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ATLAS Collaboration; Aad, G.

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Search for Higgs Boson Decays into a Z Boson and a Light Hadronically Decaying Resonance Using 13 TeV pp Collision Data from the ATLAS Detector

G. Aad *et al.**
(ATLAS Collaboration)

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A search for Higgs boson decays into a Z boson and a light resonance in two-lepton plus jet events is performed, using a pp collision dataset with an integrated luminosity of 139 fb^{-1} collected at $\sqrt{s} = 13 \text{ TeV}$ by the ATLAS experiment at the CERN LHC. The resonance considered is a light boson with a mass below 4 GeV from a possible extended scalar sector or a charmonium state. Multivariate discriminants are used for the event selection and for evaluating the mass of the light resonance. No excess of events above the expected background is found. Observed (expected) 95% confidence-level upper limits are set on the Higgs boson production cross section times branching fraction to a Z boson and the signal resonance, with values in the range 17–340 pb (16_{-5}^{+6} – 320_{-90}^{+130} pb) for the different light spin-0 boson mass and branching fraction hypotheses, and with values of 110 and 100 pb (100_{-30}^{+40} and 100_{-30}^{+40} pb) for the η_c and J/ψ hypotheses, respectively.

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The structure of the standard model (SM) scalar sector is the subject of intense scrutiny by the ATLAS [1] and CMS [2] Collaborations at the CERN Large Hadron Collider (LHC) [3]. At the current level of precision, all of the measured properties of the Higgs boson (H) [4,5] are found to be consistent with their SM predictions [6–10], and no additional Higgs boson has been observed to date. However, given the small natural decay width of the Higgs boson, even small additional contributions from physics beyond the SM can lead to final states with substantial, and thus possibly detectable, branching fractions (\mathcal{B}) [11]. This Letter presents a search for Higgs boson decays into a Z boson and a hadronically decaying light resonance in events with a same-flavor lepton pair (electrons or muons) and a jet in the ATLAS detector. Hadronic decays of an η_c or of a J/ψ charmonium resonance (Q), or of a light spin-0 boson from an extended Higgs sector with a mass up to 4 GeV, are considered and are reconstructed as a single jet.

The Yukawa sector of the SM [12] does not provide an explanation for the observed fermion mass hierarchy. As a result, a wide range of new physics scenarios have been proposed, including the Froggatt-Nielsen mechanism [13] and the Higgs-dependent Yukawa couplings model [14]; for a recent overview, see Ref. [15]. The couplings of the

Higgs boson to the third-generation fermions [16–21] have been observed, and a program to probe its couplings to the first- and second-generation charged leptons has been established [22–25]. For its couplings to first- and second-generation quarks, several approaches are being explored. Focusing on the Higgs boson’s coupling to the charm quark, direct searches have been performed for Higgs boson decays into charm quarks [26,27] and for exclusive decays into a J/ψ and a photon [28,29], with no excess observed. Constraints from differential cross section measurements of Higgs boson production versus transverse momentum (p_T) have also been derived [30,31]. Higgs boson decays into a gauge boson and a charmonium state, including an η_c or a J/ψ , have been proposed as another way to access the coupling of the Higgs boson to the charm quark [32–34] and to probe the nature of the Higgs boson [35]. This search follows the last approach and maximizes the signal acceptance by focusing on inclusive hadronic final states of the mesons in $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$ decays, which have SM branching fractions of 1.4×10^{-5} and 2.2×10^{-6} [35], respectively.

While the SM posits a single complex Higgs doublet field [36,37], extended Higgs sectors are motivated [38] and provide a rich phenomenology of additional scalars. Two such models discussed here are the two-Higgs-doublet model (2HDM) [11,39] and the 2HDM with an additional scalar singlet (2HDM + S) [11,40]. These represent two of the simplest extensions of the scalar sector, and with their type-II fermion couplings they are necessary to generate the masses in the minimal supersymmetric SM and the next-to-minimal supersymmetric SM, respectively [41]. Both of these models can include additional light pseudoscalars (a) with significant $\mathcal{B}(H \rightarrow Za)$ or $\mathcal{B}(H \rightarrow aa)$ [11]. In the

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2HDM(+S), these two \mathcal{B} values can be adjusted independently, therefore searches for $H \rightarrow aa$ do not constrain $\mathcal{B}(H \rightarrow Za)$, so that searches for the latter decay are required [11,34]. Despite the Yukawa nature of the a to fermion couplings, there are large regions of parameter space depending on the mass of a and the ratio of the vacuum expectation values of the two Higgs-doublet fields ($\tan\beta$) [11], where these pseudoscalars decay mainly to gluons and light up-type quarks, as the decays into down-type fermions are suppressed. These experimental signatures are also relevant in axion models [42–44], models of electroweak baryogenesis [45], neutrino mass models [46], dark-matter models [46,47], and models of grand unification [48]. Previous searches for Higgs boson decays into light scalars have been performed at the Tevatron [49] and the LHC [50–59]. However, these were mostly focused on searches for $H \rightarrow aa$, in final states including leptons, photons, or bottom quarks. By targeting the $H \rightarrow Za$, $a \rightarrow$ hadrons decay channel, this search accesses new, previously unexplored regions of the parameter space.

Searches for hadronic decays of light resonances are challenging at the LHC due to the large multijet background. However, substantial progress has been made in the use of jet substructure techniques in boosted final states [60], typically in searches or measurements involving heavy resonances [61,62]. In this Letter, jet substructure variables enable the reconstruction of a light, boosted, hadronic final state. Information from the individual substructure variables is combined using machine learning techniques. Specifically, for event selection, a multilayer perceptron (MLP) [63] classifier is employed. Given the range of masses considered, the classifier is provided with resonance-mass-related information from a separate MLP-based mass estimator, which results in improved classification performance over the full mass range.

This search is performed using the complete run 2 pp collision dataset, produced between 2015 and 2018 at a center-of-mass energy $\sqrt{s} = 13$ TeV by the LHC. The data were collected by the ATLAS detector [1] and correspond to an integrated luminosity of 139 fb^{-1} .

Monte Carlo (MC) samples of simulated events are used to model the signal selection efficiency. The signal samples were generated via the gluon-gluon fusion process using POWHEG-BOX v2 [64–66], with the CT10 next-to-leading order (NLO) parton distribution function (PDF) set [67]. Particle decays, hadronization, parton showers, and the underlying event were modeled using PYTHIA v8.212 [68] and Evt Gen v1.6.0 [69], interfaced to the AZNLO [70] set of tuned parameters and the CTEQ6 L1 PDF set [71]. Next-to-next-to-leading order (NNLO) corrections are applied to the p_T distribution of the Higgs boson. The a branching fractions were determined using PYTHIA 8 [68] with a 2HDM $\tan\beta$ value of 1, which predicts $a \rightarrow gg$ to be the dominant decay mode until $a \rightarrow c\bar{c}$ becomes kinematically accessible. The signal MC samples used in this analysis

have a masses of 0.5, 0.75, 1, 1.5, 2, 2.5, 3, 3.5, and 4 GeV. The Z boson is required to decay into pairs of electrons, muons, or τ leptons.

The background is dominated by Z + jets events, modeled using SHERPA 2.2.1 [72] interfaced to the NNPDF 3.0 (NNLO) PDF set [73]. The inclusive production cross sections are known to NNLO in QCD [74]. The ZZ , ZW , and $t\bar{t}$ processes contribute $< 1\%$ of the total background in this search. The diboson backgrounds were modeled using SHERPA 2.2.1 interfaced to the NNPDF 3.0 (NNLO) PDF set, except for gluon-induced ZZ production, which was modeled using SHERPA 2.2.2 [72]. All of the SHERPA samples used a set of tuned parameters developed by the SHERPA authors. The $t\bar{t}$ process was modeled using POWHEG-BOX v2, while the subsequent decay, hadronization, parton shower, and underlying event were modeled using PYTHIA v8.230 and EvtGen v1.6.0. The NNPDF 2.3 (LO) PDF set [75] and the A14 set of tuned parameters [76] were used.

The simulation of the ATLAS detector [77] in GEANT4 [78] was used to model the interaction of particles with the detector in all the MC samples. Data-driven corrections are applied to the event-level trigger efficiencies, the jet vertex tagging efficiency [79], the electron [80] reconstruction, identification, and isolation efficiencies, and the muon [81] reconstruction, isolation, and track-to-vertex association efficiencies.

Events are selected by a combination of single electron or muon triggers for each data-taking period [82–85], and the online lepton reconstructed by the trigger is required to be within $\Delta R = 0.1$ [86] of an off-line reconstructed lepton. Events are required to have at least one reconstructed primary interaction vertex [87]. Electron candidates are reconstructed by matching tracks in the inner detector to topological energy clusters in the electromagnetic calorimeter [80] and must pass a likelihood-based selection, which requires the shower profile to be compatible with that of an electromagnetic shower. Muons are reconstructed using tracks in the muon spectrometer, matched to tracks in the inner detector where available [88]. Electrons and muons are each required to have $p_T > 18$ GeV, and at least one must have $p_T > 27$ GeV. Electrons (muons) are required to be reconstructed within $|\eta| < 2.47$ ($|\eta| < 2.7$), but electrons within $1.37 < |\eta| < 1.52$ are excluded. The transverse energy sum in a cone of size $\Delta R = 0.2$ around the electron [muon] in the calorimeter must be less than 20% (30%) of the lepton's p_T , and the summed p_T of tracks within a cone of variable size $\Delta R = \min(0.2, 10 \text{ GeV}/p_T)$ [$\Delta R = \min(0.15, 10 \text{ GeV}/p_T)$] around the electron [muon] must be less than 15% of its p_T . Contributions from nearby electrons and muons are removed from these cones. If an inner detector track is present, muons must also have a longitudinal impact parameter $|z_0 \sin\theta| < 0.5$ mm and a transverse impact parameter $|d_0| < 1$ mm relative to the primary interaction vertex. At least two same-flavor

opposite-sign electrons or muons are required to pass this selection and have an invariant mass compatible with the mass of the Z boson: $81 < m_{\ell\ell} < 101$ GeV. If multiple same-flavor opposite-sign lepton pairs fulfill this requirement, the pairing with an invariant mass closest to that of the Z boson is chosen. $Z \rightarrow \tau\tau$ decays are reconstructed through the leptonic decays of the τ leptons.

The hadronically decaying resonance is reconstructed as a single jet using the anti- k_r jet algorithm [89,90] with a radius parameter of 0.4, formed from topological calorimeter energy clusters [91,92] and calibrated to the electromagnetic energy scale. Jet energies are corrected for contributions from simultaneous inelastic pp interactions (pileup) using a jet-area-based technique [93,94] and calibrated [95,96] using p_T - and η -dependent correction factors determined from simulation, with residual corrections from *in situ* measurements applied to data and internal jet properties. Jets are required to have $p_T > 20$ GeV and $|\eta| < 2.5$ and satisfy a jet cleaning requirement [97]. To reject jets from pileup interactions, jets with $p_T < 60$ GeV and $|\eta| < 2.4$ are required to pass a “jet vertex tagger” [79] requirement. An overlap removal procedure resolves cases in which multiple electrons, muons, or jets are reconstructed from the same detector signature. Higgs boson candidates are reconstructed from the lepton pair and jet system, which is required to have an invariant mass passing a loose preselection requirement: $m_{\ell\ell j} < 250$ GeV. If multiple jets satisfy these requirements, the jet with the highest p_T is selected. The acceptance for this preselection,

evaluated using generator-level MC samples, varies between 28% and 29% for the different Q/a signal hypotheses.

MLPs [63] are used to select signal events passing this preselection. The MLP input variables are built using tracks matched to the calorimeter jet by ghost association [93], in which the tracks are included in the jet clustering process as with negligible energy and their angles from the jet axis. This allows the MLP to benefit from the high resolution of the tracking detector. These tracks must have $p_T > 500$ MeV and $|\eta| < 2.5$ and pass loose quality and track-to-vertex association requirements [98] to reject fake tracks from the reconstruction and tracks from pileup, respectively. Six dimensionless variables are constructed using these tracks: the ratio of the p_T of the highest p_T track to the p_T of the ghost-associated track system; the angular separation ΔR between the highest- p_T track and the calorimeter jet axis; NSubJettiness 2 [99], using exclusive- k_r subjet axes with radius parameters of 0.2, and a jet axis radius parameter of 0.4; angularity(2) [100]; and $U_1(0.7)$ and $M_2(0.3)$, which are modified energy correlation functions [101] designed for quark-gluon discrimination and to target two-pronged substructure, respectively. These variables primarily capitalize on the presence of a narrow resonance or two-pronged substructure in the track system. Initially, a regression MLP [63], using four hidden layers of 12 nodes, is trained using the above input variables and the a signal samples to estimate the mass of a , as shown in Fig. 1(a). This estimated mass is then

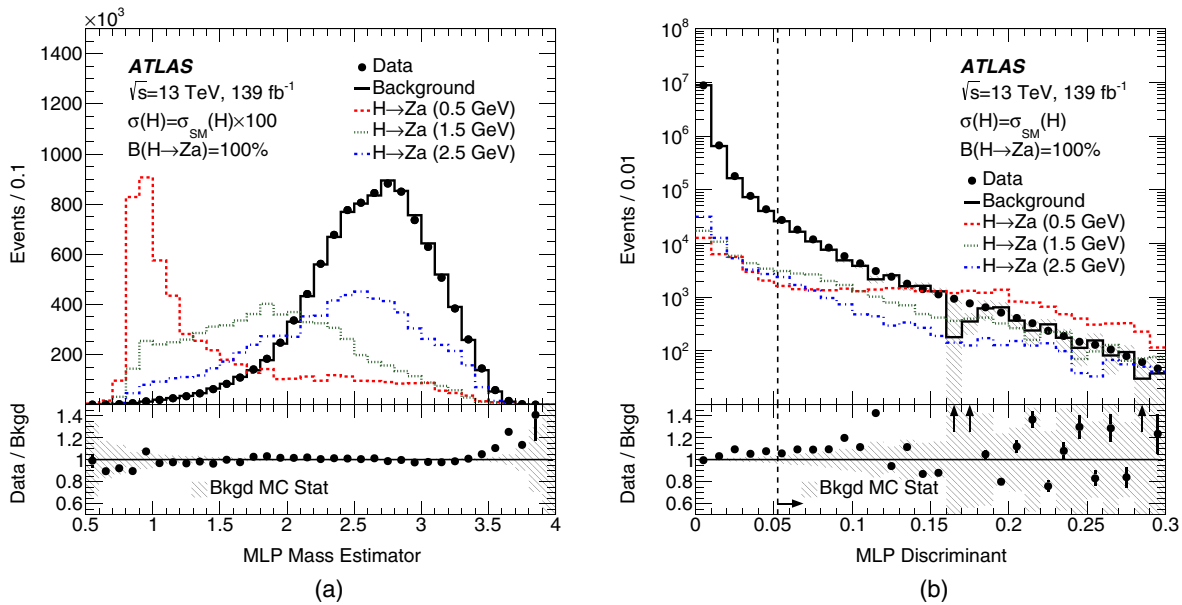


FIG. 1. Output of (a) the regression and (b) the classification MLPs, for data, background, and three signal hypotheses. Events are required to pass the complete event selection, including the $120 < m_{\ell\ell j} < 135$ GeV requirement, but not the requirement on the classification MLP output variable. The background normalization is set equal to that of the data, and the signal normalizations assume the SM Higgs boson inclusive production cross section and $\mathcal{B}(H \rightarrow Za) = 100\%$, and in (a) the signal normalization is scaled up by a factor of 100. The error bars (hatched regions) represent the data (MC) sample statistical uncertainty, in both the histograms and the ratio plots. In (b) the region to the right of the dashed line is the signal region.

provided alongside the six input variables to a *classification* MLP [63], to inform the classifier about the part of the hadronic resonance mass spectrum where the specific event lies. This classification MLP has two hidden layers of six and five nodes and is trained using the a signal samples and the background samples. The 0.75 GeV a signal sample is excluded from the training of the classification MLP to ensure an even spacing between the a mass hypotheses, so the training is not biased toward lower masses. Both MLPs use sigmoidal response functions with summed inputs and are trained using backpropagation with a mean-square estimator [63], as these resulted in optimal discrimination without overtraining. The addition of the regression MLP was found to result in about a 13% improvement in the S/\sqrt{B} of the classification MLP, where S and B are the expected numbers of signal and background events passing the MLP requirement, respectively. The classification MLP output variable (M) is shown in Fig. 1(b).

The signal region (SR) for this search is defined by the requirements $120 < m_{\ell\ell j} < 135$ GeV and $M > 0.0524$, chosen to maximize the expected S/\sqrt{B} , averaged over the various a mass hypotheses. The efficiency of this MLP requirement for events passing the preselection is $(0.761 \pm 0.020)\%$ for the background, $(5.89 \pm 0.24)\%$ and $(6.66 \pm 0.26)\%$ for $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$, respectively, and between $(1.88 \pm 0.15)\%$ and $(45.9 \pm 0.8)\%$ for $H \rightarrow Za$. The efficiencies for the complete selection are estimated using MC samples and are $(0.545 \pm 0.022)\%$ and $(0.560 \pm 0.022)\%$ for $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$, respectively, and range between $(0.140 \pm 0.011)\%$ and $(3.27 \pm 0.06)\%$ for $H \rightarrow Za$. The efficiencies are highest for the lowest a mass hypotheses, due to higher probabilities to pass the MLP requirement. The efficiency for $H \rightarrow Z\eta_c$ events to pass the MLP requirement is lower than that of $H \rightarrow ZJ/\psi$ events, as J/ψ decays tend to have a lower charged hadron multiplicity. Using the predicted cross section for inclusive SM Higgs boson production of $55.7^{+3.0}_{-3.9}$ pb [102], and $\mathcal{B}[H \rightarrow Z(Q/a)] = 100\%$, gives expected signal yields of 4260 and 4370 for $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$, respectively, and between 1090 and 25600 for $H \rightarrow Za$.

A “modified ABCD estimate” of the total background in the SR is derived using four regions: A , defined by $0.0341 < M < 0.0524$, expected to contain about 10% of the total background, and $155 < m_{\ell\ell j} < 175$ GeV; B , defined by the $m_{\ell\ell j}$ requirement of the SR and the M requirement of region A ; C , defined by the M requirement of the SR and the $m_{\ell\ell j}$ requirement of A ; and D , which is the SR. An initial data-driven background estimate in the SR is calculated as $D = BC/A$, then MC samples, reweighted to match data, are used to correct this estimate for the 13% correlation between the $m_{\ell\ell j}$ and M variables. This reweighting is performed in the p_T of the calorimeter jet, the number of ghost-associated tracks and $U_1(0.7)$. This background estimate is 82400 ± 2900 events in the SR, where the uncertainty is due to the limited data and MC

sample statistics. The background estimation method is found to be consistent with data within 1.7 times the total statistical and systematic uncertainty in 14 validation regions, defined in regions of the $m_{\ell\ell j}$ and M variables.

A measure of $\sigma(pp \rightarrow H)\mathcal{B}[H \rightarrow Z(Q/a)]$ is extracted for a given signal hypothesis using a maximum-likelihood fit [103] to the number of events observed in the SR. The systematic uncertainties are included in the likelihood fit as nuisance parameters, which modify the signal efficiencies or the simulation-based correction used to calculate the expected background yield. These systematic uncertainties include uncertainties in the signal and background modeling and experimental uncertainties. The sources of modeling uncertainty include the limited MC sample statistics, renormalization scale and choice of MC generator for the signal and background, and a signal uncertainty to account for the extrapolation from gluon-gluon fusion signal samples to the inclusive Higgs boson production cross section. The effects of factorization scale and PDF uncertainties are found to be negligible. The experimental uncertainties considered are due to the luminosity [104], pileup [105], triggers, lepton [81,106,107], and jet [96] reconstruction. The total uncertainty on the extracted signal yield is dominated by the background modeling uncertainties, the largest being due to limited MC sample statistics. The total uncertainty on the background in the SR is 3700 events, where the uncertainty due to the limited data and MC sample statistics is 2900 and the modeling uncertainty is 2300. The data statistical uncertainty corresponds to approximately 8% of the total uncertainty on the extracted signal yield.

The SR contains 82 908 data events. This result is compatible with the SM background-only expectation, and the three-body mass distribution is shown in Fig. 2. Upper limits at 95% confidence level (CL) are set on $\sigma(pp \rightarrow H)\mathcal{B}[H \rightarrow Z(Q/a)]$ for the various signal hypotheses, using the profile-likelihood test statistic [103] and the CLs technique [108]. The observed (expected) upper limits for the $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$ hypotheses are 110 and 100 pb (100^{+40}_{-30} and 100^{+40}_{-30} pb), respectively, while the upper limits for the $H \rightarrow Za$ signal hypotheses are given in Table I. In the absence of systematic uncertainties, these limits would range between 1.9 and 55 pb for the different signal hypotheses. To simplify the interpretation, the upper limits are quoted for $\mathcal{B}(a \rightarrow gg) = 100\%$ and $\mathcal{B}(a \rightarrow s\bar{s}) = 100\%$. Because of the Yukawa ordering of the decays of Higgs bosons, only decays into gluon and strange quark pairs are considered. The tighter limits for the $a \rightarrow s\bar{s}$ decays are due to a higher MLP selection efficiency. The systematic uncertainties for $a \rightarrow gg$ and $a \rightarrow s\bar{s}$ decay hypotheses are estimated using the inclusive decays as modeled in PYTHIA 8, which is a good approximation due to the dominance of the background modeling

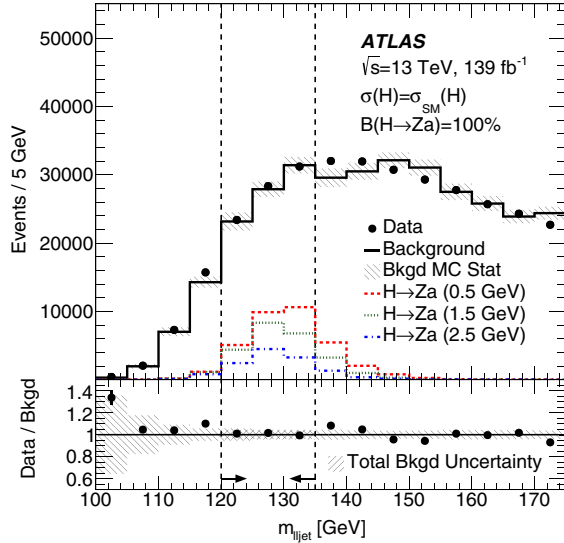


FIG. 2. Invariant mass of the lepton pair plus jet system, for data, background, and three signal hypotheses. Events are required to pass the complete event selection, including the MLP output variable requirement, but not the $120 < m_{\ell\ell j} < 135$ GeV requirement. The background normalization is defined by the background estimate in the signal region, and the signal normalizations assume the SM Higgs boson inclusive production cross section and $\mathcal{B}(H \rightarrow Za) = 100\%$. The error bars (hatched regions) represent the data (MC) sample statistical uncertainty, in both the histograms and the ratio plots. The region between the vertical dashed lines is the signal region. The total background uncertainty in the signal region is also indicated.

uncertainties. This assumption allows a limit to be set on the decay of a into gg or $s\bar{s}$ final states, in any ratio, by using a weighted sum of the two limits. Higgs boson decays to a Z boson and a quarkonium state other than the considered signal process are not included in the statistical interpretation.

TABLE I. Expected (Exp) and observed (Obs) 95% CL upper limits on $\sigma(pp \rightarrow H)\mathcal{B}(H \rightarrow Za)$ /pb. These results are quoted for $\mathcal{B}(a \rightarrow gg) = 100\%$ and $\mathcal{B}(a \rightarrow s\bar{s}) = 100\%$ for each signal sample.

a mass (GeV)	$a \rightarrow gg$		$a \rightarrow s\bar{s}$	
	Exp	Obs	Exp	Obs
0.5	16^{+6}_{-5}	17		
0.75	19^{+7}_{-5}	20		
1.0	17^{+7}_{-5}	18		
1.5	20^{+8}_{-6}	22	19^{+7}_{-5}	20
2.0	26^{+10}_{-7}	27	23^{+9}_{-6}	24
2.5	38^{+15}_{-11}	40	32^{+12}_{-9}	33
3.0	75^{+29}_{-21}	78	65^{+25}_{-18}	68
3.5	110^{+40}_{-30}	120		
4.0	320^{+130}_{-90}	340		

In conclusion, a search has been performed for Higgs boson decays into a Z boson and either a η_c or J/ψ charmonium state, or a light spin-0 boson. No excess is found, and 95% CL upper limits are set on $\sigma(pp \rightarrow H)\mathcal{B}[H \rightarrow Z(Q/a)]$, with values of 110 and 100 pb for the $H \rightarrow Z\eta_c$ and $H \rightarrow ZJ/\psi$ hypotheses, respectively, and with values in the range 17–340 pb for the $H \rightarrow Za$ signal hypotheses. Assuming the SM prediction for inclusive Higgs boson production, the limits on charmonium decay modes correspond to branching fraction limits in excess of 100%. This is the first direct limit on decays of the observed Higgs boson to light scalars, decaying to light quarks or gluons. Because of the large value of $\mathcal{B}(a \rightarrow \text{hadrons})$ over the entire 2HDM(+S) parameter space, these limits represent tight, direct constraints for low (high) $\tan\beta$ in the type-II and type-III (type-VI) 2HDM + S [109].

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Asawatavonvanich,¹⁶⁴ N. Asbah,⁵⁹ E. M. Asimakopoulou,¹⁷¹ L. Asquith,¹⁵⁵ J. Assahsah,^{35d} K. Assamagan,²⁹ R. Aсталos,^{28a} R. J. Atkin,^{33a} M. Atkinson,¹⁷² N. B. Atlay,¹⁹ H. Atmani,⁶⁵ K. Augsten,¹⁴¹ V. A. Austrup,¹⁸¹ G. Avolio,³⁶ M. K. Ayoub,^{15a} G. Azuelos,^{110,d} H. Bachacou,¹⁴⁴ K. Bachas,¹⁶¹ M. Backes,¹³⁴ F. Backman,^{45a,45b} P. Bagnaia,^{73a,73b} M. Bahmani,⁸⁵ H. Bahrasemani,¹⁵¹ A. J. Bailey,¹⁷³ V. R. Bailey,¹⁷² J. T. Baines,¹⁴³ C. Bakalis,¹⁰ O. K. Baker,¹⁸² P. J. Bakker,¹²⁰ E. Bakos,¹⁶ D. Bakshi Gupta,⁸ S. Balaji,¹⁵⁶ E. M. Baldin,^{122b,122a} P. Balek,¹⁷⁹ F. Balli,¹⁴⁴ W. K. Balunas,¹³⁴ J. Balz,¹⁰⁰ E. Banas,⁸⁵ M. Bandieramonte,¹³⁸ A. Bandyopadhyay,²⁴ Sw. Banerjee,^{180,e} L. Barak,¹⁶⁰ W. M. Barbe,³⁸ E. L. Barberio,¹⁰⁵ D. Barberis,^{55b,55a} M. Barbero,¹⁰² G. Barbour,⁹⁵ T. Barillari,¹¹⁵ M-S. Barisits,³⁶ J. Barkeloo,¹³¹ T. Barklow,¹⁵² R. Barnea,¹⁵⁹ B. M. Barnett,¹⁴³ R. M. Barnett,¹⁸ Z. Barnovska-Blenessy,^{60a} A. Baroncelli,^{60a} G. Barone,²⁹ A. J. Barr,¹³⁴ L. Barranco Navarro,^{45a,45b} F. Barreiro,⁹⁹ J. Barreiro Guimarães da Costa,^{15a} U. Barron,¹⁶⁰ S. Barsov,¹³⁷ F. Bartels,^{61a} R. Bartoldus,¹⁵² G. Bartolini,¹⁰² A. E. Barton,⁹⁰ P. Bartos,^{28a} A. Basalaeu,⁴⁶ A. Basan,¹⁰⁰ A. Bassalat,^{65,f} M. J. Basso,¹⁶⁶ R. L. Bates,⁵⁷ S. Batlamous,^{35e} J. R. Batley,³² B. Batool,¹⁵⁰ M. Battaglia,¹⁴⁵ M. Bause,^{73a,73b} F. Bauer,¹⁴⁴ K. T. Bauer,¹⁷⁰ H. S. Bawa,³¹ J. B. Beacham,⁴⁹ T. Beau,¹³⁵ P. H. Beauchemin,¹⁶⁹ F. Becherer,⁵² P. Bechtel,²⁴ H. C. Beck,⁵³ H. P. Beck,^{20,g} K. Becker,¹⁷⁷ C. Becot,⁴⁶ A. Beddall,^{12d} A. J. Beddall,^{12a} V. A. Bednyakov,⁸⁰ M. Bedognetti,¹²⁰ C. P. Bee,¹⁵⁴ T. A. Beermann,¹⁸¹ M. Begalli,^{81b} M. Begel,²⁹ A. Behera,¹⁵⁴ J. K. Behr,⁴⁶ F. Beisiegel,²⁴ M. Belfkir,⁵ A. S. Bell,⁹⁵ G. Bella,¹⁶⁰ L. Bellagamba,^{23b} A. Bellerive,³⁴ P. Bellos,⁹ K. Beloborodov,^{122b,122a} K. Belotskiy,¹¹² N. L. Belyaev,¹¹² D. Benckekroun,^{35a} N. Benekos,¹⁰ Y. Benhammou,¹⁶⁰ D. P. Benjamin,⁶ M. Benoit,⁵⁴ J. R. Bensinger,²⁶ S. Bentvelsen,¹²⁰ L. Beresford,¹³⁴ M. Beretta,⁵¹ D. Berge,¹⁹ E. Bergeas Kuutmann,¹⁷¹ N. Berger,⁵ B. Bergmann,¹⁴¹ L. J. Bergsten,²⁶ J. Beringer,¹⁸ S. Berlendis,⁷ G. Bernardi,¹³⁵ C. Bernius,¹⁵² F. U. Bernlochner,²⁴ T. Berry,⁹⁴ P. Berta,¹⁰⁰ C. Bertella,^{15a} A. Berthold,⁴⁸ I. A. Bertram,⁹⁰ O. Bessidskaia Bylund,¹⁸¹ N. Besson,¹⁴⁴ A. Bethani,¹⁰¹ S. Bethke,¹¹⁵ A. Betti,⁴² A. J. Bevan,⁹³ J. Beyer,¹¹⁵ D. S. Bhattacharya,¹⁷⁶ P. Bhattarai,²⁶ V. S. Bhopatkar,⁶ R. Bi,¹³⁸ R. M. Bianchi,¹³⁸ O. Biebel,¹¹⁴ D. Biedermann,¹⁹ R. Bielski,³⁶ K. Bierwagen,¹⁰⁰ N. V. Biesuz,^{72a,72b} M. Biglietti,^{75a} T. R. V. Billoud,¹¹⁰ M. Bindi,⁵³ A. Bingul,^{12d} C. Bini,^{73a,73b} S. Biondi,^{23b,23a} C. J. Birch-sykes,¹⁰¹ M. Birman,¹⁷⁹ T. Bisanz,³⁶ J. P. Biswal,³ D. Biswas,^{180,e} A. Bitadze,¹⁰¹ C. Bittrich,⁴⁸ K. Björke,¹³³ T. Blazek,^{28a} I. Bloch,⁴⁶ C. Blocker,²⁶ A. Blue,⁵⁷ U. Blumenschein,⁹³ G. J. Bobbink,¹²⁰ V. S. Bobrovnikov,^{122b,122a} S. S. Bocchetta,⁹⁷ D. Bogavac,¹⁴ A. G. Bogdanchikov,^{122b,122a} C. Bohm,^{45a} V. 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W. K. Brooks,^{146d} E. Brost,²⁹ P. A. Bruckman de Renstrom,⁸⁵ B. Brüers,⁴⁶ D. Bruncko,^{28b} A. Bruni,^{23b} G. Bruni,^{23b} L. S. Bruni,¹²⁰ S. Bruno,^{74a,74b} M. Bruschi,^{23b} N. Bruscolo,^{73a,73b} L. Bryngemark,¹⁵² T. Buanes,¹⁷ Q. Buat,³⁶ P. Buchholz,¹⁵⁰ A. G. Buckley,⁵⁷ I. A. Budagov,⁸⁰ M. K. Bugge,¹³³ F. Bühner,⁵² O. Bulekov,¹¹² B. A. Bullard,⁵⁹ T. J. Burch,¹²¹ S. Burdin,⁹¹ C. D. Burgard,¹²⁰ A. M. Burger,¹²⁹ B. Burghgrave,⁸ J. T. P. Burr,⁴⁶ C. D. Burton,¹¹ J. C. Burzynski,¹⁰³ V. Büscher,¹⁰⁰ E. Buschmann,⁵³ P. J. Bussey,⁵⁷ J. M. Butler,²⁵ C. M. Buttar,⁵⁷ J. M. Butterworth,⁹⁵ P. Butti,³⁶ W. Buttinger,³⁶ C. J. Buxo Vazquez,¹⁰⁷ A. Buzatu,¹⁵⁷ A. R. Buzykaev,^{122b,122a} G. Cabras,^{23b,23a} S. Cabrera Urbán,¹⁷³ D. Caforio,⁵⁶ H. Cai,¹³⁸ V. M. M. Cairo,¹⁵² O. Cakir,^{4a} N. Calace,³⁶ P. Calafiura,¹⁸ G. Calderini,¹³⁵ P. Calfayan,⁶⁶ G. Callea,⁵⁷ L. P. Caloba,^{81b} A. Caltabiano,^{74a,74b} S. Calvente Lopez,⁹⁹ D. Calvet,³⁸ S. Calvet,³⁸ T. P. Calvet,¹⁰² M. Calvetti,^{72a,72b} R. 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L. Dufloc, ⁶⁵ M. Dührssen, ³⁶ C. Dülsen, ¹⁸¹ M. Dumancic, ¹⁷⁹ A. E. Dumitriu, ^{27b} A. K. Duncan, ⁵⁷ M. Dunford, ^{61a}
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G. Gaudio, ^{71a} I. L. Gavrilenko, ¹¹¹ A. Gavriilyuk, ¹²⁴ C. Gay, ¹⁷⁴ G. Gaycken, ⁴⁶ E. N. Gazis, ¹⁰ A. A. Geanta, ^{27b} C. M. Gee, ¹⁴⁵
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M. Ghneimat, ¹⁵⁰ A. Ghosh, ⁶⁵ A. Ghosh, ⁷⁸ B. Giacobbe, ^{23b} S. Giagu, ^{73a,73b} N. Giangiacomi, ^{23b,23a} P. Giannetti, ^{72a}
A. Giannini, ^{70a,70b} G. Giannini, ¹⁴ S. M. Gibson, ⁹⁴ M. Gignac, ¹⁴⁵ D. T. Gil, ^{84b} D. Gillberg, ³⁴ G. Gilles, ¹⁸¹ D. M. Gingrich, ^{3,d}
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L. Han, ^{60a} S. Han, ¹⁸ Y. F. Han, ¹⁶⁶ K. Hanagaki, ^{82,v} M. Hance, ¹⁴⁵ D. M. Handl, ¹¹⁴ M. D. Hank, ³⁷ R. Hankache, ¹³⁵
E. Hansen, ⁹⁷ J. B. Hansen, ⁴⁰ J. D. Hansen, ⁴⁰ M. C. Hansen, ²⁴ P. H. Hansen, ⁴⁰ E. C. Hanson, ¹⁰¹ K. Hara, ¹⁶⁸ T. Harenberg, ¹⁸¹
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D. Hayden, ¹⁰⁷ C. Hayes, ¹⁰⁶ R. L. Hayes, ¹⁷⁴ C. P. Hays, ¹³⁴ J. M. Hays, ⁹³ H. S. Hayward, ⁹¹ S. J. Haywood, ¹⁴³ F. He, ^{60a}

M. P. Heath,⁵⁰ V. Hedberg,⁹⁷ S. Heer,²⁴ A. L. Heggelund,¹³³ C. Heidegger,⁵² K. K. Heidegger,⁵² W. D. Heidorn,⁷⁹ J. Heilman,³⁴ S. Heim,⁴⁶ T. Heim,¹⁸ B. Heinemann,^{46,w} J. G. Heinlein,¹³⁶ J. J. Heinrich,¹³¹ L. Heinrich,³⁶ J. Hejbal,¹⁴⁰ L. Helary,^{61b} A. Held,¹²⁵ S. Hellesund,¹³³ C. M. Helling,¹⁴⁵ S. Hellman,^{45a,45b} C. Helsens,³⁶ R. C. W. Henderson,⁹⁰ Y. Heng,¹⁸⁰ L. Henkelmann,³² A. M. Henriques Correia,³⁶ H. Herde,²⁶ Y. Hernández Jiménez,^{33e} H. Herr,¹⁰⁰ M. G. Herrmann,¹¹⁴ T. Herrmann,⁴⁸ G. Hertzen,⁵² R. Hertzenberger,¹¹⁴ L. Hervas,³⁶ T. C. Herwig,¹³⁶ G. G. Hesketh,⁹⁵ N. P. Hessey,^{167a} H. Hibi,⁸³ A. Higashida,¹⁶² S. Higashino,⁸² E. Higón-Rodríguez,¹⁷³ K. Hildebrand,³⁷ J. C. Hill,³² K. K. Hill,²⁹ K. H. Hiller,⁴⁶ S. J. Hillier,²¹ M. Hils,⁴⁸ I. Hinchliffe,¹⁸ F. Hinterkeuser,²⁴ M. Hirose,¹³² S. Hirose,⁵² D. Hirschbuehl,¹⁸¹ B. Hiti,⁹² O. Hladik,¹⁴⁰ D. R. Hlaluku,^{33e} J. Hobbs,¹⁵⁴ N. Hod,¹⁷⁹ M. C. Hodgkinson,¹⁴⁸ A. Hoecker,³⁶ D. Hohn,⁵² D. Hohov,⁶⁵ T. Holm,²⁴ T. R. Holmes,³⁷ M. Holzbock,¹¹⁴ L. B. A. H. Hommels,³² T. M. Hong,¹³⁸ J. C. Honig,⁵² A. Hönle,¹¹⁵ B. H. Hooberman,¹⁷² W. H. Hopkins,⁶ Y. Horii,¹¹⁷ P. Horn,⁴⁸ L. A. Horyn,³⁷ S. Hou,¹⁵⁷ A. Hoummada,^{35a} J. Howarth,⁵⁷ J. Hoya,⁸⁹ M. Hrabovsky,¹³⁰ J. Hrdinka,⁷⁷ J. Hrivnac,⁶⁵ A. Hrynevich,¹⁰⁹ T. Hryn'ova,⁵ P. J. Hsu,⁶⁴ S.-C. Hsu,¹⁴⁷ Q. Hu,²⁹ S. Hu,^{60c} Y. F. Hu,^{15a,15d,x} D. P. Huang,⁹⁵ Y. Huang,^{60a} Y. Huang,^{15a} Z. Hubacek,¹⁴¹ F. Hubaut,¹⁰² M. Huebner,²⁴ F. Huegging,²⁴ T. B. Huffman,¹³⁴ M. Huhtinen,³⁶ R. Hulsken,⁵⁸ R. F. H. Hunter,³⁴ P. Huo,¹⁵⁴ N. Huseynov,^{80,y} J. Huston,¹⁰⁷ J. Huth,⁵⁹ R. Hyneman,¹⁰⁶ S. Hyrych,^{28a} G. Iacobucci,⁵⁴ G. Iakovidis,²⁹ I. Ibragimov,¹⁵⁰ L. Iconomidou-Fayard,⁶⁵ P. Iengo,³⁶ R. Ignazzi,⁴⁰ O. Igonkina,^{120,a,z} R. Iguchi,¹⁶² T. Iizawa,⁵⁴ Y. Ikegami,⁸² M. Ikeno,⁸² D. Iliadis,¹⁶¹ N. Ilic,^{119,166,l} F. Iltzsche,⁴⁸ H. Imam,^{35a} G. Introzzi,^{71a,71b} M. Iodice,^{75a} K. Iordanidou,^{167a} V. Ippolito,^{73a,73b} M. F. Isacson,¹⁷¹ M. Ishino,¹⁶² W. Islam,¹²⁹ C. Issever,^{19,46} S. Istin,¹⁵⁹ F. Ito,¹⁶⁸ J. M. Iturbe Ponce,^{63a} R. Iuppa,^{76a,76b} A. Ivina,¹⁷⁹ H. Iwasaki,⁸² J. M. Izen,⁴³ V. Izzo,^{70a} P. Jacka,¹⁴⁰ P. Jackson,¹ R. M. Jacobs,⁴⁶ B. P. Jaeger,¹⁵¹ V. Jain,² G. Jäkel,¹⁸¹ K. B. Jakobi,¹⁰⁰ K. Jakobs,⁵² T. Jakoubek,¹⁷⁹ J. Jamieson,⁵⁷ K. W. Janas,^{84a} R. Jansky,⁵⁴ M. Janus,⁵³ P. A. Janus,^{84a} G. Jarlskog,⁹⁷ A. E. Jaspan,⁹¹ N. Javadov,^{80,y} T. Javůrek,³⁶ M. Javurkova,¹⁰³ F. Jeanneau,¹⁴⁴ L. Jeanty,¹³¹ J. Jejelava,^{158a} P. Jenni,^{52,aa} N. Jeong,⁴⁶ S. Jézéquel,⁵ H. Ji,¹⁸⁰ J. Jia,¹⁵⁴ H. Jiang,⁷⁹ Y. Jiang,^{60a} Z. Jiang,¹⁵² S. Jiggins,⁵² F. A. Jimenez Morales,³⁸ J. Jimenez Pena,¹¹⁵ S. Jin,^{15c} A. Jinaru,^{27b} O. Jinnouchi,¹⁶⁴ H. Jivan,^{33e} P. Johansson,¹⁴⁸ K. A. Johns,⁷ C. A. Johnson,⁶⁶ R. W. L. Jones,⁹⁰ S. D. Jones,¹⁵⁵ T. J. Jones,⁹¹ J. Jongmanns,^{61a} J. Jovicevic,³⁶ X. Ju,¹⁸ J. J. Junggeburth,¹¹⁵ A. Juste Rozas,^{14,t} A. Kaczmarska,⁸⁵ M. Kado,^{73a,73b} H. Kagan,¹²⁷ M. Kagan,¹⁵² A. Kahn,³⁹ C. Kahra,¹⁰⁰ T. Kaji,¹⁷⁸ E. Kajomovitz,¹⁵⁹ C. W. Kalderon,²⁹ A. Kaluza,¹⁰⁰ A. Kamenshchikov,¹²³ M. Kaneda,¹⁶² N. J. Kang,¹⁴⁵ S. Kang,⁷⁹ Y. Kano,¹¹⁷ J. Kanzaki,⁸² L. S. Kaplan,¹⁸⁰ D. Kar,^{33e} K. Karava,¹³⁴ M. J. Kareem,^{167b} I. Karkanias,¹⁶¹ S. N. Karpov,⁸⁰ Z. M. Karpova,⁸⁰ V. Kartvelishvili,⁹⁰ A. N. Karyukhin,¹²³ A. Kastanas,^{45a,45b} C. Kato,^{60d,60c} J. Katzy,⁴⁶ K. Kawade,¹⁴⁹ K. Kawagoe,⁸⁸ T. Kawaguchi,¹¹⁷ T. Kawamoto,¹⁴⁴ G. Kawamura,⁵³ E. F. Kay,¹⁷⁵ S. Kazakos,¹⁴ V. F. Kazanin,^{122b,122a} R. Keeler,¹⁷⁵ R. Kehoe,⁴² J. S. Keller,³⁴ E. Kellermann,⁹⁷ D. Kelsey,¹⁵⁵ J. J. Kempster,²¹ J. Kendrick,²¹ K. E. Kennedy,³⁹ O. Kepka,¹⁴⁰ S. Kersten,¹⁸¹ B. P. Kerševan,⁹² S. Ketabchi Haghighat,¹⁶⁶ M. Khader,¹⁷² F. Khalil-Zada,¹³ M. Khandoga,¹⁴⁴ A. Khanov,¹²⁹ A. G. Kharlamov,^{122b,122a} T. Kharlamova,^{122b,122a} E. E. Khoda,¹⁷⁴ A. Khodinov,¹⁶⁵ T. J. Khoo,⁵⁴ G. Khoraiuli,¹⁷⁶ E. Khramov,⁸⁰ J. Khubua,^{158b} S. Kido,⁸³ M. Kiehn,⁵⁴ C. R. Kilby,⁹⁴ E. Kim,¹⁶⁴ Y. K. Kim,³⁷ N. Kimura,⁹⁵ B. T. King,^{91,a} A. Kirchhoff,⁵³ D. Kirchmeier,⁴⁸ J. Kirk,¹⁴³ A. E. Kiryunin,¹¹⁵ T. Kishimoto,¹⁶² D. P. Kisliuk,¹⁶⁶ V. Kitali,⁴⁶ C. Kitsaki,¹⁰ O. Kivernyk,²⁴ T. Klapdor-Kleingrothaus,⁵² M. Klassen,^{61a} C. Klein,³⁴ M. H. Klein,¹⁰⁶ M. Klein,⁹¹ U. Klein,⁹¹ K. Kleinknecht,¹⁰⁰ P. Klimek,¹²¹ A. Klimentov,²⁹ T. Klingl,²⁴ T. Klioutchnikova,³⁶ F. F. Klitzner,¹¹⁴ P. Kluit,¹²⁰ S. Kluth,¹¹⁵ E. Kneringer,⁷⁷ E. B. F. G. Knoops,¹⁰² A. Knue,⁵² D. Kobayashi,⁸⁸ T. Kobayashi,¹⁶² M. Kobel,⁴⁸ M. Kocian,¹⁵² T. Kodama,¹⁶² P. Kodys,¹⁴² D. M. Koeck,¹⁵⁵ P. T. Koenig,²⁴ T. Koffas,³⁴ N. M. Köhler,³⁶ M. Kolb,¹⁴⁴ I. Koletsou,⁵ T. Komarek,¹³⁰ T. Kondo,⁸² K. Köneke,⁵² A. X. Y. Kong,¹ A. C. König,¹¹⁹ T. Kono,¹²⁶ V. Konstantinides,⁹⁵ N. Konstantinidis,⁹⁵ B. Konya,⁹⁷ R. Kopeliansky,⁶⁶ S. Koperny,^{84a} K. Korcyl,⁸⁵ K. Kordas,¹⁶¹ G. Koren,¹⁶⁰ A. Korn,⁹⁵ I. Korolkov,¹⁴ E. V. Korolkova,¹⁴⁸ N. Korotkova,¹¹³ O. Kortner,¹¹⁵ S. Kortner,¹¹⁵ V. V. Kostyukhin,^{148,165} A. Kotskechagia,⁶⁵ A. Kotwal,⁴⁹ A. Koulouris,¹⁰ A. Kourkoumeli-Charalampidi,^{71a,71b} C. Kourkoumelis,⁹ E. Kourlitis,⁶ V. Kouskoura,²⁹ R. Kowalewski,¹⁷⁵ W. Kozanecki,¹⁰¹ A. S. Kozhin,¹²³ V. A. Kramarenko,¹¹³ G. Kramberger,⁹² D. Krasnopevtsev,^{60a} M. W. Krasny,¹³⁵ A. Krasznahorkay,³⁶ D. Krauss,¹¹⁵ J. A. Kremer,¹⁰⁰ J. Kretzschmar,⁹¹ P. Krieger,¹⁶⁶ F. Krieter,¹¹⁴ A. Krishnan,^{61b} K. Krizka,¹⁸ K. Kroeninger,⁴⁷ H. Kroha,¹¹⁵ J. Kroll,¹⁴⁰ J. Kroll,¹³⁶ K. S. Krowpman,¹⁰⁷ U. Kruchonak,⁸⁰ H. Krüger,²⁴ N. Krumnack,⁷⁹ M. C. Kruse,⁴⁹ J. A. Krzysiak,⁸⁵ O. Kuchinskaia,¹⁶⁵ S. Kudah,^{4b} D. Kuechler,⁴⁶ J. T. Kuechler,⁴⁶ S. Kuehn,³⁶ A. Kugel,^{61a} T. Kuhl,⁴⁶ V. Kukhtin,⁸⁰ Y. Kulchitsky,^{108,bb} S. Kuleshov,^{146b} Y. P. Kulinich,¹⁷² M. Kuna,⁵⁸ T. Kunigo,⁸⁶ A. Kupco,¹⁴⁰ T. Kupfer,⁴⁷ O. Kuprash,⁵² H. Kurashige,⁸³ L. L. Kurchaninov,^{167a} Y. A. Kurochkin,¹⁰⁸ A. Kurova,¹¹² M. G. Kurth,^{15a,15d} E. S. Kuwertz,³⁶ M. Kuze,¹⁶⁴ A. K. Kvam,¹⁴⁷ J. Kvita,¹³⁰ T. Kwan,¹⁰⁴ F. La Ruffa,^{41b,41a} C. Lacasta,¹⁷³ F. Lacava,^{73a,73b} D. P. J. Lack,¹⁰¹ H. Lacker,¹⁹ D. Lacour,¹³⁵ E. Ladygin,⁸⁰ R. Lafaye,⁵

B. Laforge,¹³⁵ T. Lagouri,^{146b} S. Lai,⁵³ I. K. Lakomic,^{84a} J. E. Lambert,¹²⁸ S. Lammers,⁶⁶ W. Lampl,⁷ C. Lampoudis,¹⁶¹ E. Lançon,²⁹ U. Landgraf,⁵² M. P. J. Landon,⁹³ M. C. Lanfermann,⁵⁴ V. S. Lang,⁵² J. C. Lange,⁵³ R. J. Langenberg,¹⁰³ A. J. Lankford,¹⁷⁰ F. Lanni,²⁹ K. Lantzsch,²⁴ A. Lanza,^{71a} A. Lapertosa,^{55b,55a} S. Laplace,¹³⁵ J. F. Laporte,¹⁴⁴ T. Lari,^{69a} F. Lasagni Manghi,^{23b,23a} M. Lassnig,³⁶ T. S. Lau,^{63a} A. Laudrain,⁶⁵ A. Laurier,³⁴ M. Lavorgna,^{70a,70b} S. D. Lawlor,⁹⁴ M. Lazzaroni,^{69a,69b} B. Le,¹⁰¹ E. Le Guirriec,¹⁰² A. Lebedev,⁷⁹ M. LeBlanc,⁷ T. LeCompte,⁶ F. Ledroit-Guillon,⁵⁸ A. C. A. Lee,⁹⁵ C. A. Lee,²⁹ G. R. Lee,¹⁷ L. Lee,⁵⁹ S. C. Lee,¹⁵⁷ S. Lee,⁷⁹ B. Lefebvre,^{167a} H. P. Lefebvre,⁹⁴ M. Lefebvre,¹⁷⁵ C. Leggett,¹⁸ K. Lehmann,¹⁵¹ N. Lehmann,²⁰ G. Lehmann Miotto,³⁶ W. A. Leight,⁴⁶ A. Leisos,^{161,cc} M. A. L. Leite,^{81d} C. E. Leitgeb,¹¹⁴ R. Leitner,¹⁴² D. Lellouch,^{179,a} K. J. C. Leney,⁴² T. Lenz,²⁴ S. Leone,^{72a} C. Leonidopoulos,⁵⁰ A. Leopold,¹³⁵ C. Leroy,¹¹⁰ R. Les,¹⁶⁶ C. G. Lester,³² M. Levchenko,¹³⁷ J. Levêque,⁵ D. Levin,¹⁰⁶ L. J. Levinson,¹⁷⁹ D. J. Lewis,²¹ B. Li,^{15b} B. Li,¹⁰⁶ C-Q. Li,^{60a} F. Li,^{60c} H. Li,^{60a} H. Li,^{60b} J. Li,^{60c} K. Li,¹⁴⁷ L. Li,^{60c} M. Li,^{15a,15d} Q. Li,^{15a,15d} Q. Y. Li,^{60a} S. Li,^{60d,60c} X. Li,⁴⁶ Y. Li,⁴⁶ Z. Li,^{60b} Z. Li,¹³⁴ Z. Li,¹⁰⁴ Z. Liang,^{15a} M. Liberatore,⁴⁶ B. Liberti,^{74a} A. Liblong,¹⁶⁶ K. Lie,^{63c} S. Lim,²⁹ C. Y. Lin,³² K. Lin,¹⁰⁷ R. A. Linck,⁶⁶ R. E. Lindley,⁷ J. H. Lindon,²¹ A. Lins,⁴⁶ A. L. Lioni,⁵⁴ E. Lipeles,¹³⁶ A. Lipniacka,¹⁷ T. M. Liss,^{172,dd} A. Lister,¹⁷⁴ J. D. Little,⁸ B. Liu,⁷⁹ B. L. Liu,⁶ H. B. Liu,²⁹ J. B. Liu,^{60a} J. K. K. Liu,³⁷ K. Liu,^{60d} M. Liu,^{60a} P. Liu,^{15a} Y. Liu,⁴⁶ Y. Liu,^{15a,15d} Y. L. Liu,¹⁰⁶ Y. W. Liu,^{60a} M. Livan,^{71a,71b} A. Lleres,⁵⁸ J. Llorente Merino,¹⁵¹ S. L. Lloyd,⁹³ C. Y. Lo,^{63b} E. M. Lobodzinska,⁴⁶ P. Loch,⁷ S. Loffredo,^{74a,74b} T. Lohse,¹⁹ K. Lohwasser,¹⁴⁸ M. Lokajicek,¹⁴⁰ J. D. Long,¹⁷² R. E. Long,⁹⁰ L. Longo,³⁶ K. A. Looper,¹²⁷ I. Lopez Paz,¹⁰¹ A. Lopez Solis,¹⁴⁸ J. Lorenz,¹¹⁴ N. Lorenzo Martinez,⁵ A. M. Lory,¹¹⁴ P. J. Lösel,¹¹⁴ A. Lösle,⁵² X. Lou,⁴⁶ X. Lou,^{15a} A. Lounis,⁶⁵ J. Love,⁶ P. A. Love,⁹⁰ J. J. Lozano Bahilo,¹⁷³ M. Lu,^{60a} Y. J. Lu,⁶⁴ H. J. Lubatti,¹⁴⁷ C. Luci,^{73a,73b} F. L. Lucio Alves,^{15c} A. Lucotte,⁵⁸ F. Luehring,⁶⁶ I. Luise,¹³⁵ L. Luminari,^{73a} B. Lund-Jensen,¹⁵³ M. S. Lutz,¹⁶⁰ D. Lynn,²⁹ H. Lyons,⁹¹ R. Lysak,¹⁴⁰ E. Lytken,⁹⁷ F. Lyu,^{15a} V. Lyubushkin,⁸⁰ T. Lyubushkina,⁸⁰ H. Ma,²⁹ L. L. Ma,^{60b} Y. Ma,⁹⁵ D. M. Mac Donell,¹⁷⁵ G. Maccarrone,⁵¹ A. Macchiolo,¹¹⁵ C. M. Macdonald,¹⁴⁸ J. C. MacDonald,¹⁴⁸ J. Machado Miguens,¹³⁶ D. Madaffari,¹⁷³ R. Madar,³⁸ W. F. Mader,⁴⁸ M. Madugoda Ralalage Don,¹²⁹ N. Madysa,⁴⁸ J. Maeda,⁸³ T. Maeno,²⁹ M. Maerker,⁴⁸ V. Magerl,⁵² N. Magini,⁷⁹ J. Magro,^{67a,67c,m} D. J. Mahon,³⁹ C. Maidantchik,^{81b} T. Maier,¹¹⁴ A. Maio,^{139a,139b,139d} K. Maj,^{84a} O. Majersky,^{28a} S. Majewski,¹³¹ Y. Makida,⁸² N. Makovec,⁶⁵ B. Malaescu,¹³⁵ Pa. Malecki,⁸⁵ V. P. Maleev,¹³⁷ F. Malek,⁵⁸ U. Mallik,⁷⁸ D. Malon,⁶ C. Malone,³² S. Maltezos,¹⁰ S. Malyukov,⁸⁰ J. Mamuzic,¹⁷³ G. Mancini,^{70a,70b} I. Mandić,⁹² L. Manhaes de Andrade Filho,^{81a} I. M. Maniatis,¹⁶¹ J. Manjarres Ramos,⁴⁸ K. H. Mankinen,⁹⁷ A. Mann,¹¹⁴ A. Manousos,⁷⁷ B. Mansoulié,¹⁴⁴ I. Manthos,¹⁶¹ S. Manzoni,¹²⁰ A. Marantis,¹⁶¹ G. Marceca,³⁰ L. Marchese,¹³⁴ G. Marchiori,¹³⁵ M. Marcisovsky,¹⁴⁰ L. Marcoccia,^{74a,74b} C. Marcon,⁹⁷ C. A. Marin Tobon,³⁶ M. Marjanovic,¹²⁸ Z. Marshall,¹⁸ M. U. F. Martensson,¹⁷¹ S. Marti-Garcia,¹⁷³ C. B. Martin,¹²⁷ T. A. Martin,¹⁷⁷ V. J. Martin,⁵⁰ B. Martin dit Latour,¹⁷ L. Martinelli,^{75a,75b} M. Martinez,^{14,t} P. Martinez Agullo,¹⁷³ V. I. Martinez Outschoorn,¹⁰³ S. Martin-Haugh,¹⁴³ V. S. Martoiu,^{27b} A. C. Martyniuk,⁹⁵ A. Marzin,³⁶ S. R. Maschek,¹¹⁵ L. Masetti,¹⁰⁰ T. Mashimo,¹⁶² R. Mashinistov,¹¹¹ J. Masik,¹⁰¹ A. L. Maslennikov,^{122b,122a} L. Massa,^{23b,23a} P. Massarotti,^{70a,70b} P. Mastrandrea,^{72a,72b} A. Mastroberardino,^{41b,41a} T. Masubuchi,¹⁶² D. Matakias,²⁹ A. Matic,¹¹⁴ N. Matsuzawa,¹⁶² P. Mättig,²⁴ J. Maurer,^{27b} B. Maček,⁹² D. A. Maximov,^{122b,122a} R. Mazini,¹⁵⁷ I. Maznas,¹⁶¹ S. M. Mazza,¹⁴⁵ J. P. Mc Gowan,¹⁰⁴ S. P. Mc Kee,¹⁰⁶ T. G. McCarthy,¹¹⁵ W. P. McCormack,¹⁸ E. F. McDonald,¹⁰⁵ J. A. Mcfayden,³⁶ G. Mchedlidze,^{158b} M. A. McKay,⁴² K. D. McLean,¹⁷⁵ S. J. McMahon,¹⁴³ P. C. McNamara,¹⁰⁵ C. J. McNicol,¹⁷⁷ R. A. McPherson,^{175,1} J. E. Mdhului,^{33e} Z. A. Meadows,¹⁰³ S. Meehan,³⁶ T. Megy,³⁸ S. Mehlhase,¹¹⁴ A. Mehta,⁹¹ B. Meirose,⁴³ D. Melini,¹⁵⁹ B. R. Mellado Garcia,^{33e} J. D. Mellenthin,⁵³ M. Melo,^{28a} F. Meloni,⁴⁶ A. Melzer,²⁴ E. D. Mendes Gouveia,^{139a,139e} L. Meng,³⁶ X. T. Meng,¹⁰⁶ S. Menke,¹¹⁵ E. Meoni,^{41b,41a} S. Mergelmeyer,¹⁹ S. A. M. Merkt,¹³⁸ C. Merlassino,¹³⁴ P. Mermod,⁵⁴ L. Merola,^{70a,70b} C. Meroni,^{69a} G. Merz,¹⁰⁶ O. Meshkov,^{113,111} J. K. R. Meshreki,¹⁵⁰ J. Metcalfe,⁶ A. S. Mete,⁶ C. Meyer,⁶⁶ J-P. Meyer,¹⁴⁴ M. Michetti,¹⁹ R. P. Middleton,¹⁴³ L. Mijović,⁵⁰ G. Mikenberg,¹⁷⁹ M. Mikesikova,¹⁴⁰ M. Mikuž,⁹² H. Mildner,¹⁴⁸ M. Milesi,¹⁰⁵ A. Milic,¹⁶⁶ C. D. Milke,⁴² D. W. Miller,³⁷ A. Milov,¹⁷⁹ D. A. Milstead,^{45a,45b} R. A. Mina,¹⁵² A. A. Minaenko,¹²³ I. A. Minashvili,^{158b} A. I. Mincer,¹²⁵ B. Mindur,^{84a} M. Mineev,⁸⁰ Y. Minegishi,¹⁶² L. M. Mir,¹⁴ M. Mironova,¹³⁴ A. Mirtó,^{68a,68b} K. P. Mistry,¹³⁶ T. Mitani,¹⁷⁸ J. Mitrevski,¹¹⁴ V. A. Mitsou,¹⁷³ M. Mittal,^{60c} O. Miu,¹⁶⁶ A. Miucci,²⁰ P. S. Miyagawa,⁹³ A. Mizukami,⁸² J. U. Mjörnmark,⁹⁷ T. Mkrtchyan,^{61a} M. Mlynarikova,¹⁴² T. Moa,^{45a,45b} S. Mobius,⁵³ K. Mochizuki,¹¹⁰ P. Mogg,¹¹⁴ S. Mohapatra,³⁹ R. Moles-Valls,²⁴ K. Mönig,⁴⁶ E. Monnier,¹⁰² A. Montalbano,¹⁵¹ J. Montejo Berlingen,³⁶ M. Montella,⁹⁵ F. Monticelli,⁸⁹ S. Monzani,^{69a} N. Morange,⁶⁵ D. Moreno,^{22a} M. Moreno Llácer,¹⁷³ C. Moreno Martinez,¹⁴ P. Morettini,^{55b} M. Morgenstern,¹⁵⁹ S. Morgenstern,⁴⁸ D. Mori,¹⁵¹ M. Morii,⁵⁹ M. Morinaga,¹⁷⁸ V. Morisbak,¹³³ A. K. Morley,³⁶ G. Mornacchi,³⁶ A. P. Morris,⁹⁵ L. Morvaj,¹⁵⁴ P. Moschovakos,³⁶ B. Moser,¹²⁰

M. Mosidze,^{158b} T. Moskalets,¹⁴⁴ H. J. Moss,¹⁴⁸ J. Moss,^{31,ee} E. J. W. Moyses,¹⁰³ S. Muanza,¹⁰² J. Mueller,¹³⁸
 R. S. P. Mueller,¹¹⁴ D. Muenstermann,⁹⁰ G. A. Mullier,⁹⁷ D. P. Mungo,^{69a,69b} J. L. Munoz Martinez,¹⁴
 F. J. Munoz Sanchez,¹⁰¹ P. Murin,^{28b} W. J. Murray,^{177,143} A. Murrone,^{69a,69b} J. M. Muse,¹²⁸ M. Muškinja,¹⁸ C. Mwewa,^{33a}
 A. G. Myagkov,^{123,i} A. A. Myers,¹³⁸ J. Myers,¹³¹ M. Myska,¹⁴¹ B. P. Nachman,¹⁸ O. Nackenhorst,⁴⁷ A. Nag Nag,⁴⁸
 K. Nagai,¹³⁴ K. Nagano,⁸² Y. Nagasaka,⁶² J. L. Nagle,²⁹ E. Nagy,¹⁰² A. M. Nairz,³⁶ Y. Nakahama,¹¹⁷ K. Nakamura,⁸²
 T. Nakamura,¹⁶² H. Nanjo,¹³² F. Napolitano,^{61a} R. F. Naranjo Garcia,⁴⁶ R. Narayan,⁴² I. Naryshkin,¹³⁷ T. Naumann,⁴⁶
 G. Navarro,^{22a} P. Y. Nechaeva,¹¹¹ F. Nechansky,⁴⁶ T. J. Neep,²¹ A. Negri,^{71a,71b} M. Negrini,^{23b} C. Nellist,¹¹⁹ C. Nelson,¹⁰⁴
 M. E. Nelson,^{45a,45b} S. Nemecek,¹⁴⁰ M. Nessi,^{36,ff} M. S. Neubauer,¹⁷² F. Neuhaus,¹⁰⁰ M. Neumann,¹⁸¹ R. Newhouse,¹⁷⁴
 P. R. Newman,²¹ C. W. Ng,¹³⁸ Y. S. Ng,¹⁹ Y. W. Y. Ng,¹⁷⁰ B. Ngair,^{35e} H. D. N. Nguyen,¹⁰² T. Nguyen Manh,¹¹⁰
 E. Nibigira,³⁸ R. B. Nickerson,¹³⁴ R. Nicolaidou,¹⁴⁴ D. S. Nielsen,⁴⁰ J. Nielsen,¹⁴⁵ M. Niemeyer,⁵³ N. Nikiforou,¹¹
 V. Nikolaenko,^{123,i} I. Nikolic-Audit,¹³⁵ K. Nikolopoulos,²¹ P. Nilsson,²⁹ H. R. Nindhito,⁵⁴ Y. Ninomiya,⁸² A. Nisati,^{73a}
 N. Nishu,^{60c} R. Nisius,¹¹⁵ I. Nitsche,⁴⁷ T. Nitta,¹⁷⁸ T. Nobe,¹⁶² D. L. Noel,³² Y. Noguchi,⁸⁶ I. Nomidis,¹³⁵ M. A. Nomura,²⁹
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 F. G. Oakham,^{34,d} H. Oberlack,¹¹⁵ J. Ocariz,¹³⁵ A. Ochi,⁸³ I. Ochoa,³⁹ J. P. Ochoa-Ricoux,^{146a} K. O'Connor,²⁶ S. Oda,⁸⁸
 S. Odaka,⁸² S. Oerdek,⁵³ A. Ogrodnik,^{84a} A. Oh,¹⁰¹ S. H. Oh,⁴⁹ C. C. Ohm,¹⁵³ H. Oide,¹⁶⁴ M. L. Ojeda,¹⁶⁶ H. Okawa,¹⁶⁸
 Y. Okazaki,⁸⁶ M. W. O'Keefe,⁹¹ Y. Okumura,¹⁶² T. Okuyama,⁸² A. Olariu,^{27b} L. F. Oleiro Seabra,^{139a} S. A. Olivares Pino,^{146a}
 D. Oliveira Damazio,²⁹ J. L. Oliver,¹ M. J. R. Olsson,¹⁷⁰ A. Olszewski,⁸⁵ J. Olszowska,⁸⁵ D. C. O'Neil,¹⁵¹ A. P. O'Neill,¹³⁴
 A. Onofre,^{139a,139e} P. U. E. Onyisi,¹¹ H. Oppen,¹³³ R. G. Oreamuno Madriz,¹²¹ M. J. Oreglia,³⁷ G. E. Orellana,⁸⁹
 D. Orestano,^{75a,75b} N. Orlando,¹⁴ R. S. Orr,¹⁶⁶ V. O'Shea,⁵⁷ R. Ospanov,^{60a} G. Otero y Garzon,³⁰ H. Otono,⁸⁸ P. S. Ott,^{61a}
 G. J. Ottino,¹⁸ M. Ouchrif,^{35d} J. Ouellette,²⁹ F. Ould-Saada,¹³³ A. Ouraou,¹⁴⁴ Q. Ouyang,^{15a} M. Owen,⁵⁷ R. E. Owen,²¹
 V. E. Ozcan,^{12c} N. Ozturk,⁸ J. Pacalt,¹³⁰ H. A. Pacey,³² K. Pachal,⁴⁹ A. Pacheco Pages,¹⁴ C. Padilla Aranda,¹⁴
 S. Pagan Griso,¹⁸ G. Palacino,⁶⁶ S. Palazzo,⁵⁰ S. Palestini,³⁶ M. Palka,^{84b} D. Pallin,³⁸ P. Palni,^{84a} C. E. Pandini,⁵⁴
 J. G. Panduro Vazquez,⁹⁴ P. Pani,⁴⁶ G. Panizzo,^{67a,67c} L. Paolozzi,⁵⁴ C. Papadatos,¹¹⁰ K. Papageorgiou,^{9,q} S. Parajuli,⁴²
 A. Paramonov,⁶ C. Paraskevopoulos,¹⁰ D. Paredes Hernandez,^{63b} S. R. Paredes Saenz,¹³⁴ B. Parida,¹⁷⁹ T. H. Park,¹⁶⁶
 A. J. Parker,³¹ M. A. Parker,³² F. Parodi,^{55b,55a} E. W. Parrish,¹²¹ J. A. Parsons,³⁹ U. Parzefall,⁵² L. Pascual Dominguez,¹³⁵
 V. R. Pascuzzi,¹⁸ J. M. P. Pasner,¹⁴⁵ F. Pasquali,¹²⁰ E. Pasqualucci,^{73a} S. Passaggio,^{55b} F. Pastore,⁹⁴ P. Pasuwan,^{45a,45b}
 S. Pataria,¹⁰⁰ J. R. Pater,¹⁰¹ A. Pathak,^{180,e} J. Patton,⁹¹ T. Pauly,³⁶ J. Pearkes,¹⁵² B. Pearson,¹¹⁵ M. Pedersen,¹³³
 L. Pedraza Diaz,¹¹⁹ R. Pedro,^{139a} T. Peiffer,⁵³ S. V. Peleganchuk,^{122b,122a} O. Penc,¹⁴⁰ H. Peng,^{60a} B. S. Peralva,^{81a}
 M. M. Perego,⁶⁵ A. P. Pereira Peixoto,^{139a} L. Pereira Sanchez,^{45a,45b} D. V. Perepelitsa,²⁹ E. Perez Codina,^{167a} F. Peri,¹⁹
 L. Perini,^{69a,69b} H. Pernegger,³⁶ S. Perrella,^{139a} A. Perrevoort,¹²⁰ K. Peters,⁴⁶ R. F. Y. Peters,¹⁰¹ B. A. Petersen,³⁶
 T. C. Petersen,⁴⁰ E. Petit,¹⁰² A. Petridis,¹ C. Petridou,¹⁶¹ F. Petrucci,^{75a,75b} M. Pettee,¹⁸² N. E. Pettersson,¹⁰³ K. Petukhova,¹⁴²
 A. Peyaud,¹⁴⁴ R. Pezoa,^{146d} L. Pezzotti,^{71a,71b} T. Pham,¹⁰⁵ F. H. Phillips,¹⁰⁷ P. W. Phillips,¹⁴³ M. W. Phipps,¹⁷²
 G. Piacquadio,¹⁵⁴ E. Pianori,¹⁸ A. Picazio,¹⁰³ R. H. Pickles,¹⁰¹ R. Piegaia,³⁰ D. Pietreanu,^{27b} J. E. Pilcher,³⁷
 A. D. Pilkington,¹⁰¹ M. Pinamonti,^{67a,67c} J. L. Pinfold,³ C. Pitman Donaldson,⁹⁵ M. Pitt,¹⁶⁰ L. Pizzimento,^{74a,74b}
 M.-A. Pleier,²⁹ V. Pleskot,¹⁴² E. Plotnikova,⁸⁰ P. Podberezko,^{122b,122a} R. Poettgen,⁹⁷ R. Poggi,⁵⁴ L. Poggioli,¹³⁵
 I. Pogrebnyak,¹⁰⁷ D. Pohl,²⁴ I. Pokharel,⁵³ G. Polesello,^{71a} A. Poley,¹⁵¹ A. Policicchio,^{73a,73b} R. Polifka,¹⁴² A. Polini,^{23b}
 C. S. Pollard,⁴⁶ V. Polychronakos,²⁹ D. Ponomarenko,¹¹² L. Pontecorvo,³⁶ S. Popa,^{27a} G. A. Popeneciu,^{27d} L. Portales,⁵
 D. M. Portillo Quintero,⁵⁸ S. Pospisil,¹⁴¹ K. Potamianos,⁴⁶ I. N. Potrap,⁸⁰ C. J. Potter,³² H. Potti,¹¹ T. Poulsen,⁹⁷ J. Poveda,¹⁷³
 T. D. Powell,¹⁴⁸ G. Pownall,⁴⁶ M. E. Pozo Astigarraga,³⁶ P. Pralavorio,¹⁰² S. Prell,⁷⁹ D. Price,¹⁰¹ M. Primavera,^{68a}
 M. L. Proffitt,¹⁴⁷ N. Proklova,¹¹² K. Prokofiev,^{63c} F. Prokoshin,⁸⁰ S. Protopopescu,²⁹ J. Proudfoot,⁶ M. Przybycien,^{84a}
 D. Pudza,¹³⁷ A. Puri,¹⁷² P. Puzo,⁶⁵ D. Pyatiizbyantseva,¹¹² J. Qian,¹⁰⁶ Y. Qin,¹⁰¹ A. Quadt,⁵³ M. Queitsch-Maitland,³⁶
 A. Qureshi,¹ M. Racko,^{28a} F. Ragusa,^{69a,69b} G. Rahal,⁹⁸ J. A. Raine,⁵⁴ S. Rajagopalan,²⁹ A. Ramirez Morales,⁹³ K. Ran,^{15a,15d}
 D. M. Rauch,⁴⁶ F. Rauscher,¹¹⁴ S. Rave,¹⁰⁰ B. Ravina,¹⁴⁸ I. Ravinovitch,¹⁷⁹ J. H. Rawling,¹⁰¹ M. Raymond,³⁶ A. L. Read,¹³³
 N. P. Radioff,⁵⁸ M. Reale,^{68a,68b} D. M. Rebuffi,^{71a,71b} G. Redlinger,²⁹ K. Reeves,⁴³ J. Reichert,¹³⁶ D. Reikher,¹⁶⁰ A. Reiss,¹⁰⁰
 A. Rej,¹⁵⁰ C. Rembser,³⁶ A. Renardi,⁴⁶ M. Renda,^{27b} M. B. Rendel,¹¹⁵ S. Resconi,^{69a} E. D. Resseguie,¹⁸ S. Rettie,⁹⁵
 B. Reynolds,¹²⁷ E. Reynolds,²¹ O. L. Rezanova,^{122b,122a} P. Reznicek,¹⁴² E. Ricci,^{76a,76b} R. Richter,¹¹⁵ S. Richter,⁴⁶
 E. Richter-Was,^{84b} M. Ridel,¹³⁵ P. Rieck,¹¹⁵ O. Rifki,⁴⁶ M. Rijssenbeek,¹⁵⁴ A. Rimoldi,^{71a,71b} M. Rimoldi,⁴⁶ L. Rinaldi,^{23b}
 T. T. Rinn,¹⁷² G. Ripellino,¹⁵³ I. Riu,¹⁴ P. Rivadeneira,⁴⁶ J. C. Rivera Vergara,¹⁷⁵ F. Rizatdinova,¹²⁹ E. Rizvi,⁹³ C. Rizzi,³⁶

S. H. Robertson,^{104,1} M. Robin,⁴⁶ D. Robinson,³² C. M. Robles Gajardo,^{146d} M. Robles Manzano,¹⁰⁰ A. Robson,⁵⁷
A. Rocchi,^{74a,74b} E. Rocco,¹⁰⁰ C. Roda,^{72a,72b} S. Rodriguez Bosca,¹⁷³ A. M. Rodríguez Vera,^{167b} S. Roe,³⁶ J. Roggel,¹⁸¹
O. Røhne,¹³³ R. Röhrig,¹¹⁵ R. A. Rojas,^{146d} B. Roland,⁵² C. P. A. Roland,⁶⁶ J. Roloff,²⁹ A. Romaniouk,¹¹² M. Romano,^{23b,23a}
N. Rompotis,⁹¹ M. Ronzani,¹²⁵ L. Roos,¹³⁵ S. Rosati,^{73a} G. Rosin,¹⁰³ B. J. Rosser,¹³⁶ E. Rossi,⁴⁶ E. Rossi,^{75a,75b}
E. Rossi,^{70a,70b} L. P. Rossi,^{55b} L. Rossini,^{69a,69b} R. Rosten,¹⁴ M. Rotaru,^{27b} B. Rottler,⁵² D. Rousseau,⁶⁵ G. Rovelli,^{71a,71b}
A. Roy,¹¹ D. Roy,^{33e} A. Rozanov,¹⁰² Y. Rozen,¹⁵⁹ X. Ruan,^{33e} F. Rühr,⁵² A. Ruiz-Martinez,¹⁷³ A. Rummler,³⁶ Z. Rurikova,⁵²
N. A. Rusakovich,⁸⁰ H. L. Russell,¹⁰⁴ L. Rustige,^{38,47} J. P. Rutherford,⁷ E. M. Rüttinger,¹⁴⁸ M. Rybar,³⁹ G. Rybkin,⁶⁵
E. B. Rye,¹³³ A. Ryzhov,¹²³ J. A. Sabater Iglesias,⁴⁶ P. Sabatini,⁵³ S. Sacerdoti,⁶⁵ H. F.-W. Sadrozinski,¹⁴⁵ R. Sadykov,⁸⁰
F. Safai Tehrani,^{73a} B. Safarzadeh Samani,¹⁵⁵ M. Safdari,¹⁵² P. Saha,¹²¹ S. Saha,¹⁰⁴ M. Sahinsoy,¹¹⁵ A. Sahu,¹⁸¹
M. Saimpert,³⁶ M. Saito,¹⁶² T. Saito,¹⁶² H. Sakamoto,¹⁶² D. Salamani,⁵⁴ G. Salamanna,^{75a,75b} A. Salmnikov,¹⁵² J. Salt,¹⁷³
A. Salvador Salas,¹⁴ D. Salvatore,^{41b,41a} F. Salvatore,¹⁵⁵ A. Salvucci,^{63a,63b,63c} A. Salzburger,³⁶ J. Samarati,³⁶ D. Sammel,⁵²
D. Sampsonidis,¹⁶¹ D. Sampsonidou,¹⁶¹ J. Sánchez,¹⁷³ A. Sanchez Pineda,^{67a,36,67c} H. Sandaker,¹³³ C. O. Sander,⁴⁶
I. G. Sanderswood,⁹⁰ M. Sandhoff,¹⁸¹ C. Sandoval,^{22a} D. P. C. Sankey,¹⁴³ M. Sannino,^{55b,55a} Y. Sano,¹¹⁷ A. Sansoni,⁵¹
C. Santoni,³⁸ H. Santos,^{139a,139b} S. N. Santpur,¹⁸ A. Santra,¹⁷³ A. Sapronov,⁸⁰ J. G. Saraiva,^{139a,139d} O. Sasaki,⁸² K. Sato,¹⁶⁸
F. Sauerburger,⁵² E. Sauvan,⁵ P. Savard,^{166,d} R. Sawada,¹⁶² C. Sawyer,¹⁴³ L. Sawyer,^{96,gg} I. Sayago Galvan,¹⁷³ C. Sbarra,^{23b}
A. Sbrizzi,^{67a,67c} T. Scanlon,⁹⁵ J. Schaarschmidt,¹⁴⁷ P. Schacht,¹¹⁵ D. Schaefer,³⁷ L. Schaefer,¹³⁶ S. Schaepe,³⁶ U. Schäfer,¹⁰⁰
A. C. Schaffer,⁶⁵ D. Schaile,¹¹⁴ R. D. Schamberger,¹⁵⁴ E. Schanet,¹¹⁴ N. Scharmberg,¹⁰¹ V. A. Schegelsky,¹³⁷
D. Scheirich,¹⁴² F. Schenck,¹⁹ M. Schernau,¹⁷⁰ C. Schiavi,^{55b,55a} L. K. Schildgen,²⁴ Z. M. Schillaci,²⁶ E. J. Schioppa,^{68a,68b}
M. Schioppa,^{41b,41a} K. E. Schleicher,⁵² S. Schlenker,³⁶ K. R. Schmidt-Sommerfeld,¹¹⁵ K. Schmieden,³⁶ C. Schmitt,¹⁰⁰
S. Schmitt,⁴⁶ J. C. Schmoedel,⁴⁶ L. Schoeffel,¹⁴⁴ A. Schoening,^{61b} P. G. Scholer,⁵² E. Schopf,¹³⁴ M. Schott,¹⁰⁰
J. F. P. Schouwenberg,¹¹⁹ J. Schovancova,³⁶ S. Schramm,⁵⁴ F. Schroeder,¹⁸¹ A. Schulte,¹⁰⁰ H.-C. Schultz-Coulon,^{61a}
M. Schumacher,⁵² B. A. Schumm,¹⁴⁵ Ph. Schune,¹⁴⁴ A. Schwartzman,¹⁵² T. A. Schwarz,¹⁰⁶ Ph. Schwemling,¹⁴⁴
R. Schwienhorst,¹⁰⁷ A. Sciandra,¹⁴⁵ G. Sciolla,²⁶ M. Scodreggio,⁴⁶ M. Scornajenghi,^{41b,41a} F. Scuri,^{72a} F. Scutti,¹⁰⁵
L. M. Scyboz,¹¹⁵ C. D. Sebastiani,⁹¹ P. Seema,¹⁹ S. C. Seidel,¹¹⁸ A. Seiden,¹⁴⁵ B. D. Seidlitz,²⁹ T. Seiss,³⁷ C. Seitz,⁴⁶
J. M. Seixas,^{81b} G. Sekhniaidze,^{70a} S. J. Sekula,⁴² N. Semprini-Cesari,^{23b,23a} S. Sen,⁴⁹ C. Serfon,²⁹ L. Serin,⁶⁵ L. Serkin,^{67a,67b}
M. Sessa,^{60a} H. Severini,¹²⁸ S. Sevova,¹⁵² F. Sforza,^{55b,55a} A. Sfyrta,⁵⁴ E. Shabalina,⁵³ J. D. Shahinian,¹⁴⁵ N. W. Shaikh,^{45a,45b}
D. Shaked Renous,¹⁷⁹ L. Y. Shan,^{15a} M. Shapiro,¹⁸ A. Sharma,¹³⁴ A. S. Sharma,¹ P. B. Shatalov,¹²⁴ K. Shaw,¹⁵⁵
S. M. Shaw,¹⁰¹ M. Shehade,¹⁷⁹ Y. Shen,¹²⁸ A. D. Sherman,²⁵ P. Sherwood,⁹⁵ L. Shi,⁹⁵ S. Shimizu,⁸² C. O. Shimmin,¹⁸²
Y. Shimogama,¹⁷⁸ M. Shimojima,¹¹⁶ I. P. J. Shipsey,¹³⁴ S. Shirabe,¹⁶⁴ M. Shiyakova,^{80,hh} J. Shlomi,¹⁷⁹ A. Shmeleva,¹¹¹
M. J. Shochet,³⁷ J. Shojaii,¹⁰⁵ D. R. Shope,¹²⁸ S. Shrestha,¹²⁷ E. M. Shrif,^{33e} E. Shulga,¹⁷⁹ P. Sicho,¹⁴⁰ A. M. Sickles,¹⁷²
E. Sideras Haddad,^{33e} O. Sidiropoulou,³⁶ A. Sidoti,^{23b,23a} F. Siegert,⁴⁸ Dj. Sijacki,¹⁶ M. Silva Jr.,¹⁸⁰ M. V. Silva Oliveira,³⁶
S. B. Silverstein,^{45a} S. Simion,⁶⁵ R. Simoniello,¹⁰⁰ C. J. Simpson-allsoy,²¹ S. Simsek,^{12b} P. Sinervo,¹⁶⁶ V. Sinetckii,¹¹³
S. Singh,¹⁵¹ M. Sioli,^{23b,23a} I. Siral,¹³¹ S. Yu. Sivoklov,¹¹³ J. Sjölin,^{45a,45b} A. Skaf,⁵³ E. Skorda,⁹⁷ P. Skubic,¹²⁸
M. Slawinska,⁸⁵ K. Sliwa,¹⁶⁹ R. Slovak,¹⁴² V. Smakhtin,¹⁷⁹ B. H. Smart,¹⁴³ J. Smiesko,^{28b} N. Smirnov,¹¹² S. Yu. Smirnov,¹¹²
Y. Smirnov,¹¹² L. N. Smirnova,^{113,ii} O. Smirnova,⁹⁷ H. A. Smith,¹³⁴ M. Smizanska,⁹⁰ K. Smolek,¹⁴¹ A. Smykiewicz,⁸⁵
A. A. Snesarev,¹¹¹ H. L. Snoek,¹²⁰ I. M. Snyder,¹³¹ S. Snyder,²⁹ R. Sobie,^{175,1} A. Soffer,¹⁶⁰ A. Søggaard,⁵⁰ F. Sohns,⁵³
C. A. Solans Sanchez,³⁶ E. Yu. Soldatov,¹¹² U. Soldevila,¹⁷³ A. A. Solodkov,¹²³ A. Soloshenko,⁸⁰ O. V. Solovyanov,¹²³
V. Solovyev,¹³⁷ P. Sommer,¹⁴⁸ H. Son,¹⁶⁹ W. Song,¹⁴³ W. Y. Song,^{167b} A. Sopczak,¹⁴¹ A. L. Soppio,⁹⁵ F. Sopkova,^{28b}
S. Sottocornola,^{71a,71b} R. Soualah,^{67a,67c} A. M. Soukharev,^{122b,122a} D. South,⁴⁶ S. Spagnolo,^{68a,68b} M. Spalla,¹¹⁵
M. Spangenberg,¹⁷⁷ F. Spanò,⁹⁴ D. Sperlich,⁵² T. M. Spieker,^{61a} G. Spigo,³⁶ M. Spina,¹⁵⁵ D. P. Spiteri,⁵⁷ M. Spousta,¹⁴²
A. Stabile,^{69a,69b} B. L. Stamas,¹²¹ R. Stamen,^{61a} M. Stamenkovic,¹²⁰ E. Stanecka,⁸⁵ B. Stanislaus,¹³⁴ M. M. Stanitzki,⁴⁶
M. Stankaityte,¹³⁴ B. Stapf,¹²⁰ E. A. Starchenko,¹²³ G. H. Stark,¹⁴⁵ J. Stark,⁵⁸ P. Staroba,¹⁴⁰ P. Starovoitov,^{61a} S. Stärz,¹⁰⁴
R. Staszewski,⁸⁵ G. Stavropoulos,⁴⁴ M. Stegler,⁴⁶ P. Steinberg,²⁹ A. L. Steinhebel,¹³¹ B. Stelzer,¹⁵¹ H. J. Stelzer,¹³⁸
O. Stelzer-Chilton,^{167a} H. Stenzel,⁵⁶ T. J. Stevenson,¹⁵⁵ G. A. Stewart,³⁶ M. C. Stockton,³⁶ G. Stoica,^{27b} M. Stolarski,^{139a}
S. Stonjek,¹¹⁵ A. Straessner,⁴⁸ J. Strandberg,¹⁵³ S. Strandberg,^{45a,45b} M. Strauss,¹²⁸ T. Strebler,¹⁰² P. Strizenec,^{28b}
R. Ströhmer,¹⁷⁶ D. M. Strom,¹³¹ R. Stroynowski,⁴² A. Strubig,⁵⁰ S. A. Stucci,²⁹ B. Stugu,¹⁷ J. Stupak,¹²⁸ N. A. Styles,⁴⁶
D. Su,¹⁵² W. Su,^{60c,147} S. Suchek,^{61a} V. V. Sulin,¹¹¹ M. J. Sullivan,⁹¹ D. M. S. Sultan,⁵⁴ S. Sultansoy,^{4c} T. Sumida,⁸⁶ S. Sun,¹⁰⁶
X. Sun,¹⁰¹ K. Suruliz,¹⁵⁵ C. J. E. Suster,¹⁵⁶ M. R. Sutton,¹⁵⁵ S. Suzuki,⁸² M. Svatos,¹⁴⁰ M. Swiatlowski,^{167a} S. P. Swift,²

T. Swirski,¹⁷⁶ A. Sydorenko,¹⁰⁰ I. Sykora,^{28a} M. Sykora,¹⁴² T. Sykora,¹⁴² D. Ta,¹⁰⁰ K. Tackmann,^{46,jj} J. Taenzer,¹⁶⁰ A. Taffard,¹⁷⁰ R. Tafirout,^{167a} H. Takai,²⁹ R. Takashima,⁸⁷ K. Takeda,⁸³ T. Takeshita,¹⁴⁹ E. P. Takeva,⁵⁰ Y. Takubo,⁸² M. Talby,¹⁰² A. A. Talyshev,^{122b,122a} K. C. Tam,^{63b} N. M. Tamir,¹⁶⁰ J. Tanaka,¹⁶² R. Tanaka,⁶⁵ S. Tapia Araya,¹⁷² S. Tapprogge,¹⁰⁰ A. Tarek Abouelfadl Mohamed,¹⁰⁷ S. Tarem,¹⁵⁹ K. Tariq,^{60b} G. Tarna,^{27b,kk} G. F. Tartarelli,^{69a} P. Tas,¹⁴² M. Tasevsky,¹⁴⁰ T. Tashiro,⁸⁶ E. Tassi,^{41b,41a} A. Tavares Delgado,^{139a} Y. Tayalati,^{35e} A. J. Taylor,⁵⁰ G. N. Taylor,¹⁰⁵ W. Taylor,^{167b} H. Teagle,⁹¹ A. S. Tee,⁹⁰ R. Teixeira De Lima,¹⁵² P. Teixeira-Dias,⁹⁴ H. Ten Kate,³⁶ J. J. Teoh,¹²⁰ S. Terada,⁸² K. Terashi,¹⁶² J. Terron,⁹⁹ S. Terzo,¹⁴ M. Testa,⁵¹ R. J. Teuscher,^{166,1} S. J. Thais,¹⁸² N. Themistokleous,⁵⁰ T. Theveneaux-Pelzer,⁴⁶ F. Thiele,⁴⁰ D. W. Thomas,⁹⁴ J. O. Thomas,⁴² J. P. Thomas,²¹ E. A. Thompson,⁴⁶ P. D. Thompson,²¹ E. Thomson,¹³⁶ E. J. Thorpe,⁹³ R. E. Ticse Torres,⁵³ V. O. Tikhomirov,^{111,ll} Yu. A. Tikhonov,^{122b,122a} S. Timoshenko,¹¹² P. Tipton,¹⁸² S. Tisserant,¹⁰² K. Todome,^{23b,23a} S. Todorova-Nova,¹⁴² S. Todt,⁴⁸ J. Tojo,⁸⁸ S. Tokár,^{28a} K. Tokushuku,⁸² E. Tolley,¹²⁷ R. Tombs,³² K. G. Tomiwa,^{33e} M. Tomoto,¹¹⁷ L. Tompkins,¹⁵² P. Tornambe,¹⁰³ E. Torrence,¹³¹ H. Torres,⁴⁸ E. Torró Pastor,¹⁴⁷ C. Tosciri,¹³⁴ J. Toth,^{102,mm} D. R. Tovey,¹⁴⁸ A. Traet,¹⁷ C. J. Treado,¹²⁵ T. Trefzger,¹⁷⁶ F. Tresoldi,¹⁵⁵ A. Tricoli,²⁹ I. M. Trigger,^{167a} S. Trincaz-Duvoid,¹³⁵ D. A. Trischuk,¹⁷⁴ W. Trischuk,¹⁶⁶ B. Trocmé,⁵⁸ A. Trofymov,⁶⁵ C. Troncon,^{69a} F. Trovato,¹⁵⁵ L. Truong,^{33c} M. Trzebinski,⁸⁵ A. Trzupek,⁸⁵ F. Tsai,⁴⁶ J. C-L. Tseng,¹³⁴ P. V. Tsiareshka,^{108,bb} A. Tsirigotis,^{161,cc} V. Tsiskaridze,¹⁵⁴ E. G. Tskhadadze,^{158a} M. Tsopoulou,¹⁶¹ I. I. Tsukerman,¹²⁴ V. Tsulaia,¹⁸ S. Tsuno,⁸² D. Tsybychev,¹⁵⁴ Y. Tu,^{63b} A. Tudorache,^{27b} V. Tudorache,^{27b} T. T. Tulbure,^{27a} A. N. Tuna,⁵⁹ S. Turchikhin,⁸⁰ D. Turgeman,¹⁷⁹ I. Turk Cakir,^{4b,nn} R. J. Turner,²¹ R. Turra,^{69a} P. M. Tuts,³⁹ S. Tzamarias,¹⁶¹ E. Tzovara,¹⁰⁰ K. Uchida,¹⁶² F. Ukegawa,¹⁶⁸ G. Unal,³⁶ M. Unal,¹¹ A. Undrus,²⁹ G. Unel,¹⁷⁰ F. C. Ungaro,¹⁰⁵ Y. Unno,⁸² K. Uno,¹⁶² J. Urban,^{28b} P. Urquijo,¹⁰⁵ G. Usai,⁸ Z. Uysal,^{12d} V. Vacek,¹⁴¹ B. Vachon,¹⁰⁴ K. O. H. Vadla,¹³³ T. Vafeiadis,³⁶ A. Vaidya,⁹⁵ C. Valderanis,¹¹⁴ E. Valdes Santurio,^{45a,45b} M. Valente,⁵⁴ S. Valentinetti,^{23b,23a} A. Valero,¹⁷³ L. Valéry,⁴⁶ R. A. Vallance,²¹ A. Vallier,³⁶ J. A. Valls Ferrer,¹⁷³ T. R. Van Daalen,¹⁴ P. Van Gemmeren,⁶ I. Van Vulpen,¹²⁰ M. Vanadia,^{74a,74b} W. Vandelli,³⁶ M. Vandenbroucke,¹⁴⁴ E. R. Vandewall,¹²⁹ A. Vaniachine,¹⁶⁵ D. Vannicola,^{73a,73b} R. Vari,^{73a} E. W. Varnes,⁷ C. Varni,^{55b,55a} T. Varol,¹⁵⁷ D. Varouchas,⁶⁵ K. E. Varvell,¹⁵⁶ M. E. Vasile,^{27b} G. A. Vasquez,¹⁷⁵ F. Vazeille,³⁸ D. Vazquez Furelos,¹⁴ T. Vazquez Schroeder,³⁶ J. Veatch,⁵³ V. Vecchio,¹⁰¹ M. J. Veen,¹²⁰ L. M. Veloce,¹⁶⁶ F. Veloso,^{139a,139c} S. Veneziano,^{73a} A. Ventura,^{68a,68b} N. Venturi,³⁶ A. Verbytskyi,¹¹⁵ V. Vercesi,^{71a} M. Verducci,^{72a,72b} C. M. Vergel Infante,⁷⁹ C. Vergis,²⁴ W. Verkerke,¹²⁰ A. T. Vermeulen,¹²⁰ J. C. Vermeulen,¹²⁰ C. Vernieri,¹⁵² M. C. Vetterli,^{151,d} N. Viaux Maira,^{146d} T. Vickey,¹⁴⁸ O. E. Vickey Boeriu,¹⁴⁸ G. H. A. Viehhauser,¹³⁴ L. Vigani,^{61b} M. Villa,^{23b,23a} M. Villaplana Perez,³ E. M. Villhauer,⁵⁰ E. Vilucchi,⁵¹ M. G. Vincter,³⁴ G. S. Virdee,²¹ A. Vishwakarma,⁵⁰ C. Vittori,^{23b,23a} I. Vivarelli,¹⁵⁵ M. Vogel,¹⁸¹ P. Vokac,¹⁴¹ S. E. von Buddenbrock,^{33e} E. Von Toerne,²⁴ V. Vorobel,¹⁴² K. Vorobev,¹¹² M. Vos,¹⁷³ J. H. Vosseveld,⁹¹ M. Vozak,¹⁰¹ N. Vranjes,¹⁶ M. Vranjes Milosavljevic,¹⁶ V. Vrba,¹⁴¹ M. Vreeswijk,¹²⁰ R. Vuillermet,³⁶ I. Vukotic,³⁷ S. Wada,¹⁶⁸ P. Wagner,²⁴ W. Wagner,¹⁸¹ J. Wagner-Kuhr,¹¹⁴ S. Wahdan,¹⁸¹ H. Wahlberg,⁸⁹ R. Wakasa,¹⁶⁸ V. M. Walbrecht,¹¹⁵ J. Walder,⁹⁰ R. Walker,¹¹⁴ S. D. Walker,⁹⁴ W. Walkowiak,¹⁵⁰ V. Wallangen,^{45a,45b} A. M. Wang,⁵⁹ A. Z. Wang,¹⁸⁰ C. Wang,^{60c} F. Wang,¹⁸⁰ H. Wang,¹⁸ H. Wang,³ J. Wang,^{63a} P. Wang,⁴² Q. Wang,¹²⁸ R.-J. Wang,¹⁰⁰ R. Wang,^{60a} R. Wang,⁶ S. M. Wang,¹⁵⁷ W. T. Wang,^{60a} W. Wang,^{15c} W. X. Wang,^{60a} Y. Wang,^{60a} Z. Wang,¹⁰⁶ C. Wanotayaroj,⁴⁶ A. Warburton,¹⁰⁴ C. P. Ward,³² D. R. Wardrope,⁹⁵ N. Warrack,⁵⁷ A. T. Watson,²¹ M. F. Watson,²¹ G. Watts,¹⁴⁷ B. M. Waugh,⁹⁵ A. F. Webb,¹¹ C. Weber,²⁹ M. S. Weber,²⁰ S. A. Weber,³⁴ S. M. Weber,^{61a} A. R. Weidberg,¹³⁴ J. Weingarten,⁴⁷ M. Weirich,¹⁰⁰ C. Weiser,⁵² P. S. Wells,³⁶ T. Wenaus,²⁹ T. Wengler,³⁶ S. Wenig,³⁶ N. Wermes,²⁴ M. Wessels,^{61a} T. D. Weston,²⁰ K. Whalen,¹³¹ N. L. Whallon,¹⁴⁷ A. M. Wharton,⁹⁰ A. S. White,¹⁰⁶ A. White,⁸ M. J. White,¹ D. Whiteson,¹⁷⁰ B. W. Whitmore,⁹⁰ W. Wiedenmann,¹⁸⁰ C. Wiel,⁴⁸ M. Wielers,¹⁴³ N. Wieseotte,¹⁰⁰ C. Wiglesworth,⁴⁰ L. A. M. Wiik-Fuchs,⁵² H. G. Wilkens,³⁶ L. J. Wilkins,⁹⁴ H. H. Williams,¹³⁶ S. Williams,³² S. Willocq,¹⁰³ P. J. Windischhofer,¹³⁴ I. Wingerter-Seez,⁵ E. Winkels,¹⁵⁵ F. Winklmeier,¹³¹ B. T. Winter,⁵² M. Wittgen,¹⁵² M. Wobisch,⁹⁶ A. Wolf,¹⁰⁰ R. Wölker,¹³⁴ J. Wollrath,⁵² M. W. Wolter,⁸⁵ H. Wolters,^{139a,139c} V. W. S. Wong,¹⁷⁴ N. L. Woods,¹⁴⁵ S. D. Worm,⁴⁶ B. K. Wosiek,⁸⁵ K. W. Woźniak,⁸⁵ K. Wraight,⁵⁷ S. L. Wu,¹⁸⁰ X. Wu,⁵⁴ Y. Wu,^{60a} J. Wuerzinger,¹³⁴ T. R. Wyatt,¹⁰¹ B. M. Wynne,⁵⁰ S. Xella,⁴⁰ L. Xia,¹⁷⁷ J. Xiang,^{63c} X. Xiao,¹⁰⁶ X. Xie,^{60a} I. Xioutidis,¹⁵⁵ D. Xu,^{15a} H. Xu,^{60a} H. Xu,^{60a} L. Xu,²⁹ T. Xu,¹⁴⁴ W. Xu,¹⁰⁶ Z. Xu,^{60b} Z. Xu,¹⁵² B. Yabsley,¹⁵⁶ S. Yacoub,^{33a} K. Yajima,¹³² D. P. Yallup,⁹⁵ N. Yamaguchi,⁸⁸ Y. Yamaguchi,¹⁶⁴ A. Yamamoto,⁸² M. Yamatani,¹⁶² T. Yamazaki,¹⁶² Y. Yamazaki,⁸³ J. Yan,^{60c} Z. Yan,²⁵ H. J. Yang,^{60c,60d} H. T. Yang,¹⁸ S. Yang,^{60a} T. Yang,^{63c} X. Yang,^{60b,58} Y. Yang,¹⁶² Z. Yang,^{60a} W-M. Yao,¹⁸ Y. C. Yap,⁴⁶ Y. Yasu,⁸² E. Yatsenko,^{60c,60d} H. Ye,^{15c} J. Ye,⁴² S. Ye,²⁹ I. Yeletsikh,⁸⁰ M. R. Yexley,⁹⁰ E. Yigitbasi,²⁵ P. Yin,³⁹ K. Yorita,¹⁷⁸ K. Yoshihara,⁷⁹ C. J. S. Young,³⁶ C. Young,¹⁵² J. Yu,⁷⁹ R. Yuan,^{60b,oo} X. Yue,^{61a} M. Zaazoua,^{35e} B. Zabinski,⁸⁵ G. Zacharis,¹⁰ E. Zaffaroni,⁵⁴ J. Zahreddine,¹³⁵ A. M. Zaitsev,^{123,i} T. Zakareishvili,^{158b}

N. Zakharchuk,³⁴ S. Zambito,³⁶ D. Zanzi,³⁶ D. R. Zaripovas,⁵⁷ S. V. Zeiβner,⁴⁷ C. Zeitnitz,¹⁸¹ G. Zemaityte,¹³⁴ J. C. Zeng,¹⁷² O. Zenin,¹²³ T. Ženiš,^{28a} D. Zerwas,⁶⁵ M. Zgubič,¹³⁴ B. Zhang,^{15c} D. F. Zhang,^{15b} G. Zhang,^{15b} J. Zhang,⁶ Kaili. Zhang,^{15a} L. Zhang,^{15c} L. Zhang,^{60a} M. Zhang,¹⁷² R. Zhang,¹⁸⁰ S. Zhang,¹⁰⁶ X. Zhang,^{60c} X. Zhang,^{60b} Y. Zhang,^{15a,15d} Z. Zhang,^{63a} Z. Zhang,⁶⁵ P. Zhao,⁴⁹ Z. Zhao,^{60a} A. Zhemchugov,⁸⁰ Z. Zheng,¹⁰⁶ D. Zhong,¹⁷² B. Zhou,¹⁰⁶ C. Zhou,¹⁸⁰ H. Zhou,⁷ M. S. Zhou,^{15a,15d} M. Zhou,¹⁵⁴ N. Zhou,^{60c} Y. Zhou,⁷ C. G. Zhu,^{60b} C. Zhu,^{15a,15d} H. L. Zhu,^{60a} H. Zhu,^{15a} J. Zhu,¹⁰⁶ Y. Zhu,^{60a} X. Zhuang,^{15a} K. Zhukov,¹¹¹ V. Zhulanov,^{122b,122a} D. Zieminska,⁶⁶ N. I. Zimine,⁸⁰ S. Zimmermann,⁵² Z. Zinonos,¹¹⁵ M. Ziolkowski,¹⁵⁰ L. Živković,¹⁶ G. Zobernig,¹⁸⁰ A. Zoccoli,^{23b,23a} K. Zoch,⁵³ T. G. Zorbas,¹⁴⁸ R. Zou,³⁷ and L. Zwalinski³⁶

(ATLAS Collaboration)

¹*Department of Physics, University of Adelaide, Adelaide, Australia*

²*Physics Department, SUNY Albany, Albany, New York, USA*

³*Department of Physics, University of Alberta, Edmonton AB, Canada*

^{4a}*Department of Physics, Ankara University, Ankara, Turkey*

^{4b}*Istanbul Aydin University, Application and Research Center for Advanced Studies, Istanbul, Turkey*

^{4c}*Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey*

⁵*LAPP, Université Grenoble Alpes, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France*

⁶*High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA*

⁷*Department of Physics, University of Arizona, Tucson, Arizona, USA*

⁸*Department of Physics, University of Texas at Arlington, Arlington, Texas, USA*

⁹*Physics Department, National and Kapodistrian University of Athens, Athens, Greece*

¹⁰*Physics Department, National Technical University of Athens, Zografou, Greece*

¹¹*Department of Physics, University of Texas at Austin, Austin, Texas, USA*

^{12a}*Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey*

^{12b}*Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul, Turkey*

^{12c}*Department of Physics, Bogazici University, Istanbul, Turkey*

^{12d}*Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey*

¹³*Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*

¹⁴*Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain*

^{15a}*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*

^{15b}*Physics Department, Tsinghua University, Beijing, China*

^{15c}*Department of Physics, Nanjing University, Nanjing, China*

^{15d}*University of Chinese Academy of Science (UCAS), Beijing, China*

¹⁶*Institute of Physics, University of Belgrade, Belgrade, Serbia*

¹⁷*Department for Physics and Technology, University of Bergen, Bergen, Norway*

¹⁸*Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA*

¹⁹*Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany*

²⁰*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*

²¹*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*

^{22a}*Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá, Colombia*

^{22b}*Departamento de Física, Universidad Nacional de Colombia, Bogotá, Colombia, Colombia*

^{23a}*INFN Bologna and Università di Bologna, Dipartimento di Fisica, Italy*

^{23b}*INFN Sezione di Bologna, Italy*

²⁴*Physikalisches Institut, Universität Bonn, Bonn, Germany*

²⁵*Department of Physics, Boston University, Boston, Massachusetts, USA*

²⁶*Department of Physics, Brandeis University, Waltham, Massachusetts, USA*

^{27a}*Transilvania University of Brasov, Brasov, Romania*

^{27b}*Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania*

^{27c}*Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania*

^{27d}*National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca, Romania*

^{27e}*University Politehnica Bucharest, Bucharest, Romania*

^{27f}*West University in Timisoara, Timisoara, Romania*

^{28a}*Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic*

^{28b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*

²⁹*Physics Department, Brookhaven National Laboratory, Upton, New York, USA*

- ³⁰*Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- ³¹*California State University, Fresno, USA*
- ³²*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ^{33a}*Department of Physics, University of Cape Town, Cape Town, South Africa*
- ^{33b}*iThemba Labs, Western Cape, South Africa*
- ^{33c}*Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa*
- ^{33d}*University of South Africa, Department of Physics, Pretoria, South Africa*
- ^{33e}*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
- ³⁴*Department of Physics, Carleton University, Ottawa ON, Canada*
- ^{35a}*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies—Université Hassan II, Casablanca, Morocco*
- ^{35b}*Faculté des Sciences, Université Ibn-Tofail, Kénitra, Morocco*
- ^{35c}*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- ^{35d}*Faculté des Sciences, Université Mohamed Premier and LTPM, Oujda, Morocco*
- ^{35e}*Faculté des sciences, Université Mohammed V, Rabat, Morocco*
- ³⁶*CERN, Geneva, Switzerland*
- ³⁷*Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA*
- ³⁸*LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France*
- ³⁹*Nevis Laboratory, Columbia University, Irvington, New York, USA*
- ⁴⁰*Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark*
- ^{41a}*Dipartimento di Fisica, Università della Calabria, Rende, Italy*
- ^{41b}*INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy*
- ⁴²*Physics Department, Southern Methodist University, Dallas, Texas, USA*
- ⁴³*Physics Department, University of Texas at Dallas, Richardson, Texas, USA*
- ⁴⁴*National Centre for Scientific Research “Demokritos”, Agia Paraskevi, Greece*
- ^{45a}*Department of Physics, Stockholm University, Sweden*
- ^{45b}*Oskar Klein Centre, Stockholm, Sweden*
- ⁴⁶*Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany*
- ⁴⁷*Lehrstuhl für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*
- ⁴⁸*Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany*
- ⁴⁹*Department of Physics, Duke University, Durham, North Carolina, USA*
- ⁵⁰*SUPA—School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁵¹*INFN e Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁵²*Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany*
- ⁵³*II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany*
- ⁵⁴*Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland*
- ^{55a}*Dipartimento di Fisica, Università di Genova, Genova, Italy*
- ^{55b}*INFN Sezione di Genova, Italy*
- ⁵⁶*II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
- ⁵⁷*SUPA—School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- ⁵⁸*LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France*
- ⁵⁹*Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA*
- ^{60a}*Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei, China*
- ^{60b}*Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao, China*
- ^{60c}*School of Physics and Astronomy, Shanghai Jiao Tong University, KLPPAC-MoE, SKLPPC, Shanghai, China*
- ^{60d}*Tsung-Dao Lee Institute, Shanghai, China*
- ^{61a}*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ^{61b}*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ⁶²*Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*
- ^{63a}*Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China*
- ^{63b}*Department of Physics, University of Hong Kong, Hong Kong, China*
- ^{63c}*Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
- ⁶⁴*Department of Physics, National Tsing Hua University, Hsinchu, Taiwan*
- ⁶⁵*IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France*
- ⁶⁶*Department of Physics, Indiana University, Bloomington, Indiana, USA*
- ^{67a}*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
- ^{67b}*ICTP, Trieste, Italy*

- ^{67c}*Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine, Italy*
^{68a}*INFN Sezione di Lecce, Italy*
- ^{68b}*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
^{69a}*INFN Sezione di Milano, Italy*
- ^{69b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
^{70a}*INFN Sezione di Napoli, Italy*
- ^{70b}*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
^{71a}*INFN Sezione di Pavia, Italy*
- ^{71b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
^{72a}*INFN Sezione di Pisa, Italy*
- ^{72b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
^{73a}*INFN Sezione di Roma, Italy*
- ^{73b}*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
^{74a}*INFN Sezione di Roma Tor Vergata, Italy*
- ^{74b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
^{75a}*INFN Sezione di Roma Tre, Italy*
- ^{75b}*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
^{76a}*INFN-TIFPA, Italy*
^{76b}*Università degli Studi di Trento, Trento, Italy*
- ⁷⁷*Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria*
⁷⁸*University of Iowa, Iowa City, Iowa, USA*
- ⁷⁹*Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA*
⁸⁰*Joint Institute for Nuclear Research, Dubna, Russia*
- ^{81a}*Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora, Brazil*
^{81b}*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
^{81c}*Universidade Federal de São João del Rei (UFSJ), São João del Rei, Brazil*
^{81d}*Instituto de Física, Universidade de São Paulo, São Paulo, Brazil*
- ⁸²*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
⁸³*Graduate School of Science, Kobe University, Kobe, Japan*
- ^{84a}*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
^{84b}*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
⁸⁵*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
⁸⁶*Faculty of Science, Kyoto University, Kyoto, Japan*
⁸⁷*Kyoto University of Education, Kyoto, Japan*
- ⁸⁸*Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan*
⁸⁹*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
⁹⁰*Physics Department, Lancaster University, Lancaster, United Kingdom*
⁹¹*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
- ⁹²*Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia*
- ⁹³*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
⁹⁴*Department of Physics, Royal Holloway University of London, Egham, United Kingdom*
⁹⁵*Department of Physics and Astronomy, University College London, London, United Kingdom*
⁹⁶*Louisiana Tech University, Ruston, Louisiana, USA*
⁹⁷*Fysiska institutionen, Lunds universitet, Lund, Sweden*
- ⁹⁸*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*
⁹⁹*Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain*
¹⁰⁰*Institut für Physik, Universität Mainz, Mainz, Germany*
- ¹⁰¹*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
¹⁰²*CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France*
- ¹⁰³*Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA*
¹⁰⁴*Department of Physics, McGill University, Montreal QC, Canada*
¹⁰⁵*School of Physics, University of Melbourne, Victoria, Australia*
¹⁰⁶*Department of Physics, University of Michigan, Ann Arbor, Michigan, USA*
- ¹⁰⁷*Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA*
¹⁰⁸*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus*
¹⁰⁹*Research Institute for Nuclear Problems of Byelorussian State University, Minsk, Belarus*
¹¹⁰*Group of Particle Physics, University of Montreal, Montreal QC, Canada*
¹¹¹*P.N. Lebedev Physical Institute of the Russian Academy of Sciences, Moscow, Russia*

- ¹¹²*National Research Nuclear University MEPhI, Moscow, Russia*
- ¹¹³*D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia*
- ¹¹⁴*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
- ¹¹⁵*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
- ¹¹⁶*Nagasaki Institute of Applied Science, Nagasaki, Japan*
- ¹¹⁷*Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan*
- ¹¹⁸*Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA*
- ¹¹⁹*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands*
- ¹²⁰*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
- ¹²¹*Department of Physics, Northern Illinois University, DeKalb, Illinois, USA*
- ^{122a}*Budker Institute of Nuclear Physics and NSU, SB RAS, Novosibirsk, Russia*
- ^{122b}*Novosibirsk State University Novosibirsk, Russia*
- ¹²³*Institute for High Energy Physics of the National Research Centre Kurchatov Institute, Protvino, Russia*
- ¹²⁴*Institute for Theoretical and Experimental Physics named by A.I. Alikhanov of National Research Centre “Kurchatov Institute”, Moscow, Russia*
- ¹²⁵*Department of Physics, New York University, New York, New York, USA*
- ¹²⁶*Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan*
- ¹²⁷*Ohio State University, Columbus, Ohio, USA*
- ¹²⁸*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA*
- ¹²⁹*Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA*
- ¹³⁰*Palacký University, RCPTM, Joint Laboratory of Optics, Olomouc, Czech Republic*
- ¹³¹*Institute for Fundamental Science, University of Oregon, Eugene, Oregon, USA*
- ¹³²*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹³³*Department of Physics, University of Oslo, Oslo, Norway*
- ¹³⁴*Department of Physics, Oxford University, Oxford, United Kingdom*
- ¹³⁵*LPNHE, Sorbonne Université, Université de Paris, CNRS/IN2P3, Paris, France*
- ¹³⁶*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA*
- ¹³⁷*Konstantinov Nuclear Physics Institute of National Research Centre “Kurchatov Institute”, PNPI, St. Petersburg, Russia*
- ¹³⁸*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA*
- ^{139a}*Laboratório de Instrumentação e Física Experimental de Partículas—LIP, Lisboa, Portugal*
- ^{139b}*Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal*
- ^{139c}*Departamento de Física, Universidade de Coimbra, Coimbra, Portugal*
- ^{139d}*Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal*
- ^{139e}*Departamento de Física, Universidade do Minho, Braga, Portugal*
- ^{139f}*Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain), Spain*
- ^{139g}*Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal*
- ^{139h}*Instituto Superior Técnico, Universidade de Lisboa, Lisboa, Portugal*
- ¹⁴⁰*Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic*
- ¹⁴¹*Czech Technical University in Prague, Prague, Czech Republic*
- ¹⁴²*Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic*
- ¹⁴³*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ¹⁴⁴*IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France*
- ¹⁴⁵*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA*
- ^{146a}*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- ^{146b}*Universidad Andres Bello, Department of Physics, Santiago, Chile*
- ^{146c}*Instituto de Alta Investigación, Universidad de Tarapacá, Chile*
- ^{146d}*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ¹⁴⁷*Department of Physics, University of Washington, Seattle, Washington, USA*
- ¹⁴⁸*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹⁴⁹*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁵⁰*Department Physik, Universität Siegen, Siegen, Germany*
- ¹⁵¹*Department of Physics, Simon Fraser University, Burnaby BC, Canada*
- ¹⁵²*SLAC National Accelerator Laboratory, Stanford, California, USA*
- ¹⁵³*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
- ¹⁵⁴*Departments of Physics and Astronomy, Stony Brook University, Stony Brook, New York, USA*
- ¹⁵⁵*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
- ¹⁵⁶*School of Physics, University of Sydney, Sydney, Australia*
- ¹⁵⁷*Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ^{158a}*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
- ^{158b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*

- ¹⁵⁹*Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel*
- ¹⁶⁰*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
- ¹⁶¹*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
- ¹⁶²*International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan*
- ¹⁶³*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
- ¹⁶⁴*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
- ¹⁶⁵*Tomsk State University, Tomsk, Russia*
- ¹⁶⁶*Department of Physics, University of Toronto, Toronto ON, Canada*
- ^{167a}*TRIUMF, Vancouver BC, Canada*
- ^{167b}*Department of Physics and Astronomy, York University, Toronto ON, Canada*
- ¹⁶⁸*Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
- ¹⁶⁹*Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA*
- ¹⁷⁰*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
- ¹⁷¹*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
- ¹⁷²*Department of Physics, University of Illinois, Urbana, Illinois, USA*
- ¹⁷³*Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia—CSIC, Valencia, Spain*
- ¹⁷⁴*Department of Physics, University of British Columbia, Vancouver BC, Canada*
- ¹⁷⁵*Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada*
- ¹⁷⁶*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany*
- ¹⁷⁷*Department of Physics, University of Warwick, Coventry, United Kingdom*
- ¹⁷⁸*Waseda University, Tokyo, Japan*
- ¹⁷⁹*Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot, Israel*
- ¹⁸⁰*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
- ¹⁸¹*Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
- ¹⁸²*Department of Physics, Yale University, New Haven, Connecticut, USA*

^aDeceased.

^bAlso at Department of Physics, King's College London, London, United Kingdom.

^cAlso at Instituto de Física Teórica, IFT-UAM/CSIC, Madrid, Spain.

^dAlso at TRIUMF, Vancouver BC, Canada.

^eAlso at Department of Physics and Astronomy, University of Louisville, Louisville, Kentucky, USA.

^fAlso at Physics Department, An-Najah National University, Nablus, Palestine.

^gAlso at Department of Physics, University of Fribourg, Fribourg, Switzerland.

^hAlso at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain.

ⁱAlso at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

^jAlso at Department of Physics, Ben Gurion University of the Negev, Beer Sheva, Israel.

^kAlso at Università di Napoli Parthenope, Napoli, Italy.

^lAlso at Institute of Particle Physics (IPP), Canada.

^mAlso at Dipartimento di Matematica, Informatica e Fisica, Università di Udine, Udine, Italy.

ⁿAlso at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

^oAlso at Borough of Manhattan Community College, City University of New York, New York, New York, USA.

^pAlso at Department of Physics, California State University, Fresno, USA.

^qAlso at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

^rAlso at Centro Studi e Ricerche Enrico Fermi, Italy.

^sAlso at Department of Physics, California State University, East Bay, USA.

^tAlso at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

^uAlso at IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France.

^vAlso at Graduate School of Science, Osaka University, Osaka, Japan.

^wAlso at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany.

^xAlso at University of Chinese Academy of Sciences (UCAS), Beijing, China.

^yAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^zAlso at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands.

^{aa}Also at CERN, Geneva, Switzerland.

^{bb}Also at Joint Institute for Nuclear Research, Dubna, Russia.

^{cc}Also at Hellenic Open University, Patras, Greece.

^{dd}Also at The City College of New York, New York, New York, USA.

^{ee}Also at Department of Physics, California State University, Sacramento, USA.

^{ff}Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland.

^{gg} Also at Louisiana Tech University, Ruston, Louisiana, USA.

^{hh} Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria.

ⁱⁱ Also at Faculty of Physics, M.V. Lomonosov Moscow State University, Moscow, Russia.

^{jj} Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

^{kk} Also at CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France.

^{ll} Also at National Research Nuclear University MEPHI, Moscow, Russia.

^{mm} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

ⁿⁿ Also at Giresun University, Faculty of Engineering, Giresun, Turkey.

^{oo} Also at Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA.