sum_{l,i,j=0}^{n-1} |T\rangle \langle T| \otimes |l\rangle \langle i| &= \sum_{x=0}^{d-1} \sum_{l,i,j=0}^{n-1} |l\rangle \otimes k_{lj} \circ q_{lj}^\dagger |x\rangle \langle x| q_{ij} \circ k_{ij}^\dagger \otimes \langle i| \\ &\stackrel{(11)}{=} \sum_{i,j=0}^{n-1} |i\rangle \otimes k_{ij} \circ q_{ij}^\dagger \otimes \langle i| \stackrel{(12)}{=} \mathbb{I}_d \otimes \mathbb{I}_n \end{aligned}$$

Equation (19) can be derived from equations (10) and (12) similarly.

We now show equation (18).

$$\begin{aligned} \sum_{x,y=0}^{d-1} \sum_{i,j=0}^{n-1} |T\rangle \langle T| \otimes |y\rangle \langle x| &= \sum_{x,y=0}^{d-1} \sum_{i,j=0}^{n-1} |y\rangle \otimes k_{ij} \circ q_{ij}^\dagger |y\rangle \langle x| q_{ij} \circ k_{ij}^\dagger \otimes \langle x| \\ &\stackrel{(15)}{=} \mathbb{I}_{d^2} \end{aligned}$$

□

Example 34. Given the orthogonal skew PPM $T = \{T_{ij} | i, j \in \{0, \dots, 7\}\}$ as in Example 31, by Theorem 33 the following three-to-one 4-valent tensor is an encoding map that detects a single qubit local error:

$$\langle T | := \sum_{i,j=0}^{n-1} |i\rangle \otimes T_{ij} \otimes |j\rangle$$

3.5 Orthogonal quantum Latin isometry squares from unitary error bases

Here we recall the standard notion of unitary error basis, and show that they can be characterized as orthogonal pairs of quantum Latin isometry squares, which are not quantum Latin squares. Unitary error bases were introduced by Werner [18], and provide the basic data for quantum teleportation, dense coding and error correction procedures [9, 16, 18].

Definition 35. For a Hilbert space \mathbb{C}^n , a *unitary error basis* is a family of unitary operators $U_i : \mathbb{C}^n \rightarrow \mathbb{C}^n$ which span the space of operators, and which are orthogonal under the trace inner product:

$$\text{Tr}(U_i \circ U_j^\dagger) = n\delta_{ij} \tag{21}$$

We have already seen in Example 21 that any family of unitaries is a quantum Latin isometry square. We now show that unitary error bases can be characterized in terms of an orthogonal pair of quantum Latin isometry squares. We first define the *identity square*.

Definition 36 (Identity square). The d -dimensional *identity square* is the quantum Latin isometry square given by $(\mathbb{I}_d \delta_{ij}, d\delta_{ij}, d)$. Note that the identity square is also a skew PPM and a PPM.

We now present the main result of this section.

Theorem 37. *The following are equivalent:*

- *the set of all n^2 -by- n^2 quantum Latin isometry squares that are orthogonal to the n^2 -dimensional identity square;*
- *the set of all unitary error basis for \mathbb{C}^n .*

Proof. First note that any quantum Latin isometry square that is orthogonal to the identity square must be of the form $(X_{ij}, d\delta_{ij}, d)$ in order to be composed. The isometries X_{ij} are linear operators and must therefore be unitary. The orthogonality condition for the quantum Latin isometry squares unpacks to the requirement that the unitary operators X_i are orthogonal and span the space of operators; this is exactly the unitary error basis condition. □

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