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Leaky-Wave Antenna with Switchable Omnidirectional Conical Radiation via Polarization Handedness

Yahong Liu, Meize Li, Kun Song, Dingshan Hu, Hongchao Liu, Xiaopeng Zhao, Shuang Zhang, and Miguel Navarro-Cía, *Senior Member, IEEE*

Abstract—Reconfigurable antennas capable of beam-steering offer an efficient solution to optimize the use of the crowded wireless medium and can serve as a multifunction antenna. Beam-steering is often achieved by antenna geometry switching at the expense of hardware complexity. Here, polarization is used to realize beam-steering without the need of antenna geometry modification. Depending on the handedness of the feed, backward or forward conical radiation is demonstrated in a $\sim 13\lambda_0$ -long short-circuited helically slotted waveguide antenna. Tapering the slit width with a Taylor distribution reduces the measured sidelobe levels by ~ 3 dB in average and results in a realized gain of 10-13 dB and 11-13 dB for right-handed (backward radiation) and left-handed circularly polarized (forward radiation) feeding, respectively, in the bandwidth from 8.5 to 9.5 GHz.

Index Terms—Beam-steering, circular polarization, helical slit antenna, leaky wave antenna, periodic structure.

I. INTRODUCTION

HELICAL antennas in axial-mode operation are the workhorse of satellite communications because of their immunity to Faraday rotation caused by the ionosphere [1]. In photonics, they have been used to realize negative refraction metamaterials [2], topological metamaterials [3],[4], broadband optical and infrared circular polarizers, absorbers, etc. [5]. When operating in the normal-mode, helical antennas are electrically short monopole antennas with omnidirectional radiation pattern in the plane perpendicular to the helix axis suitable for mobile and portable communication equipment [6] and rocket-borne radiosonde [7]. Helical antennas can also support a conical mode that has found fewer applications in

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communications, but holds promise for generation of orbital angular momentum [8] and Cherenkov radiation applications [9],[10].

The underlying physics of the conical mode of operation is a leaky-wave [1],[6],[11]. Hence, the beam direction can be steered by sweeping the operating frequency. Although frequency-scanning is used in sensing/radar applications [12]-[15] it is not a desirable feature in modern communication systems, which usually require pattern reconfigurability with fixed frequency operation [16],[17].

A leaky-wave antenna that can provide both frequency-scanning and beam-steering for fixed frequency could be an elegant solution for multifunction radiofrequency systems. Such antenna would facilitate the integration of functionalities by reducing the number of antennas and radiofrequency units to be installed. However, such leaky-wave antenna has not been reported yet in the literature. With this motivation in mind and inspired by work done in the field of metamaterials, we propose a leaky-wave antenna with switchable (*front and back semi-infinite planes*) omnidirectional conical radiation via polarization handedness. Polarization has been largely overlooked for reconfigurability purposes in microwave engineering, but is well documented at optics for metamaterials [18]. The proposed leaky-wave antenna is a short-circuited helically slotted waveguide antenna (referred as helical tape waveguide, HTW, from now on). A uniform and a tapered HTW to reduce sidelobe levels (SLL) are reported here. Frequency-scanning of 15° (in an 11% fractional bandwidth) is achieved due to the inherent leaky-wave nature of the antenna. The 15° angle scan occurs in the forward/backward direction for left-handed/right-handed circularly-polarized feeding mode.

II. STRUCTURAL DESIGN

Conventional helical antennas use a metal strip as a radiating element [1]. We will look, however, at their complementary counterpart, the HTW antenna, because it can be fed naturally with a cylindrical waveguide supporting left-handed (LCP) and right-handed circularly-polarized (RCP) modes – *such modes are the superposition of the two degenerate TE_{11} modes with spatially orthogonal fields and ± 90 deg phase difference, respectively* –. Notice that this feeding has not been reported before for helically slotted waveguide antennas

[7],[19], except in our previous work devoted to polarization-induced waveguide filtering and orbital angular momentum beam generation [20]. Unlike Ref. [20], the HTW is terminated in a short circuit to enable leaky-wave functionality for both RCP and LCP. For leaky-wave radiation to occur, the waveguide modes must have opposite handedness to the helix. The waveguide mode that initially has the same handedness (RCP in this work) travels until the end of the HTW and, upon reflection by the terminated mirror, turns into the opposite handedness to the helix, which can radiate into free space. This will be illustrated subsequently in III. Results and Discussion. The dispersion diagram of this short-circuited HTW antenna calculated via the ABCD matrix method [21] can be found in the Appendix. The same concept of handedness matching upon reflection has been exploited for optical tractor beam applications [22].

Two HTW antennas are designed and fabricated with constant pitch $a = 20$ mm ($\sim 0.6\lambda_0$, where λ_0 is the free-space wavelength), diameter $d = 21.6$ mm ($\sim 0.65\lambda_0$) and physical length $l = 430$ mm ($\sim 13\lambda_0$). The pitch angle is $\psi = \arctan(a/\pi d) = 16.4^\circ$. The helix tape is made of copper with thickness $t = 0.5$ mm and wraps a circular resin tube (i.e. hollow cylinder). One HTW antenna is based on a uniform slit of width $s = 5$ mm (Fig. 1(a)), chosen so that approximately 90-95% of the power is radiated [11]. With such uniform slit, the leakage constant α does not change along the length of the HTW, and the aperture distribution has an exponential amplitude variation and a constant phase. Such an aperture distribution results in a high sidelobe level. To reduce sidelobes for both LCP and RCP modes, thereby yielding higher main lobe gain, the second HTW antenna has a tapered slit width $s(z)$ (Fig. 1(b)). Unlike monofilar axial-mode helical antennas [1], the HTW antenna leakage characteristics in the conical mode operation are very sensitive to even moderate changes in dimensions [20], and thus, the effect of moderate departures from uniformity is large. This greatly facilitates the realization of the taper with a Taylor amplitude distribution for the leakage constant α .

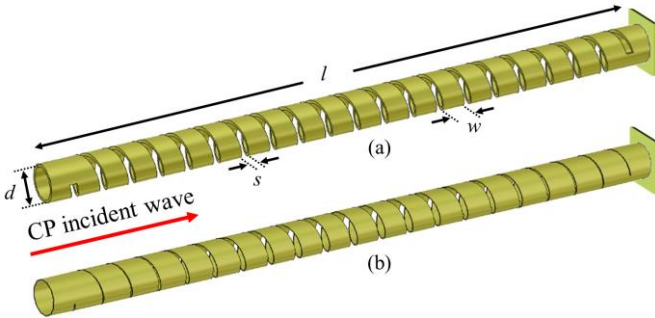


Fig. 1. Sketch of the 20-turn uniform (a) and tapered (b) HTW antenna. The uniform HTW has constant slit width, metallic width and thus lattice constant of $s = 5$ mm, $w = 15$ mm, and $a = s + w = 20$ mm, respectively. The tapered HTW antenna keep constant the spacing between turns $a = 20$ mm, while s and w vary. The dimension of the mirror is $26 \times 26 \times 0.5$ mm³.

The procedure to find out the leakage constant $\alpha(z)$ for the non-uniform HTW relies on the ABCD matrix method [21] and assumes that the local periodicity approximation holds.

This assumption is reasonable here given the adiabatic Taylor distribution envisioned for the non-uniform HTW. Several uniform HTWs are simulated with constant a and d , but different s to find the required value of α for the Taylor distribution. Notice that the circular resin tube is not considered in this design procedure and that metal is modelled as a perfect electric conductor. Hence, the attenuation computed via the ABCD matrix method is indeed the leakage constant and is identical for LCP and RCP modes. The effect of the circular resin tube and dielectric absorption in the radiation characteristics are discussed in the Appendix. Alternatively, the leakage constant could have been calculated as $\alpha = -\frac{1}{2l} \ln\left(\frac{P(l)}{P(0)}\right)$, where $P(l)$ is the power remaining in the waveguide at a distance $z = l$ along the length of the HTW and $P(0)$ is the power input at $z = 0$ [20]. The calculated frequency-dependent leakage constant $\alpha(z)$ for the uniform and non-uniform HTW antenna can be found in Fig. 2. The corresponding $s(z)$ at the beginning of each turn can be found in the Appendix.

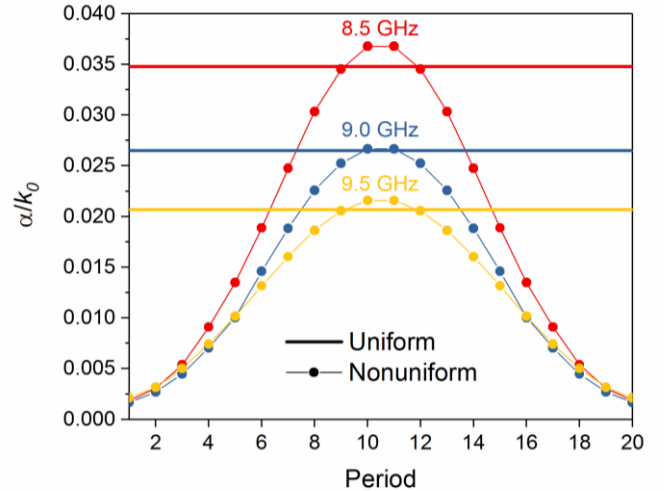


Fig. 2. Simulated leakage constant $\alpha(z)$ for uniform and non-uniform slit width.

III. RESULTS AND DISCUSSION

A. Underlying mechanism and S_{11}

Prior to the measurements, the leaky-wave and polarization switch-ability characteristics of the proposed HTWs can be demonstrated unequivocally with the full-wave simulations. Figure 3 displays the electric field distribution of the two HTW antennas. For the LCP feeding, the energy leaks through the slits with a conical radiation pattern akin to the HTW without mirror [16] and Bull's Eye antennas [23]-[27]. For RCP feeding, the energy travels first along the waveguide almost unperturbedly and after reflecting on the end mirror starts leaking through the slits as it travels toward the feeding with a conical radiation pattern opposite to the LCP feeding. Hence, a standing-wave pattern is observed inside the HTW under RCP feeding unlike the LCP case.

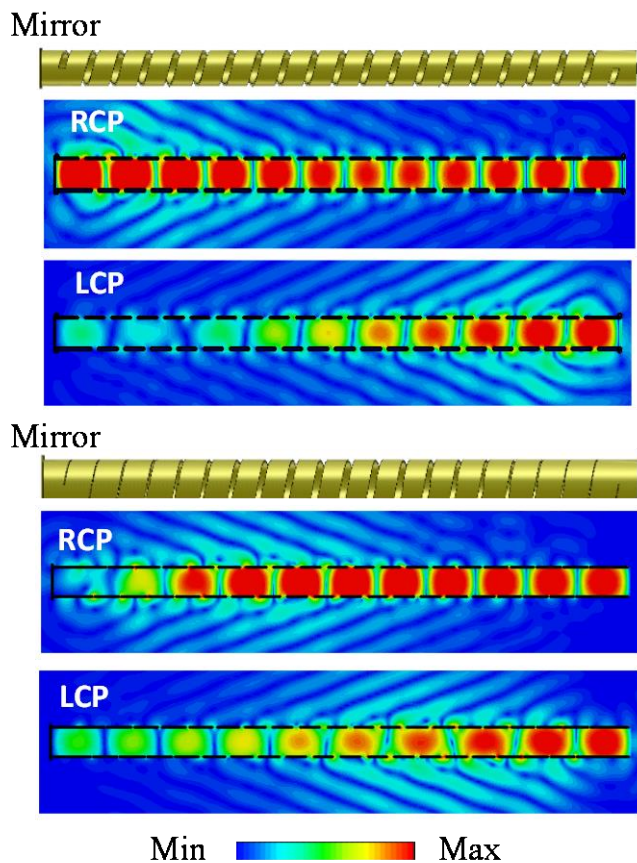


Fig. 3. Simulated electric field distributions for RCP and LCP feeding at 9 GHz: (top) Uniform and (bottom) non-uniform leaky-wave antennas.

Figure 4 shows the measured reflection coefficient S_{11} for both HTW designs under the two circular polarization handedness. The reflection coefficient is plotted from 8.4 GHz upwards since the cutoff frequency of the circular waveguide feeding the HTW antennas is 8.14 GHz. The reflection coefficient for the uniform HTW antenna is below -20 dB for both LCP and RCP feeding. Tapering the slit width worsens the S_{11} , which however still remains below -10 dB in the spectral window shown for both LCP and RCP feeding. Simulations are in reasonable agreement with the measurements given that the full setup and the circular resin tube are not modelled as illustrated in Fig. 8 of the Appendix.

B. Radiation characteristics

For small s , the mode with opposite handedness to that of the helix (RCP feeding here) travels unperturbedly along the HTW before being reflected. However, for moderate values of s such as 5 mm ($\sim 0.15\lambda_0$), this mode can also leak before being reflected by the end mirror. This can be observed indeed in the top contour plots of Fig. 5 where the numerically-computed realized gain is plotted as a function of frequency and radiation angle for the uniform HTW antenna. For RCP (Fig. 5(b)), the main lobes emerge near -70° and 70° . Two prominent sidelobes appear near -110° and 110° . Since these angles are exactly those observed for the main lobes with LCP feeding (Fig. 5(a)), we can conclude that they are residual leakage radiation of the RCP mode as it travels from the

generator toward the end of the waveguide. Further evidence supporting this conclusion comes from the observation that the LCP case does not have any prominent sidelobes near -70° and 70° ; for LCP feeding negligible energy reaches the end of the waveguide and is reflected back to generate such sidelobes.

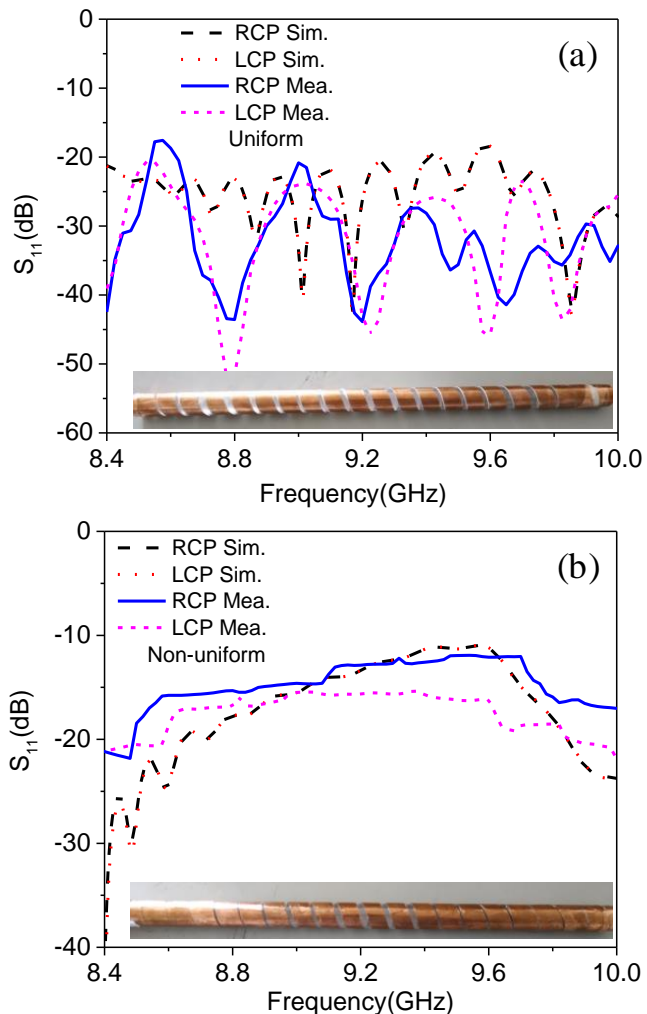


Fig. 4. Reflection coefficient under RCP and LCP feeding for (a) uniform and (b) non-uniform HTW antenna. Insets: picture of the fabricated samples.

When comparing Fig. 5(a) with Fig. 5(b), one can perceive the benefits of the taper. Firstly, the prominent sidelobes for RCP feeding near -110° and 110° are significantly reduced. The Taylor distribution is not meant to address this; this reduction is a byproduct of using consistently smaller s than 5 mm, except for the two central turns. Secondly, the overall background level is reduced. Even though the non-uniform HTW antenna has a worse S_{11} than the uniform one, it displays higher realized gain for both polarizations due to the significant sidelobe level and background reduction achieved with the Taylor taper.

To facilitate the comparison with measurements – done in an anechoic chamber with an AV 3629 vector network analyzer, and a linearly polarized horn antenna (1-18 GHz) –, Fig. 6 presents the far-field radiation patterns in polar form for the uniform and non-uniform HTW antennas at the

representative frequencies of 8.5 GHz, 9 GHz and 9.5 GHz. Overall, simulations agree well with measurements for all cases.

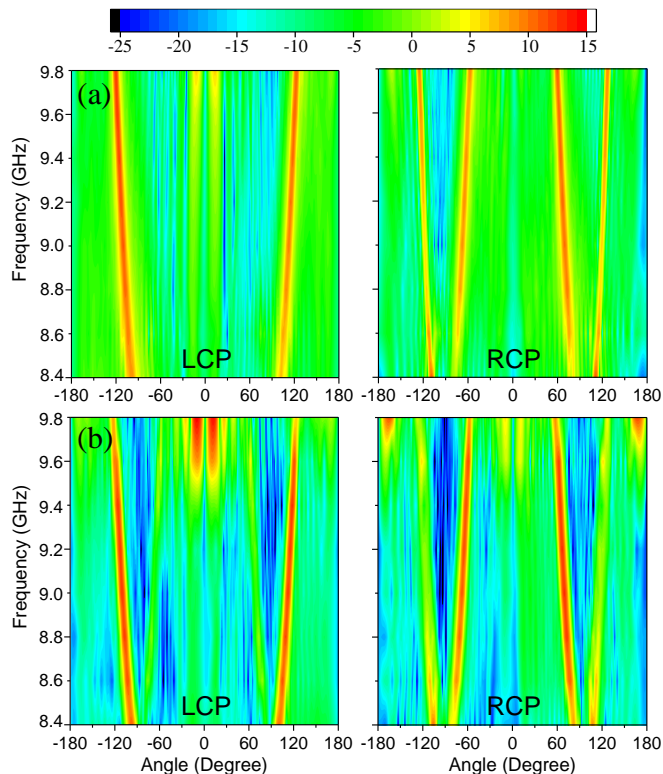


Fig. 5. Realized gain in dB vs. radiation angle and frequency for LCP (left) and RCP (right) feeding: (a) Uniform and (b) non-uniform HTW antenna.

The simulated main lobe gains for the uniform HTW antenna are 11.4 dB and 11.7 dB at 113° and 246° (-114°) for LCP feeding at the central frequency of 9 GHz (left-hand side of Fig. 6(a)). The measured counterparts are 11.1 dB and 11.1 dB at 115° and 250° (-110°) (left-hand side of Fig. 6(c)). For RCP feeding, the main lobe gains are 11.3 dB and 11.6 dB at 67° and 294° (-66°) according to the simulation (right-hand side of Fig. 6(a)), and 11.4 dB and 8.9 dB at 66° and 294° (-66°) according to the measurement (right-hand side of Fig. 6(c)).

As anticipated from Fig. 5, gain levels increase for the non-uniform HTW antenna despite its worse S_{11} . At 9 GHz, the simulated main lobe gains go up to 10.6 dB (111°) and 12.5 dB ($249^\circ = -111^\circ$) for LCP feeding, and 11.3 dB (69°) and 11.2 dB ($292^\circ = -68^\circ$) for RCP feeding (Fig. 6(b)). The measured counterparts (Fig. 6(d)) are 11.8 dB (115°) and 12.6 dB ($250^\circ = -110^\circ$) for LCP feeding, and 12.1 dB (68°) and 10.5 dB ($294^\circ = -66^\circ$) for RCP feeding.

Similar results are obtained at the frequencies of 8.5 GHz and 9.5 GHz. Hence, they are all tabulated in Table I along with the sidelobe levels. Notice that the tabulated gain is the

average gain between the two measured main lobes, whereas the SLL represents the worst case (i.e. the sidelobe level with respect to the minimum gain of either of the two main lobes). The general gain reduction for RCP compared to LCP reported in Table I can arguably link to the additional dissipation loss that the RCP mode experienced, as it has to travel an extra $l = 430$ mm distance before being able to radiate.

TABLE I
MEASURED RADIATION PATTERN CHARACTERISTICS

FREQ. (GHz)	FORWARD BEAM (LCP)				BACKWARD BEAM (RCP)			
	Average Gain (dB)		SLL (dB)		Average Gain (dB)		SLL (dB)	
	uni.	non.	uni.	non.	uni.	non.	uni.	non.
8.5	10.8	11.5	-15.8	-19.3	10.9	10.7	-5.7	-8.3
9.0	11.1	12.2	-14.5	-17.3	10.2	11.3	-12	-11.3
9.5	11.3	12.6	-10.1	-17.1	10.3	12	-8.5	-12.5

Similar conical radiation is achieved with Bull's Eye [23]-[29], and Fabry-Perot antennas [30],[31], albeit, for the reported implementation of those antennas so far, in half of the radiation space without the ability to switch radiation to the other half-space. Setting aside this key aspect, the reported maximum gains of the above-mentioned antennas [23]-[31] overpass the one reported here, but the frequency-scanning within 3 dB scan losses are similar in all cases.

IV. CONCLUSION

To conclude, we have designed and fabricated two leaky-wave antennas based on a short-circuited HTW with a uniform and non-uniform slit operating in the X band. In addition to a frequency-scanning of 15° (in a 11% fractional bandwidth) – 23° in a 17.3% fractional bandwidth if considering the simulations not shown here –, we have demonstrated numerically and experimentally that the radiation direction of the proposed antennas can be switched with the polarization of the feeding. As the LCP mode is incident into the proposed antennas, the power largely leaks through the slits and leads to forward leakage radiation with conical pattern with a measured gain ranging from 11 to 12 dB for the uniform case, and 11-13 dB for the non-uniform case. As the RCP mode is incident on the proposed leaky-wave antenna, the power largely leaks through the slits into a backward conical radiation pattern with a measured gain ranging from 9 to 11 dB for the uniform case, and 10-13 dB for the non-uniform case. The reconfigurability of the proposed HTW antennas without additional hardware complexity and their simple waveguide design make them attractive for the next generation of communication and sensing systems at (sub-)millimeter-waves.

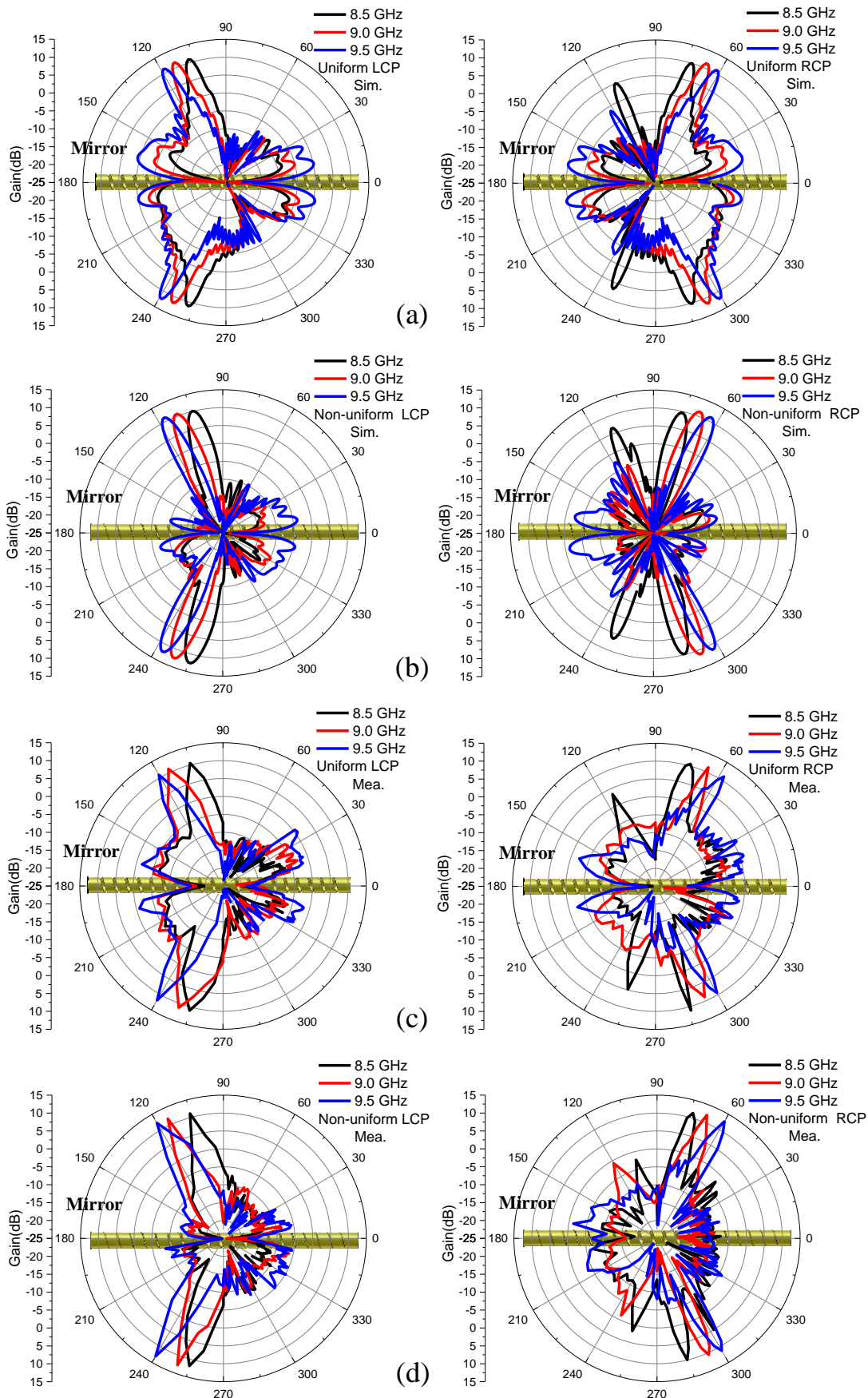


Fig. 6. (a) Far-field radiation patterns at different frequencies for LCP (left) and RCP (right) incident waveguide mode. Simulated (a) uniform and (b) non-uniform leaky-wave HTW antennas, and measured (c) uniform and (d) non-uniform leaky-wave HTW antennas.

APPENDIX

A. Dispersion diagram

The complex dispersion diagram shown below (Fig. 7) is computed with the frequency-domain solver of CST Microwave Studio® and using the ABCD matrix method from a single unit cell and with a macro-cell made up of 2, 3 and 5 unit cells to [21]. More technical details about the simulation can be found in Appendix B. Figure 7 reveals that the phase constant is more robust against changes in s than the leakage constant. This facilitates the design of the non-uniform HTW antenna since we can assume that each period will radiate at almost the same angle. Indeed, the maximum deviation expected for the radiation angle due to the range of propagation constants for $s(z)$ is 2.5° , 1.5° and 1.2° at 8.5, 9 and 9.5 GHz, respectively. Specifically, for 9 GHz, this dispersion diagram predicts the main lobes to emerge between $114.9^\circ/245.1^\circ$ and $116.4^\circ/243.6^\circ$ for LCP, whereas for RCP they are expected between $63.6^\circ/294.9^\circ$ and $65.1^\circ/296.4^\circ$, which agree very well with the time-domain simulations of the fabricated samples and the measurements.

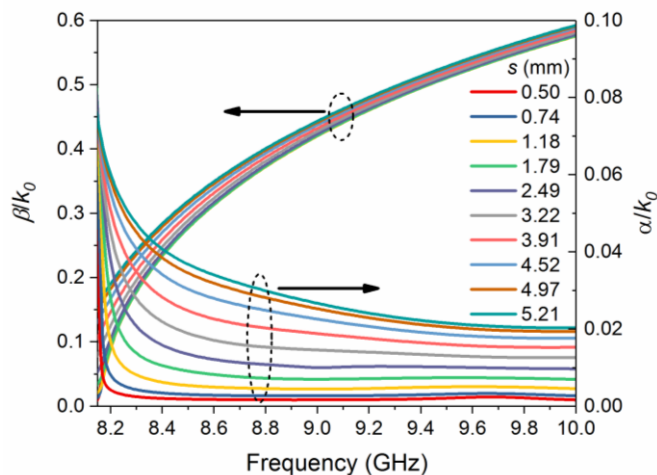


Fig. 7. Dispersion diagram of the short-circuited HTW antenna along with the light line.

B. Simulations

We use the commercial software CST Microwave Studio® to design and model the HTW antennas. Both, frequency-domain and transient solvers are used for the results reported here. The Eigen-mode solver is also used for validation purposes, but its results are not reported. Notice that the coaxial-to-rectangular transition and the rectangular-to-circular transition used in the measurements (Fig. 8) are not modelled numerically with either of the solvers, which may be the reason of the minor disagreements between simulations and measurements in Figs. 4 and 6. Open boundary conditions are applied to the x -, y -, and z -directions. Metal is modelled as a perfect electric conductor. Given the 0.22 mm thick resin tube's small electrical thickness ($\sim 0.007\lambda_0$), it is considered unnecessary to model the tube in the simulations for the main body of the manuscript such to reduce computational burden. Indeed, such negligible influence is demonstrated in Table II

where the radiation pattern characteristics with and without lossy resin tube ($\epsilon = 2.2 - j0.001$) are tabulated. The solver's adaptive mesh (tetrahedral and hexahedral in the frequency- and time-domain solvers, respectively) is employed with a convergence accuracy of the S-parameters of 0.001.

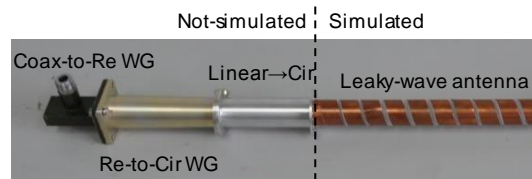


Fig. 8. Picture of the measured prototype highlighting the simulated section.

TABLE II
SIMULATED RADIATION PATTERN CHARACTERISTICS WITH AND WITHOUT LOSSY RESIN TUBE

FREQ. (GHz)	FORWARD BEAM (LCP)				BACKWARD BEAM (RCP)			
	Average Gain (dB)		SLL (dB)		Average Gain (dB)		SLL (dB)	
	w/	w/o	w/	w/o	w/	w/o	w/	w/o
8.5	11.4	11.3	-27.5	-28.0	9.7	9.7	-26.5	-27.0
9.0	12.1	11.6	-15.9	-15.6	11.4	11.3	-15.0	-15.6
9.5	12.3	11.7	-10.9	-11.3	11.6	11.3	-11.4	-11.3

C. Slit width dimensions for the tapered HTW antenna

TABLE III
SLIT WIDTH AT THE BEGINNING OF EACH PERIOD

TURN	Slit width (mm)
1, 20	0.5
2, 19	0.74
3, 18	1.18
4, 17	1.79
5, 16	2.49
6, 15	3.22
7, 14	3.91
8, 13	4.52
9, 12	4.97
10, 11	5.21

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