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# On the Synthesis of Morphology Controlled Transition Metal Dichalcogenides via CVD for Electrochemical Hydrogen Generation

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## Abstract

Shape engineered atomically thin transition metal dichalcogenide (TMD) crystals are highly intriguing systems in regards to both fundamental and applied science. Here, a chemical vapour deposition assisted generalized synthesis strategy for the triangular and dendritic shaped TMDs and their ternary alloys is proposed, and the TMD structures' potential for electrocatalytic hydrogen evolution reaction (HER) applications is demonstrated. The alloy formation is confirmed *via* micro-Raman and photoluminescence studies, and further verified by transmission electron microscopy and X-ray photoelectron spectroscopy. The HER activities of MoS<sub>2</sub> and MoSe<sub>2</sub> triangles are compared with those of their dendritic structures, and an enormous improvement in terms of overpotential and current density is observed for the dendritic structures. A further enhancement of the HER activity is observed in MoS<sub>2(1-x)</sub>Se<sub>2x</sub> triangular and dendritic structures, with dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> providing the best activity. The demonstrated non-equilibrium growth technique opens new avenues for the synthesis of morphology controlled, large area, complex, atomically thin TMD structures s, which can have unprecedented properties such as the enormous catalytic activity, tunable luminescence etc., as presented here.

Engineering the structure and morphology of atomically thin layers (two dimensional/2D materials) is receiving tremendous attention due to both fundamental and applied interest. The quantum confinement effects can yield drastic and unprecedented changes in the physical properties of 2D layers. It is found that the nature and the position of Dirac cone in graphene, the indirect to direct band gap transition, and the spin orbit coupling of transition metal dichalcogenides (TMDs) are tunable by controlling the morphology. Through this, structure and morphology control of TMDs can have high impact on their electronic, optical, and photo-/electrocatalytic properties. Hence, methods for the controlled growth of complex TMD structures are in high demand.

Several self-organized patterns of atomic layers were reported in the recent past, including structures such as snowflake-like graphene structures, [10] dendritic graphene [11] and TMD structures, [18] polygons of TMDs [12] and their alloys. [13] Growing these complex structures of micrometer length scales in a controllable and uniform fashion is rather challenging. [10, 14] Developing their ternary, quaternary, or higher order mixtures having different metal/chalcogen dopants in order to engineer their physical properties is an active area of research. [15] Top-down growth methods such as liquid exfoliation, [16] mechanical exfoliation [17] etc. are not viable for the morphology controlled 2D material synthesis since most of these engineered structures are only form under non-equilibrium growth conditions. The chemical vapour deposition (CVD) technique is ideal for such morphology controlled structures since the growth process in this case can be tuned both kinetically and thermodynamically by controlling the pressure temperature profile. [11-12]

TMDs, especially MoS<sub>2</sub>, are well researched for their catalytic activity, particularly for their electro-/photo-electro catalytic hydrogen generation (HER).<sup>[18]</sup> Though thermodynamically

many TMDs favor HER ( $\Delta G_{H^{\sim}}$  0eV), [19] their catalytic activity is limited kinetically due to the basal plane's inactivity. [20] To circumvent this issue, either the basal plane has to be made active or more edge exposure needs to be achieved for enhancing the exchange current density. [8, 20] The development of size confined structures such as quantum dots is an approach to make large edge exposed TMDs, but the chemical techniques employed to generate such confined structures result in oxidized edge states, hence hamper the charge transfer properties.<sup>[21]</sup> On the other hand, Wenshuo Xu et al. have recently reported a CVD based growth technique for dendritic MoS<sub>2</sub> structures, exhibiting an enhanced number of edge sites. [22] They controlled the Mo/S vapour concentration by varying the height of the Si/SiO<sub>2</sub> substrate in the CVD chamber and demonstrated that a low concentration of Mo/S vapors gives rise to dendritic structures. Further, it was also found that the dendritic structures of TMDs have higher HER electro-catalytic activity than their triangular structures (the most common polygon shape obtained during monolayer MoS<sub>2</sub> growth) because of the presence of enhanced number of edge sites in these structures. Similarly, dendritic MoS2 structures can also be generated by varying the growth temperature, where it is found that MoS<sub>2</sub> flakes can vary from highly fractal to relatively compact as the growth temperature varies from low to high (730 °C to 880 °C). It is believed that higher growth temperatures enhance the surface migration of precursors and the crystallization of MoS<sub>2</sub> flakes, thus yielding more compact structures. [23]

Along with the edge plane electrocatalytic activity of TMDs, there are both theoretical and experimental reports on the tunability of the basal plane's activity of TMDs by forming their ternary or quaternary alloys (in-plane mixtures of different metals and/or chalcogens in different ratios)<sup>[13, 20]</sup>. Lei Yang *et al.* have reported the formation  $MoS_{2(1-x)}Se_{2x}$  alloys using CVD, where x is varied from 0 to 1 depending on the amount of S and Se in the CVD chamber.<sup>[24]</sup> Their

studies showed that fractional values of x show improved catalytic activity compared to integral x. The comparison of their HER activity with that of pristine samples (MoS<sub>2</sub>/MoSe<sub>2</sub>) showed that the enhanced activity contribution is predominantly coming from the basal plane of the alloys. Hence, it can be concluded that the HER activity of TMDs can be improved by making alloys as well as engineering the edges in the structure. In the similar direction, Feng *et. al* reported a physical vapor deposition method for morphologically controlled Mo-zigzag and S/Se zigzag edged oriented single layer crystals by tuning the growth temperature and temperature gradient in the quartz tube.<sup>[13]</sup> But the morphologies reported by Feng *et al.* were straight edged triangles, hexagons, and triangles with inward curving edges. A generalized method for the growth of morphology controlled different TMDs structures is still lacking and here we are proposing a protocol for the growth of several microns (area) sized triangle and dendritic structures of TMDs and their alloys.

Atomically thin, triangle structured MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> were grown over 300 nm thick oxidized Si wafers (Si/ SiO<sub>2</sub>). A customized two zone furnace based CVD set up was used for the growth, as shown in figure 1. Ar/H<sub>2</sub> (10% H<sub>2</sub>) gas was used as the carrier gas and also providing an inert atmosphere. The flow rate was kept as 200 sccm during the entire growth process. The chalcogen S and Se powders were placed in the lower temperature zone with the temperature of 200 °C and 350 °C for MoS<sub>2</sub> and MoSe<sub>2</sub>, respectively. 6 mg of MoO<sub>3</sub> powder in an alumina crucible and the SiO<sub>2</sub>/Si wafer suspended over it were kept in the high temperature zone (as shown in figure 1) with a temperature of 650 °C and 750 °C for MoS<sub>2</sub> and MoSe<sub>2</sub>, respectively. The times taken for zone-II to reach 750 °C (for MoSe<sub>2</sub> or 650 °C for MoS<sub>2</sub>) and zone-I to reach 350 °C (for Se or 200 °C for sulfur) were kept as 25 minutes. The growth time was kept as 20 minutes for all the triangle systems. For alloys (triangle), Se was

placed at 350 °C and S was placed at the corner of lower temperature zone-I (same side) where the temperature was optimized to be 200 °C. The high temperature zone was set as 720 °C where the MoO<sub>3</sub> - SiO<sub>2</sub>/Si substrate system was placed as mentioned before. The growth time was kept as 20 minute for the alloys too. In all cases, the furnace was suddenly opened after 20 minutes to arrest the growth of TMDs by sudden cooling.

For the growth of dendrites, the respective growth temperatures were kept the same as above (650 °C for MoS<sub>2</sub>, 750 °C for MoSe<sub>2</sub>, and 720 °C for alloys) while the growth time and the flow rates were changed. In this method, the flow of Ar/H<sub>2</sub> was carried out from both the sides during the growth, as shown in the figure 1a & 1b. A higher flow rate of 350 sccm was kept during the growth of dendritic structures and the entire growth time was kept as 9 minutes. The precursor arrangement for MoS<sub>2</sub>/MoSe<sub>2</sub> dendrite was kept exactly the same as that of triangle growth while a higher flow rate was kept as mentioned before. After the growth of initial 6 minutes, the direction of Ar/H<sub>2</sub> flow (indicated by red arrows) was reverted immediately to the opposite direction (green arrows) for the next two minutes. The growth was then arrested by the sudden cooling of the furnace. For the alloy dendrite growth, the precursors were kept in the same manner as for the triangled alloy structures, while the flow rate and the direction were tuned as mentioned for MoS<sub>2</sub>/MoSe<sub>2</sub> dendritic growth.

The as grown TMDs were transferred from the SiO<sub>2</sub>/Si substrate to glassy carbon (GC, 3 mm diameter) electrodes using a polymer assisted transfer technique. In this method, the CVD grown TMD/SiO<sub>2</sub>/Si substrate was spin-coated with polymethyl methacrylate (PMMA) and then kept for drying. The PMMA coated TMDs/SiO<sub>2</sub>/Si substrate was dipped in 2 M KOH solution for one day to etch away the SiO<sub>2</sub> layer. This will result the separation of PMMA/TMDs from the Si substrate (the Si substrate will then sink in the KOH solution container while the

PMMA/TMDs film will float on the top). PMMA/TMDs was further scooped out using the GC electrode. The scooped out PMMA/TMDs thin film on GC electrode was thoroughly washed in triple distilled water to remove the KOH. The GC electrode was then left to air dry and the PMMA was washed off with acetone and isopropanol, successively. These GC/TMDs systems were used as working electrode in the electrochemical HER measurements (see method section for the details of electrochemical set up).

Water bubblers are connected *via* exhaust tubes to both ends of the tubular chamber to trap the toxic by-products formed during the reaction. In order to achieve the dendritic structures low vapour pressure of Mo/S is required <sup>[22]</sup>. This has been achieved by a simple modification of the conventional CVD based TMD growth technique, as shown in the schematic figure 1a. The time-temperature/flow rate profile for the dendrite growth is shown in figure 1b. The opposite flow is indicated by the negative flow rate values in the profile. The suddenly changed opposite direction of the Ar/H<sub>2</sub> gas flow will lead to the low Mo/S vapour pressure and thus it is hypothesized to form dendritic structures. In general, single layer TMDs are formed in a triangular fashion, which is thermodynamically favourable, while an immediate change in the gas flow direction will result in a growth mode that is kinetically favoured.

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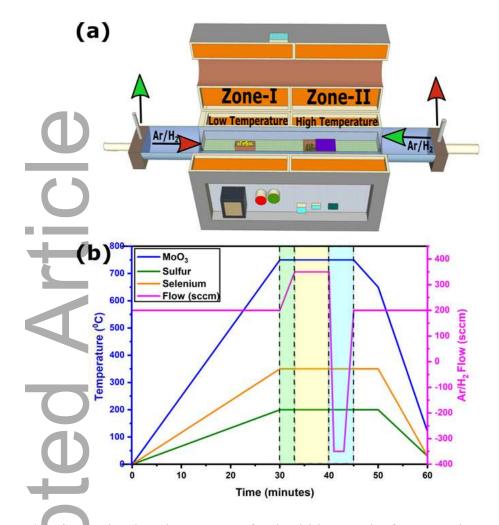


Figure 1. a) In house developed CVD setup for dendritic growth of TMDs. The arrows represent the Ar/H<sub>2</sub> flow and its exit in each directions. b) Time-temperature profile for the growth of dendrites. Note that the negative flow rate of Ar/H<sub>2</sub> indicates opposite direction of Ar/H<sub>2</sub> flow.

Figure 2 (a-c) shows optical images of the resultant  $MoS_2$ ,  $MoSe_2$ , and  $MoS_{2(1-x)}Se_{2x}$ , with triangular shaped grains whereas figure (d-f) represents the corresponding dendritic morphology structures. Figure 2(g-i) shows scanning electron microscope (SEM) images of  $MoS_2$ ,  $MoSe_2$ , and  $MoS_{2(1-x)}Se_{2x}$  dendrites, respectively. The optical images show that these dendrites are uniformly monolayer (by optical contrast) with a large number of edge sites. From the SEM analyses, lateral size of the triangles and dendrites of the TMDs are found to be  $\sim 30$ -

100 μm. The large area coverage and uniform growth are further verified by low resolution optical and SEM images, as shown in the supporting information, figure S1 and S2 respectively.

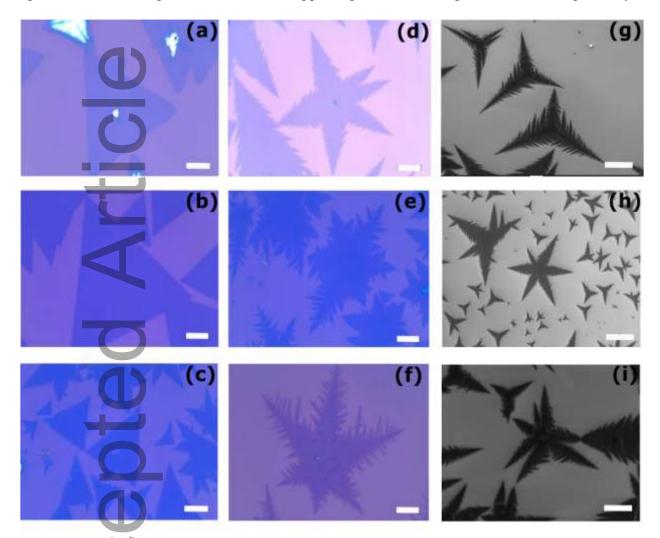
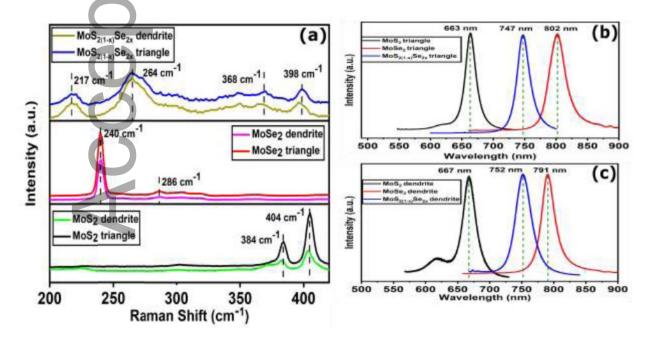


Figure 2: Optical microscope images of triangular morphology of (a) MoS<sub>2</sub> (b) MoSe<sub>2</sub> and (c) MoS<sub>2(1-x)</sub>Se<sub>2x</sub> and dendritic Morphology of (d) MoS<sub>2</sub> (e) MoSe<sub>2</sub> and (f) MoS<sub>2(1-x)</sub>Se<sub>2x</sub>. SEM images of (g) MoS<sub>2</sub>, (h) MoSe<sub>2</sub> (i) MoS<sub>2(1-x)</sub>Se<sub>2x</sub> dendrites. All scale bars are 20 μm.

The crystalline quality and alloying of the TMD structures are confirmed using micro Raman spectroscopy and mapping. Figure 3a shows the comparison of Raman spectra from triangular and dendritic  $MoS_2$ ,  $MoSe_2$ , and  $MoS_{2(1-x)}Se_{2x}$ . TMDs show significant and

dominating Raman peaks originating from the in-plane ( $E_{2g}$ ) and out of plane ( $A_{1g}$ ) vibration modes of chalcogens.<sup>[25]</sup> It is found that there are no significant changes in peak position and line widths of the Raman modes from dendrite to triangle structures of the same alloy composition. The separation between the  $E_{2g}$  and  $A_{1g}$  vibrational modes is very sensitive to the number of layers of a particular TMD. In the case of  $MoS_2$  and  $MoSe_2$ , the separations between these two vibration modes are found to be  $\sim 20$  and  $\sim 46$  cm<sup>-1</sup>, respectively, indicating the presence of respective single layers.<sup>[26]</sup>

The  $MoS_{2(1-x)}Se_{2x}$  alloy has four dominating peaks of which two are the  $E_{2g}$  and  $A_{1g}$  peaks from S-Mo--S bonds at 368 cm<sup>-1</sup> and 398 cm<sup>-1</sup>, respectively. The other two peaks at 217 cm<sup>-1</sup> and 264 cm<sup>-1</sup> are due to vibrations from partially selenized  $MoS_2$  (S--Mo--Se). [27] Even though the alloy structures are found to be single monolayers, the separation between  $E_{2g}$  and  $A_{1g}$  modes of  $MoS_2$  is around 30 cm<sup>-1</sup> due to strain induced effects in the lattice *via* Se doping. Hence, the Raman studies confirm the formation of alloys, both in dendritic and triangle shaped structures.



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Figure 3: a) The micro-Raman spectra of MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> triangular and dendrite structures(532 nm laser excitation for all ). Photoluminescence (PL) from MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> (b) triangular, and (c) dendritic structures. (633 nm laser (excitation) was used for MoSe<sub>2</sub> and MoSe<sub>2</sub> and MoSe<sub>2</sub> and MoSe<sub>2</sub> and Solution for MoSe<sub>2</sub> PL).

One of the important features of atomically thin TMDs, which makes them invaluable materials for optoelectronic materials, is their high photoluminescence (PL) due to the conversion of indirect bandgap to direct one while limiting the number of layers to one monolayer. [28] It is found that  $MoS_2$ ,  $MoSe_2$ , and  $MoS_{2(1-x)}Se_{2x}$  show PL at 663, 802 and 747 nm respectively, which shows that these alloys have great potential in the band gap tunability of TMDs. The PL from the  $MoS_2$  has two clear peaks A (dominating one) and B (small peak), correspond to direct band transitions at the K point of Brillouin zone. [29] Two peaks come from the valence band splitting in monolayer TMDs due to the presence of strong spin-orbit coupling. The difference between these two peaks (A and B) yields a valence band splitting of  $\sim 0.2$  eV [29] and 0.25 eV[30] for  $MoS_2$  and  $MoSe_2$ , respectively. The PL mapping is carried out on the dendritic structures (supplementary information, figure S3) indicating the uniformity of these structures at a micrometer scale.

In order to further estimate the presence of 'Se' and the local environment of chemical composition in pristine  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$  alloys, X-ray photoelectron spectroscopy (XPS) measurements were performed. Figure 4 shows the core level XPS bands of Mo-3d, S-2p, Se-3d, and O-1s in  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$ . A XPS survey spectrum and the corresponding peak positions and their binding energies are shown in the supporting information (figure S4), indicating the presence of Mo, S, Se, and O. Figure 4a shows the comparison of Mo-3d core level XPS spectra in  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$ , respectively. As shown in figure 4a, the strongest

Mo-3d doublet peaks for [MoS<sub>2</sub>] and [MoS<sub>2(1-x)</sub>Se<sub>2x</sub>] are detected at [230.0 eV (3d<sub>5/2</sub>) & 233.1 eV (3d<sub>3/2</sub>)], and [229.6 eV & 232.8 eV], respectively corresponding to the 4+ oxidation state. The characteristic peaks corresponding to S  $2p_{3/2}$  and S  $2p_{1/2}$  orbitals of the divalent sulphide ions (S<sup>2</sup>-) in [MoS<sub>2</sub>] and [MoS<sub>2(1-x)</sub>Se<sub>2x</sub>] are located at [162.8 eV & 164.0 eV] and [162.5 eV & 163.7 eV, respectively. Besides, an additional peak is observed at 161.2 eV, which may be due to the formation of S-Se and amorphous sulphur in the  $MoS_{2(1-x)}Se_{2x}$  system. [31] Interestingly, the observed negative binding energy shifts in both Mo-3d  $(3d_{5/2}, 3d_{3/2})$  and S-2p  $(2p_{3/2}, 2p_{1/2})$  around 0.3 eV in  $MoS_{2(1-x)}Se_{2x}$  with respect to  $MoS_2$  can be due to the alloying of 'Se' with  $MoS_2$  basal plane lattice (which is later proven using TEM). It is well known that when a less electronegative element is added to the system there will be negative binding energy shift.<sup>[32]</sup> Since Se is less electronegative than S, the shift in XPS can be attributed to the alloying. In particular, the formation of MoS<sub>2(1-x)</sub>Se<sub>2x</sub> is identified by Se-3d, which exhibits two doublet peaks 3d<sub>5/2</sub>, 3d<sub>3/2</sub> detected at 54.9 eV, 55.8 eV, respectively, as shown in figure 4c. The percentage (atomic%) of Se on  $MoS_2$  is found to be ~34%, indicating that x is 0.34. Figure 4d displays the O-1s spectra from the  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$  samples. The O-1s spectrum of  $MoS_2$  is peaked at 529.3 eV whereas that in MoS<sub>2(1-x)</sub>Se<sub>2x</sub> alloy is peaked at 529.5 eV. The relative intensities of the O-1s spectra in the MoS<sub>2</sub> and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> samples are almost identical. The nature of the spectra show that the presence of oxygen might be due to the formation of physisorbed oxygen during the atmospheric CVD growth process.

The Se distribution over the  $MoS_2$  basal plane is further estimated from high angle annular dark field (HAADF) - scanning transmission electron microscope (STEM), operating at 80 kV, as shown in figure 4(e-f). Figure 4(e) shows the HAADF-STEM images of pristine monolayer  $MoS_2$  and that of  $MoS_{2(1-x)}Se_{2x}$  alloy is shown in figure 4(f). The bright spots

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correspond to Mo sites and columns are filled with S atoms in case of  $MoS_2$  as shown in figure 4e. It is to be noted that the HAADF intensity of S atoms is rather low, which may be due to the marginal contamination from PMMA residue during the electron beam exposure. It is very important to understand the Se distribution in the MoS<sub>2</sub> layers as it affects the structure's properties. For example, a material in which subsequent atomic layers have different dopant concentrations will have different properties from the one consisting of identical alloy layers. [33] Atomic scale chemical analysis using HAADF-STEM imaging is more challenging for thicker samples than monolayers due to the overlapping of atoms in the different layers.<sup>[34]</sup> An example of consecutive mono and bi layer AB stacked Se doped MoS2 is shown in the supporting information in figure S5(a-b), which shows that a drastic variation of image contrast (see dark and bright field STEM image) can be observed at the atomic sites which makes visual interpretation in stacked layers challenging. Note that the HAADF-STEM imaging is performed on several locations of the monolayer  $MoS_{2(1-x)}Se_{2x}$  sample to estimate Se doping as shown in the supporting information figure S6. HAADF-STEM imaging clearly shows the formation of atomically thin crystalline layers.

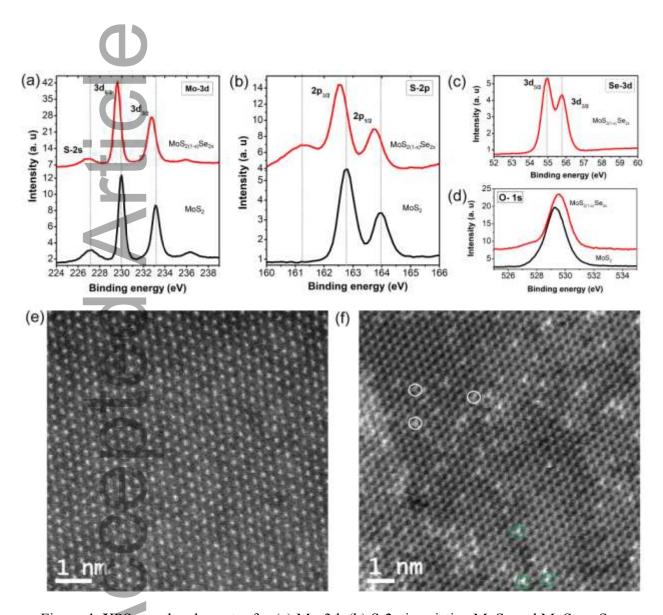


Figure 4: XPS core level spectra for (a) Mo-3d, (b) S-2p in pristine  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$  alloys, (c) Se 3d in  $MoS_{2(1-x)}Se_{2x}$ . (d) O-1s core level XPS spectra of  $MoS_2$  and  $MoS_{2(1-x)}Se_{2x}$ . HAADF-STEM images of a monolayer (e)  $MoS_2$  and (f)  $MoS_{2(1-x)}Se_{2x}$ . Note that the white and green circles denote the Se+S and  $Se_2$  atomic sites, respectively.

Gong *et. al* reported the quantitative site-separated histogram analysis for Se distribution in mono and bilayer MoS<sub>2</sub> using HAADF-STEM image intensity from atomic arrangement of Mo, S<sub>2</sub>, Se+S, and Se<sub>2</sub> sites.<sup>[33]</sup> Based on their analysis, Mo atoms display brighter intensity than the S<sub>2</sub> sites in pristine MoS<sub>2</sub> hexagonal lattice while in the case of Se doped MoS<sub>2</sub> some S<sub>2</sub> sites become much brighter and display intensities close to or even higher than the Mo sites with the intensities of the Mo sites remaining unchanged. They quantitatively evaluated that higher image intensity at the S<sub>2</sub> sites arises from Se substitution (Se+S) and Se<sub>2</sub>, which are brighter than Mo.<sup>[33]</sup> Interestingly, we have found similar observations from our HAADF-STEM image analysis as shown in figure 4 (e-f), in which S<sub>2</sub> is substituted with single Se (Se+S) highlighted with white rings and double sites (Se<sub>2</sub>) are highlighted with green rings (see figure 4(e)).

In order to assess the impact of structural (morphological) variations on the inherent electrocatalytic hydrogen evolution reaction (HER) properties of these TMDs, the CVD grown monolayer triangular as well as dendritic structures of MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> were first transferred as such onto GC electrodes (as mentioned before). For instance, figure S7 shows the optical and SEM images of transferred dendritic  $MoS_{2(1-x)}Se_{2x}$  structure from  $SiO_2/Si$  substrate to GC electrode along with their Raman and PL spectra. Remarkably, the shapes of the dendritic structure seem nearly intact before to (figure S7a and b) and after the transfer (figure S7c and d) process, suggested the perfect preservation of the edge sites. Similarly, the Raman (figure S7e) and PL ((figure S7f) spectra of transferred  $MoS_{2(1-x)}Se_{2x}$  flakes on GC are matching well with that of monolayer  $MoS_{2(1-x)}Se_{2x}$  as grown on  $SiO_2/Si$ . The red shift in A excitonic peak in PL spectra before and after the transfer process can be attributed to the released interfacial strain effect during wet transfer procedure or changes in the dielectric environment. The TMDs transferred onto GC electrodes were further systematically explored for HER studies using linear

scan voltammetry (LSVs) (details in experimental section) using a three electrode system with graphite rod (Pt counter electrode is also checked and got similar results) as counter electrode. It is to be noted that the LSV polarisation curves presented here are not iR compensated, since it is known that uncompensated data will be more meaningful while comparing the activity of two different materials at the same current density using the similar electrochemical set up.<sup>[35]</sup>

The HER polarization curves were taken in 0.5 M H<sub>2</sub>SO<sub>4</sub> at a scan rate of 2 mV s<sup>-1</sup> revealing the respective HER performances, as depicted in figure 5a. The current densities were normalized by the electrochemical active surface area (calculated using Randles-Sevcik equation, details in supporting information). The HER performances of the bare GC and benchmarked commercially available Pt/C catalyst are also performed comparably. As observed (figure 5a), the TMDs modified GC electrodes are HER active, presenting less onset overpotential  $(\eta_0)$  and high current densities in compare to bare GC. Among the TMDs, the monolayer triangle  $MoS_2$  exhibits a poor HER activity (high  $\eta_0$ ), whereas the monolayer dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> alloy exhibits the highest HER activity and have the lowest onset overpotential and high current densities than others. The variations of onset overpotential  $(\eta_0)$ and overpotential at current density 10 mA cm<sup>-2</sup> ( $\eta_{10}$ ) of all electrodes are summarized in the figure 5b. It is also worth noting that the triangular shaped monolayer TMD flakes are less HER active than the respective dendritic structures. To achieve a current density of 10 mA cm<sup>-2</sup> (a benchmarked value while comparing the solar to hydrogen conversion), the triangle shaped MoS<sub>2</sub>, MoSe<sub>2</sub> and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> require overpotentials of 540, 485 and 349 mV, respectively whereas for the same current density, the dendritic structure of MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> require overpotentials of only 394, 326 and 241 mV, respectively. As observed, the  $\eta_{10}$  shifts positively when structure of TMD flakes changes from triangular to dendritic structure, which is

expected due to fact that the electrocatalytic activities of TMDs are strongly related to it exposed edge sites.<sup>[22-23]</sup> It has been well reported theoretically as well as experimentally that edge sites of 2D TMDs are catalytic active for HER due to their low hydrogen binding energy (ΔG<sub>H</sub>) whereas the basal planes are catalytically inert. [20] Additionally, it has been also reported that the ternary TMD alloys exhibits better HER performance than that of their binary counterparts due to presence of different radii of atoms on the basal planes which can induced the lattice distortion and can create many more unsaturated active sites. [36] Thus, lattice distortion in  $MoS_{2(1-x)}Se_{2x}$ alloy may further leads the active basal planes too along with its active edge site. Therefore, the high activity of MoS<sub>2(1-x)</sub>Se<sub>2x</sub> dendritic structure can be ascribed due to the synergic effects of abundant catalytic active edge sites as well as their active basal plane. Furthermore, in order to confirm the inherent HER activity of TMDs is not due to other artefacts, the current densities are also calculated using the geometric surface area of the electrode (supporting information figure S8), where it also shows a similar trend of HER activity as that shown in figure 5a. All these measurements are carried out using graphite counter electrode and the HER activity of benchmarked Pt/C is also shown in figure 5a. The comparison of LSV polarization curves of different TMD structures with current density calculated using both electrochemical surface area and geometrical area are given in figure S9 (the x-axis scaling is same for all the curves). The enhanced HER activity (decrease in over potential and increase in current density) while moving from triangle to dendrite is evident in both MoS<sub>2</sub> and MoSe<sub>2</sub>, indicating the enhanced activity due to edges. Further, the activity of MoS<sub>2(1-x)</sub>Se<sub>2x</sub> triangle is much higher than that of MoS<sub>2</sub> triangle and dendrite, indicting the activity enhancement due to basal plane Se doping (which is proven via STEM analyses). Moreover, the activity of MoS<sub>2(1-x)</sub>Se<sub>2x</sub> triangle is further enhanced

when the morphology was changed from triangle to dendrite (figure S9F), indicating the role of edges in this basal plane Se doped MoS<sub>2</sub> dendrite.

Further, to evaluate the efficiency of the catalysts, Tafel slope values were extracted from the respective Tafel plots (figure 5c). The dendritic  $MoS_{2(1-x)}Se_{2x}$  structure showed the lowest Tafel slope (67 mV/dec) than those of triangle shape  $MoS_2$  (164 mV/dec),  $MoSe_2$  (148 mV/dec), and  $MoS_{2(1-x)}Se_{2x}$  (101 mV /dec) triangle flakes as well as dendritic structures of  $MoS_2$  (122 mV/dec) and  $MoSe_2$  (91 mV/dec) indicating its excellent kinetics behaviour among all. The Tafel slope of pristine  $MoS_2$  is observed to be 164 mV/dec which is in tune with that reported for CVD grown  $MoS_2$  (137) A significant enhancement in HER activity was observed for  $MoS_2$  dendrite where the Tafel slopes show a significant decrease of ~25%. Similar trend was observed in case of  $MoSe_2$  and  $MoS_{2(1-x)}Se_{2x}$  too, which is tabulated in table 1. This indicates the edge induced activity changes in TMDs. Secondly, while doing the Tafel analyses on doped TMDs, it is observed a further enhancement in activity because of doping (decrease in Tafel slope), as shown in table 2. This indicates the augmented HER kinetics due to the doping.

Table 1. Comparison of Tafel slopes of triangles and dendrites.

	Tafel slope (mV/dec)		% change
MoS <sub>2</sub>	Triangle	164	~ 25 %
	Dendrite	122	
MoSe <sub>2</sub>	Triangle	148	~ 38 %
	Dendrite	91	
$MoS_{2(1-x)}Se_{2x}$	Triangle	105	~ 36 %
	Dendrite	67	

Table 2. Comparison of Tafel slopes of undoped and doped (alloy) TMDs..

		Alloying	Tafel Slope change (%)
(1)			
Triangle	$MoS_2$	$MoS_{2(1-x)}Se_{2x}$	~ 36 %
0	MoSe <sub>2</sub>		~ 29 %
Dendrite	MoS <sub>2</sub>	$MoS_{2(1-x)}Se_{2x}$	~ 45 %
t	MoSe <sub>2</sub>		~ 26 %

To further confirm the observed HER activity trends of TMDs, electrochemical impedance spectroscopy (EIS) studies were performed in HER conditions at an overpotential of 250 mV, and the results are shown in figure 5d. The charge transfer resistance ( $R_{ct}$ ) values across the electrode-electrolyte interface observed after fitting Randles plot, for triangle shape MoS<sub>2</sub>, MoSe<sub>2</sub> and MoS<sub>2(1-x)</sub>Se<sub>2x</sub>, were found to be 10.1, 6.9 and 1.6 K $\Omega$ , respectively whereas that for the dendritic structures of MoS<sub>2</sub>, MoSe<sub>2</sub> and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> are found to be 4.1, 1.1 and 0.4 K $\Omega$ , respectively. The lowest  $R_{ct}$  value of dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> alloy among all the other TMD structures indicates the augmented electron transfer kinetics of this structure, in tune with the LSV studies. The EIS curves of MoS<sub>2</sub>, MoSe<sub>2</sub>, and MoS<sub>2(1-x)</sub>Se<sub>2x</sub> triangles and dendrites are compared in figure S10, and it is clear that dendritic structures have low  $R_{ct}$  values than the corresponding triangles while doped structures have the lowest  $R_{ct}$  values in comparison to the undoped structures. These studies show that the HER activity can be tuned both by morphology control (with larger number of edges) and basal plane doping.

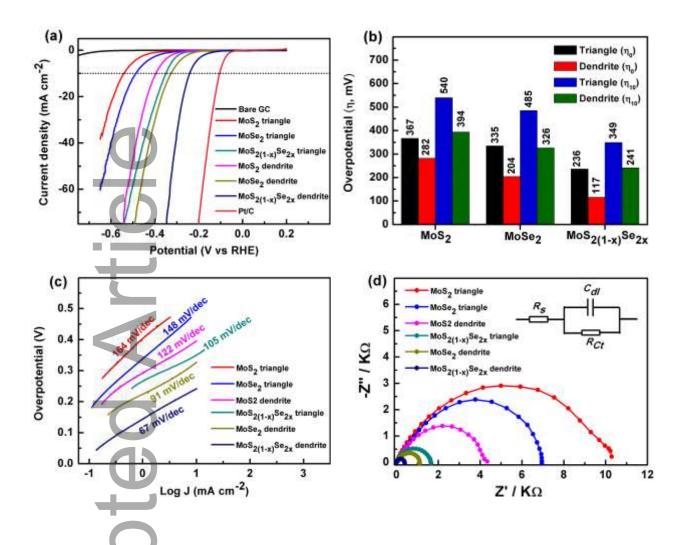


Figure 5: HER activities of triangular and dendritic shape binary and ternary TMDs on glassy carbon electrode (a) LSV curves measured in 0.5 M H<sub>2</sub>SO<sub>4</sub> at scan rate of 2 mV s<sup>-1</sup>. (b) Variation of onset overpotential and overpotential at 10 mA cm<sup>-2</sup>. (C) Tafel plots extracted from the corresponding polarization curves. (d) Nyquist plot carried out at overpotential of 250 mV (*vs.* RHE) with an inset of best fitted electrical equivalent (Randles) circuit. All the measurements were carried out with graphite rod as counter electrode.

Apart from the HER activity, the operational durability of an electrocatalyst is another vital aspect to realize its capability for the long term performance. To this purpose, the stability of an optimum ternary dendritic  $MoS_{2(1-x)}Se_{2x}$  electrocatalyst was tested by chronoamperometry

and the LSVs in 0.5 M H<sub>2</sub>SO<sub>4</sub>. Figure 6a shows the chronoamperometry response of the dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> electrocatalyst performed at an overpotential of 250 mV. The observed chronoamperometric current response of the dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> electrocatalyst is nearly constant even after more than 20,000 seconds of continuous operation indicating its high electrocatalytic stability for the long term HER performance in acidic condition. Furthermore, the electrochemical cycling stability of the dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> is also tested by measuring the polarization curves before and after 1000 cyclic voltammetry (CV) runs (figure 6b). After 1000 CVs, the dendritic MoS<sub>2(1-x)</sub>Se<sub>2x</sub> exhibits the reproducible HER polarization curve (only 14 mV difference in overpotential at 10 mA cm<sup>-2</sup>), indicating their high stability towards HER performance. A small increment in the current density was observed in the initial part of the chronoamperometry test (figure 6a), similar to that observed by Xu *et al.*, [22] where it can be due to the presence of sulfur vacancies and their possible aggregation during the initial electrochemical process.

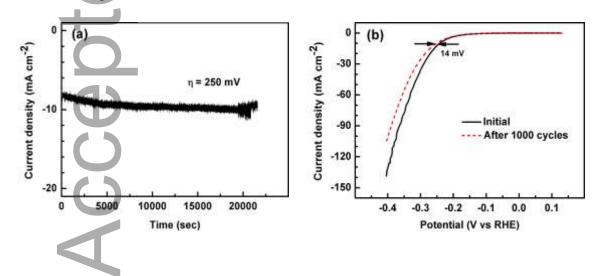


Figure 6: Stability test for dendritic  $MoS_{2(1-x)}Se_{2x}$  alloy catalyst (a) current-time response at an applied potential of -0.250 V (vs. RHE) and (b) LSV curves taken before and after 1000 voltammetry cycles at a scan rate of 2 mV s<sup>-1</sup>.

Different methods are adopted to enhance the inherent HER activity of 2H-MoS<sub>2</sub> (which is thermodynamically more stable than its structural polymorph 1T MoS<sub>2</sub>).<sup>[38]</sup> But, the 2H phase shows semiconducting nature while the 1T phase has a metallic nature. [38-39] Further, it is reported that 1T-MoS<sub>2</sub> has higher hydrophilicity which eventually directs the maximum interaction with the electrolyte. [40] Furthermore, unlike in 2H-MoS<sub>2</sub>, the basal plane of 1T-MoS<sub>2</sub> is also HER active. [38] Apart from the enhanced HER activity of 1T-MoS<sub>2</sub> atomically thin layers (inherent activity), recently nanostructured (enhanced active sites) 1T-MoS<sub>2</sub> flowers were shown for their HER properties. [38] The overpotential required for 10 mA/cm<sup>2</sup> current density ( $\eta_{10}$ ) was reported for this structures is found to be 252 mV, which is higher than that of present 2H- $MoS_{2(1-x)}Se_{2x}$  structures (241 mV). Further, the performance of  $MoS_{2(1-x)}Se_{2x}$  dendrite is compared with the other recent reports on CVD grown TMDs on different substrates, and the comparison is shown in table S2. This indicates that within the spectrum of thermodynamically stable TMD structures itself one can tune the HER activity of a parent crystal by controlling morphology and dopants, and this opens plethora of opportunities for engineered TMD structures for catalysis and other opto-electronic applications.

In conclusion, a novel method is proposed for the morphology controlled TMD structures of  $MoS_2$  and  $MoSe_2$ . Atomically thin several micrometer sized triangle and dendritic  $MoS_2$  and  $MoSe_2$  are developed by this generalized CVD protocol, and similar morphology controlled structures of their alloy, namely  $2H-MoS_{2(1-x)}Se_{2x}$ , are also synthesized, where the XPS analyses show that x is 0.34. Apart from the extensive HAADF-STEM based analyses, crystallinity and ultra-thinness (monolayer,  $\sim$ 0.83 nm) of the  $2H-MoS_{2(1-x)}Se_{2x}$  are also proven using X-ray diffraction (XRD, figure S11) and atomic force microscopy (AFM, figure S12) studies. The doping and alloying of Se in  $MoS_2$  are proven via different techniques and the possibility of

tuning the bandgap by alloying is demonstrated. HER properties of these structures in acidic medium are studied and the role of morphology in the HER performance is unravelled. The HER activities (in terms of both overpotential and Tafel slope) of triangle shaped 2H-MoS<sub>2</sub> and MoSe<sub>2</sub> are found to be augmented by changing their morphology to dendritic, and further enhancement in the HER activity can be achieved by doping 2H-MoS<sub>2</sub> with Se using the same synthesis set up. The 'best' HER activity is observed for the 2H-MoS<sub>2(1-x)</sub>Se<sub>2x</sub> dendritic structure where it is found to be even better than the recently reported nanostructured (flower shaped) inherently more active 1T-MoS<sub>2</sub> structures. This synthesis strategy opens intriguing opportunities for the growth of morphology controlled large area complex structures of TMDs *via* non-equilibrium CVD, with the potential of yielding structures with enhanced catalytic and tuneable optoelectronic properties, as demonstrated here.

#### **Experimental methods:**

#### **Materials:**

Molybdenum (VI) tri oxide (MoO<sub>3</sub>, 99.95%) from SRL, sulfur (S, 99.98%) and selenium (Se, 99.99%) were purchased from sigma Aldrich.

#### **Characterization:**

Crystallinity and quality of pristine and dendritic  $MoS_2$ ,  $MoSe_2$  and  $MoS_{2(1-x)}Se_{2x}$  were studied by XRD (Cu K $\alpha$ radiation,  $\succcurlyeq$ 1.5418 Å ) and micro-Raman spectroscopy (using Invia Raman spectrometer equipped with 532 and 633 nm laser excitation wavelengths) studies. The samples were focused with 100x objective lens at a laser power  $\sim$  1 mw. The SEM images were taken using JEOL JSM7200F SEM. XPS was performed to estimate chemical composition and

atomic percentage of Se in  $MoS_{2(1-x)}Se_{2x}$  alloy using KratosS5 Axis Ultra DLD spectrometer using the Al K $\alpha$ X -ray source (1486.6 eV). The HAADF-STEM studies are conducted using an aberration-corrected STEM (JEOL 2100F, 80 kV).

All the electrochemical measurements were performed in a three electrode set up using a Bio Logic SP-300 electrochemical workstation at room temperature. Transferred TMDs sample on GC electrodes were served as working electrodes, graphite rod as a counter electrode (Pt is also used to verify the data, but data is not shown), and Ag/AgCl was used as a reference electrode. The hydrogen evolution performance was carried out in 0.5 M  $H_2SO_4$  solution using LSV technique at a scan rate of 2 mV s<sup>-1</sup>. Prior to electrochemical measurements, the electrolyte was degassed by purging pure  $N_2$  for 30 minutes. The EIS measurements were performed in the frequency ranged of 0.1 Hz to 100 KHz at an overpotential of 250 mV. All the potentials presented in this work were referenced to the reversible hydrogen electrode (RHE) using equation  $E_{RHe} = E_{Ag/AgCl} + 0.197 \text{ V (for 0.5 M H}_2SO_4)$ .

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## References:

- [1] Z. Lin, B. R. Carvalho, E. Kahn, R. Lv, R. Rao, H. Terrones, M. A. Pimenta, M. Terrones, *2D Mater.* **2016**, *3*, 022002.
- [2] M. G. Stanford, P. D. Rack, D. Jariwala, npj 2D Materials and Applications 2018, 2, 20.
- [3] T.-T. Jia, M.-M. Zheng, X.-Y. Fan, Y. Su, S.-J. Li, H.-Y. Liu, G. Chen, Y. Kawazoe, *Sci. Rep.* **2016**, *6*, 18869.
- [4] Y. L. Huang, Y. Chen, W. Zhang, S. Y. Quek, C.-H. Chen, L.-J. Li, W.-T. Hsu, W.-H. Chang, Y. J. Zheng, W. Chen, A. T. S. Wee, *Nat. Commun.* **2015**, *6*, 6298.
- [5] A. Kormányos, V. Zólyomi, N. D. Drummond, G. Burkard, *Phys. Rev. X* **2014**, *4*, 011034.
- [6] D. Kotekar-Patil, J. Deng, S. L. Wong, C. S. Lau, K. E. J. Goh, *Appl. Phys. Lett.* **2019**, *114*, 013508.
- [7] J. Kim, W. S. Yun, J. D. Lee, J. Phys. Chem. C 2015, 119, 13901.
- [8] J. Ma, X. Li, L. Gan, S. Zhang, Y. Cao, Z. Nie, X. Wang, D. Ma, L. He, J. Nie, C. Xiong, R. Dou, *2D Mater* . **2018**, *5*, 031015.
- [9] G. Zhang, J. Wang, Z. Wu, R. Shi, W. Ouyang, A. Amini, B. N. Chandrashekar, N. Wang, C. Cheng, *ACS Appl. Mater. Interfaces* **2017**, *9*, 763.
- [10] B. Wu, D. Geng, Z. Xu, Y. Guo, L. Huang, Y. Xue, J. Chen, G. Yu, Y. Liu, *Npg Asia Materials* **2013**, *5*, e36.
- [11] W. Fei, J. Yin, X. Liu, W. Guo, Mater. Lett. 2013, 110, 225.
- [12] S. Wang, Y. Rong, Y. Fan, M. Pacios, H. Bhaskaran, K. He, J. H. Warner, *Chem. Mater.* **2014**, *26*, 6371.
- [13] Q. Feng, N. Mao, J. Wu, H. Xu, C. Wang, J. Zhang, L. Xie, ACS Nano 2015, 9, 7450.
- [14] H. Yu, M. Liao, W. Zhao, G. Liu, X. J. Zhou, Z. Wei, X. Xu, K. Liu, Z. Hu, K. Deng, S. Zhou, J.-A. Shi, L. Gu, C. Shen, T. Zhang, L. Du, L. Xie, J. Zhu, W. Chen, R. Yang, D. Shi, G. Zhang, ACS Nano 2017, 11, 12001.
- [15] S. Susarla, A. Kutana, J. A. Hachtel, V. Kochat, A. Apte, R. Vajtai, J. C. Idrobo, B. I. Yakobson, C. S. Tiwary, P. M. Ajayan, *Adv. Mater* .2017, 29, 1702457.
- [16] J. N. Coleman, M. Lotya, A. O'Neill, S. D. Bergin, P. J. King, U. Khan, K. Young, A. Gaucher, S. De, R. J. Smith, I. V. Shvets, S. K. Arora, G. Stanton, H.-Y. Kim, K. Lee, G. T. Kim, G. S. Duesberg, T. Hallam, J. J. Boland, J. J. Wang, J. F. Donegan, J. C. Grunlan, G. Moriarty, A. Shmeliov, R. J. Nicholls, J. M. Perkins, E. M. Grieveson, K. Theuwissen, D. W. McComb, P. D. Nellist, V. Nicolosi, *Science* 2011, 331, 568.
- [17] L. Yuan, J. Ge, X. Peng, Q. Zhang, Z. Wu, Y. Jian, X. Xiong, H. Yin, J. Han, *AIP Adv.* **2016**, *6*, 125201.
- [18] a) R. K. Biroju, D. Das, R. Sharma, S. Pal, L. P. L. Mawlong, K. Bhorkar, P. K. Giri, A. K. Singh, T. N. Narayanan, *ACS Energy Lett.* **2017**, *2*, 1355; b) P. K. Rastogi, S. Sarkar, D. Mandler, *Appl. Mater. Today* **2017**, *8*, 44.

- [19] R. Kronberg, M. Hakala, N. Holmberg, K. Laasonen, *Phys. Chem. Chem. Phys.* **2017**, 19, 16231.
- [20] T. H. M. Lau, X. Lu, J. Kulhavý, S. Wu, L. Lu, T.-S. Wu, R. Kato, J. S. Foord, Y.-L. Soo, K. Suenaga, S. C. E. Tsang, *Chem Sci.* **2018**, *9*, 4769.
- [21] S. Pal, K. K. Tadi, P. M. Sudeep, S. Radhakrishnan, T. N. Narayanan, *Mater. Chem. Front.***2017**, *1*, 319.
- [22] W. Xu, S. Li, S. Zhou, J. K. Lee, S. Wang, S. G. Sarwat, X. Wang, H. Bhaskaran, M. Pasta, J. H. Warner, *ACS Appl. Mater. Interfaces* **2018**, *10*, 4630.
- [23] Y. Zhang, Q. Ji, G.-F. Han, J. Ju, J. Shi, D. Ma, J. Sun, Y. Zhang, M. Li, X.-Y. Lang, Y. Zhang, Z. Liu, *ACS Nano* **2014**, *8*, 8617.
- [24] L. Yang, Q. Fu, W. Wang, J. Huang, J. Zhang, B. Xiang, *Nanoscale* **2015**, 7, 10490.
- [25] R. Saito, Y. Tatsumi, S. Huang, X. Ling, M. S. Dresselhaus, *J. Phys. Condens. Matter* **2016**, *28*, 353002.
- [26] a) K. F. Mak, C. Lee, J. Hone, J. Shan, T. F. Heinz, *Phys. Rev. Lett.* 2010, 105, 136805;
  b) S. Tongay, J. Zhou, C. Ataca, K. Lo, T. S. Matthews, J. Li, J. C. Grossman, J. Wu, *Nano Lett* 2012, 12, 5576.
- [27] S.-H. Su, Y.-T. Hsu, Y.-H. Chang, M.-H. Chiu, C.-L. Hsu, W.-T. Hsu, W.-H. Chang, J.-H. He, L.-J. Li, *Small* **2014**, *10*, 2589.
- [28] C. Ruppert, O. B. Aslan, T. F. Heinz, *Nano Letters* **2014**, *14*, 6231.
- [29] G. Eda, H. Yamaguchi, D. Voiry, T. Fujita, M. Chen, M. Chhowalla, *Nano Lett Nano Lett* **2011**, *11*, 5111.
- [30] P. Tonndorf, R. Schmidt, P. Böttger, X. Zhang, J. Börner, A. Liebig, M. Albrecht, C. Kloc, O. Gordan, D. R. T. Zahn, S. Michaelis de Vasconcellos, R. Bratschitsch, *Opt. Express* **2013**, *21*, 4908.
- [31] K. K. Paul, N. Sreekanth, R. K. Biroju, T. N. Narayanan, P. K. Giri, *Sol. Energ. Mat. Sol. C.***2018**, *185*, 364.
- [32] v. d. Heide, in *X- Ray Photoelectron Spectroscopy*, DOI: 10.1002/9781118162897.fmatter (Ed: P. v. d. Heide) **2011**.
- [33] Y. Gong, Z. Liu, A. R. Lupini, G. Shi, J. Lin, S. Najmaei, Z. Lin, A. L. Elías, A. Berkdemir, G. You, H. Terrones, M. Terrones, R. Vajtai, S. T. Pantelides, S. J. Pennycook, J. Lou, W. Zhou, P. M. Ajayan, *Nano Lett* **2014**, *14*, 442.
- [34] a) O. L. Krivanek, M. F. Chisholm, V. Nicolosi, T. J. Pennycook, G. J. Corbin, N. Dellby, M. F. Murfitt, C. S. Own, Z. S. Szilagyi, M. P. Oxley, S. T. Pantelides, S. J. Pennycook, *Nature* 2010, 464, 571; b) W. Zhou, M. P. Oxley, A. R. Lupini, O. L. Krivanek, S. J. Pennycook, J.-C. Idrobo, *Microscopy and Microanalysis* 2012, 18, 1342.
- [35] S. Anantharaj, S. R. Ede, K. Karthick, S. Sam Sankar, K. Sangeetha, P. E. Karthik, S. Kundu, *Energy Environ. Sci.*, **2018**, *11*, 744.
- [36] Q. Gong, L. Cheng, C. Liu, M. Zhang, Q. Feng, H. Ye, M. Zeng, L. Xie, Z. Liu, Y. Li, *ACS Catal.* **2015**, *5*, 2213.
- [37] Z. Wang, Q. Li, H. Xu, C. Dahl-Petersen, Q. Yang, D. Cheng, D. Cao, F. Besenbacher, J. Lauritsen, S. Helveg, M. Dong, *Nano energy*, **2018**, *49*, 634-643
- [38] S. Venkateshwaran, S. M. Senthil Kumar, ACS Sustain. Chem. Eng. 2019, 7, 2008.
- [39] L. Chacko, P. K. Rastogi, P. M. Aneesh, J. Electrochem. Soc. 2019, 166, H263.

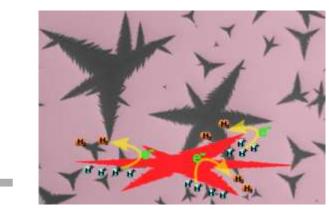
[40] X. Geng, W. Sun, W. Wu, B. Chen, A. Al-Hilo, M. Benamara, H. Zhu, F. Watanabe, J. Cui, T.-p. Chen, *Nat Commun.* **2016**, *7*, 10672.

**Keyword:** CVD, Transition Metal Dichalcogenide, Dendritic Structures, Hydrogen Evolution, Atomically Thin Alloys.

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On the Synthesis of Morphology Controlled Transition Metal Dichalcogenides via CVD for

#### **Electrochemical Hydrogen Generation**



A generalized CVD based synthesis strategy for morphology controlled TMDs and their alloys, and the enhanced hydrogen evolution activities of dendritic TMD alloys:

Optical images of dendritic  $2H\text{-MoS}_{2(1-x)}Se_{2x}$  structures are shown in the figure with their photoluminescence properties indicated by the red colour. Augmented electro-catalytic hydrogen generation is also shown in the schematic.