

# First narrow-band search for continuous gravitational waves from known pulsars in advanced detector data

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# First narrow-band search for continuous gravitational waves from known pulsars in advanced detector data

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Spinning neutron stars asymmetric with respect to their rotation axis are potential sources of continuous gravitational waves for ground-based interferometric detectors. In the case of known pulsars a fully coherent search, based on matched filtering, which uses the position and rotational parameters obtained from electromagnetic observations, can be carried out. Matched filtering maximizes the signal-to-noise (SNR) ratio, but a large sensitivity loss is expected in case of even a very small mismatch between the assumed and the true signal parameters. For this reason, *narrow-band* analysis methods have been developed, allowing a fully coherent search for gravitational waves from known pulsars over a fraction of a hertz and several spin-down values. In this paper we describe a narrow-band search of 11 pulsars using data from Advanced LIGO's first observing run. Although we have found several initial outliers, further studies show no significant evidence for the presence of a gravitational wave signal. Finally, we have placed upper limits on the signal strain amplitude lower than the spin-down limit for 5 of the 11 targets over the bands searched; in the case of J1813-1749 the spin-down limit has been beaten for the first time. For an additional 3 targets, the median upper limit across the search bands is below the spin-down limit. This is the most sensitive narrow-band search for continuous gravitational waves carried out so far.

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## I. INTRODUCTION

On September 14, 2015, the gravitational wave (GW) signal emitted by a binary black hole merger was detected by the LIGO interferometers (IFOs) [1] followed on December 26, 2015, by the detection of a second event again associated to a binary black hole merger [2], thus opening the era of gravitational waves astronomy. More recently, the detection of a third binary black hole merger on January 4, 2017, was announced [3]. Binary black hole mergers, however, are not the only detectable sources of GW. Among the potential sources of GW there are also spinning neutron stars (NS) asymmetric with respect to their rotation axis. These sources are expected to emit nearly monochromatic continuous waves (CW), with a frequency at a given fixed ratio with respect to the star's rotational frequency, e.g. two times the rotational frequency for an asymmetric NS rotating around one of its principal axes of inertia. Different flavors of CW searches exist, depending on the degree of knowledge on the source parameters. *Targeted* searches assume source position and rotational parameters to be known with high accuracy, while *all-sky* searches aim at neutron stars with no observed electromagnetic counterpart. Various intermediate searches have also been developed. Among these, *narrow-band* searches are an extension of targeted searches for which the position of the source is accurately known but the rotational parameters are slightly uncertain. Narrow-band searches

allow for a possible small mismatch between the GW rotational parameters and those inferred from electromagnetic observations. This can be crucial if, for instance, the CW signal is emitted by a freely precessing neutron star [4], or in the case no updated ephemeris is available for a given pulsar. In both cases a targeted search could assume wrong rotational parameters, resulting in a significant sensitivity loss. In this paper we present the results of a fully coherent, narrow-band search for 11 known pulsars using data from the first observation run (O1) of the Advanced LIGO detectors [5]. The paper is organized as follows. In Sec. II we briefly summarize the main concepts of the analysis. Section III is dedicated to an outline of the analysis method. Section IV describes the selected pulsars. In Sec. V we discuss the analysis results, while the reader can refer to the Appendix for some technical details on the computation of upper limits. Finally, Sec. VI is dedicated to the conclusions and future prospects.

## II. BACKGROUND

The GW signal emitted by an asymmetric spinning NS can be written, following the formalism first introduced in [6], as the real part of

$$h(t) = H_0(H^+A_+(t) + H^\times A_\times(t))e^{2\pi i f_{\text{gw}}(t)t + i\phi_0} \quad (1)$$

where  $f_{\text{gw}}(t)$  is the GW frequency,  $\phi_0$  an initial phase. The polarization amplitudes  $H^+$ ,  $H^\times$  are given by

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\*Full author list given at the end of the article.

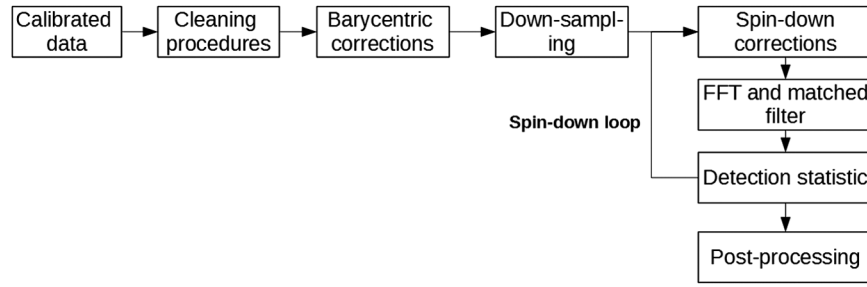


FIG. 1. Simplified flowchart of the narrow-band search pipeline for CW. The method relies on the use of FFTs to simultaneously compute the detection statistic, for each given spin-down value, over the full explored frequency range. See [13] for more details on the method.

$$H^+ = \frac{\cos(2\psi) - i\eta \sin(2\psi)}{\sqrt{1 + \eta^2}},$$

$$H^\times = \frac{\sin(2\psi) - i\eta \cos(2\psi)}{\sqrt{1 + \eta^2}},$$

$\eta$  being the ratio of the polarization ellipse semiminor to semimajor axis and  $\psi$  the polarization angle, defined as the direction of the major axis with respect to the celestial parallel of the source (measured counterclockwise). The detector *sidereal response* to the GW polarizations is encoded in the functions  $A_+(t), A_\times(t)$ . It can be shown that the waveform defined by Eq. (1) is equivalent to the GW signal expressed in the more standard formalism of [7], given by the following relations:

$$\eta = -\frac{2 \cos \iota}{1 + \cos^2 \iota}, \quad (2)$$

where  $\iota$  is the angle between the line of sight and the star rotation axis, and

$$H_0 = h_0 \sqrt{\frac{1 + 6\cos^2 \iota + \cos^4 \iota}{4}} \quad (3)$$

with

$$h_0 = \frac{1}{d} \frac{4\pi^2 G}{c^4} I_{zz} f_{\text{gw}}^2 \epsilon, \quad (4)$$

where  $d, I_{zz}$  and  $\epsilon$  are respectively the star's distance, its moment of inertia with respect to the rotation axis and the *ellipticity*, which measures the star's degree of asymmetry. The signal at the detector is not monochromatic; i.e. the frequency  $f_{\text{gw}}(t)$  in Eq. (1) is a function of time. In fact the signal is modulated by several effects, such as the *Römer delay* due to the detector motion and the source's intrinsic spin-down due to the rotational energy loss from the source. In order to recover all the signal-to-noise ratio all these effects must be properly taken into account. If we have a measure of the pulsar rotational frequency  $f_{\text{rot}}$ , frequency derivative  $\dot{f}_{\text{rot}}$  and distance  $d$ , the GW signal amplitude can be constrained, assuming that all the rotational energy is lost via gravitational radiation. This strict upper limit, called the “spin-down limit,” is given by [8]

$$h_{sd} = 8.06 \times 10^{-19} I_{38}^{1/2} \left[ \frac{1 \text{ kpc}}{d} \right] \left[ \frac{\dot{f}_{\text{rot}}}{\text{Hz/s}} \right]^{1/2} \left[ \frac{\text{Hz}}{f_{\text{rot}}} \right]^{1/2} \quad (5)$$

where  $I_{38}$  is the star's moment of inertia in units of  $10^{38} \text{ kgm}^2$ . The corresponding spin-down limit on the star's equatorial fiducial ellipticity can be easily obtained from Eq. (4):

$$\epsilon_{sd} = 0.237 I_{38}^{-1} \left[ \frac{h_{sd}}{10^{-24}} \right] \left[ \frac{\text{Hz}}{f_{\text{rot}}} \right]^2 \left[ \frac{d}{1 \text{ kpc}} \right]. \quad (6)$$

Even in the absence of a detection, establishing an amplitude upper limit below the spin-down limit for a given source is an important milestone, as it allows us to put a nontrivial constraint on the fraction of rotational energy lost through GWs.

### III. THE ANALYSIS

The results discussed in this paper have been obtained by searching for CW signals from 11 known pulsars using data from the O1 run from the Advanced LIGO detectors [Hanford (LIGO H) and Livingston (LIGO L) jointly]. The run started on September 12, 2015, at 01:25:03 UTC and 18:29:03 UTC, respectively, and finished on January 19, 2016, at 17:07:59. LIGO H had a duty cycle of  $\sim 60\%$  and LIGO L had a duty cycle of  $\sim 51\%$ , which correspond respectively to 72 and 62 days of science data available for the analysis. In this paper we have used an initial calibration of the data [9]. In order to perform a joint search between the two detectors a common period from September 13, 2015, to January 12, 2016,<sup>1</sup> with a total observation time of about  $T_{\text{obs}} \approx 121$  days is selected. The natural frequency and spin-down grid spacings of the search are  $\delta f = 1/T_{\text{obs}} \approx 9.5 \times 10^{-8} \text{ Hz}$  and  $\delta \dot{f} = 1/T_{\text{obs}}^2 \approx 4.57 \times 10^{-15} \text{ Hz/s}$ . A follow-up analysis based on the LIGO's second observation run (O2) has been carried out. For this data set we have analyzed data from December 16, 2016, to May 8, 2017; more details will be given in Appendix C. The analysis pipeline consists of several steps, schematically depicted in Fig. 1, which we summarize here. The starting point is a

<sup>1</sup>An exception is pulsar J0205 + 6449; see later.

TABLE I. This table reports the explored range for the rotational parameters of each pulsar. The columns are the central frequency of the search ( $f_0$ ), explored frequency band ( $\Delta f$ ), central spin-down value of the search ( $\dot{f}_0$ ), explored spin-down band ( $\Delta \dot{f}_0$ ), the number of frequency bins explored ( $n_f$ ), and the number of spin-down values explored ( $n_{\dot{f}}$ ). All the rotational parameters are scaled at the common reference time on September 12, 2015.

Name	$f_0$ (Hz)	$\Delta f$ (Hz)	$\dot{f}_0$ (Hz/s)	$\Delta \dot{f}$ (Hz/s)	$n_f$	$n_{\dot{f}}$
J0205 + 6449	30.4095820	0.03	$-8.9586 \times 10^{-11}$	$1.75 \times 10^{-13}$	$2.5 \times 10^6$	19
J0534 + 2200 (Crab)	59.32365204	0.10	$-7.3883 \times 10^{-10}$	$1.48 \times 10^{-12}$	$18.5 \times 10^6$	161
J0835-4510 (Vela)	22.3740981	0.03	$-3.1191 \times 10^{-11}$	$6.43 \times 10^{-14}$	$2.5 \times 10^6$	7
J1400-6326	64.1253722	0.07	$-8.0017 \times 10^{-11}$	$1.75 \times 10^{-13}$	$6.5 \times 10^6$	19
J1813-1246	41.6010333	0.04	$-1.2866 \times 10^{-11}$	$6.43 \times 10^{-14}$	$3.4 \times 10^6$	7
J1813-1749	44.7128464	0.05	$-1.5000 \times 10^{-10}$	$3.03 \times 10^{-13}$	$2.5 \times 10^6$	33
J1833-1034	32.2940958	0.04	$-1.0543 \times 10^{-10}$	$2.11 \times 10^{-13}$	$3.4 \times 10^6$	23
J1952 + 3252	50.5882336	0.05	$-7.4797 \times 10^{-12}$	$6.43 \times 10^{-14}$	$4.3 \times 10^6$	7
J2022 + 3842	41.1600845	0.04	$-7.2969 \times 10^{-11}$	$1.60 \times 10^{-13}$	$3.4 \times 10^6$	17
J2043 + 2740	20.8048628	0.05	$-3.4390 \times 10^{-11}$	$6.43 \times 10^{-14}$	$4.3 \times 10^6$	7
J2229 + 6114	38.7153156	0.06	$-5.8681 \times 10^{-11}$	$1.19 \times 10^{-13}$	$5.1 \times 10^6$	13

collection of FFTs obtained from several interlaced data chunks [the short FFT Database (SFDB)] built from calibrated detector data chunks of duration 1024 seconds [10]. At this stage, a first cleaning procedure is applied to the data in order to remove large, short-duration disturbances, that could reduce the search sensitivity. A frequency band is then extracted from the SFDBs covering typically a range larger (of the order of a factor of 2) than the frequency region analyzed in the narrow-band search. The actual search frequency and spin-down bands,  $\Delta f$  and  $\Delta \dot{f}$ , around the reference values,  $f_0$  and  $\dot{f}_0$ , have been chosen according to the following relations [11]:

$$\Delta f = 2f_0\delta \quad (7)$$

$$\Delta \dot{f} = 2\dot{f}_0\delta, \quad (8)$$

$\delta$  being a factor parametrizing a possible discrepancy between the GW rotation parameters and those derived from electromagnetic observations. Previous narrow-band searches used values of  $\delta$  of the order  $\sim \mathcal{O}(10^{-4})$ , motivated partly by astrophysical considerations [4], and partly by computational limitations [12]. Here we exploit the high computation efficiency of our pipeline to enlarge the search somewhat, depending on the pulsar, to a range between  $\delta \sim 10^{-4}$  and  $10^{-3}$ . The frequency and spin-down ranges explored in this analysis are listed in Table I.

The narrow-band search is performed using a pipeline based on the *five-vector method* [12] and, in particular, its latest implementation, fully described in [13], to which the reader is referred for more details. The basic idea is that of exploring a range of frequency and spin-down values by properly applying barycentric and spin-down corrections to the data in such a way that a signal would become monochromatic apart from the sidereal modulation. While a single barycentric correction applied in the time domain holds for all the explored frequency bands, several

spin-down corrections, one for each point in the spin-down grid, are needed. A detection statistic (DS) is then computed for each point of the explored parameter space. By using the FFT algorithm for each given spin-down value it is possible to compute the statistic simultaneously over the whole range of frequencies; this process is done for each detector, and then data are combined. The frequency/spin-down plane is then divided into frequency subbands ( $10^{-4}$  Hz) and, for each of them, the local maximum, over

TABLE II. Distance and spin-down limit on the GW amplitude and ellipticity for the 11 selected pulsars. Distance and spin-down limit uncertainties refer to the  $1\sigma$  confidence level.

Name	Distance (kpc)	$h_{sd} \times 10^{-25}$	$e_{sd} \times 10^{-4}$
J0205 + 6449 <sup>a</sup>	$2.0 \pm 0.3^b$	$6.9 \pm 1.1$	14
J0534 + 2200 (Crab)	$2.0 \pm 0.5^c$	$14 \pm 3.5$	7.6
J0835-4510 (Vela)	$0.28 \pm 0.02^c$	$34 \pm 2.4$	18
J1400-6326	$10 \pm 3^d$	$0.90 \pm 0.27$	2.1
J1813-1246	$>2.5^e$	$<1.8$	$<2.4$
J1813-1749	$4.8 \pm 0.3^f$	$3.0 \pm 0.2$	7.0
J1833-1034	$4.8 \pm 0.4^g$	$3.1 \pm 0.3$	13
J1952 + 3252	$3.0 \pm 0.5^h$	$1.0 \pm 0.2$	1.1
J2022 + 3842	$10 \pm 2^i$	$1.0 \pm 0.3$	6.0
J2043 + 2740	$1.5 \pm 0.6^j$	$6.9 \pm 2.8$	23
J2229 + 6114	$3.0 \pm 2^c$	$3.4 \pm 2.2$	6.2

<sup>a</sup>This pulsar had a glitch on November 11, 2015.

<sup>b</sup>Distance from neutral hydrogen absorption of pulsar wind nebula 3C 58 [14].

<sup>c</sup>Distance taken from independent measures reported in the ATNF catalog; see text for references.

<sup>d</sup>Distance from dispersion measures [15].

<sup>e</sup>Lower limit of [16].

<sup>f</sup>Distance from Chandra and XMM-Newton from [17].

<sup>g</sup>Distance from the Parkes telescope [18].

<sup>h</sup>Distance from the kinematic distance of the associated supernova remnant [19].

<sup>i</sup>Distance of the hosting supernova remnant [20]. In some papers a distance value of  $\sim 10$  kpc is considered [21].

<sup>j</sup>Distances taken from v1.56 of the ATNF Pulsar Catalog [22].

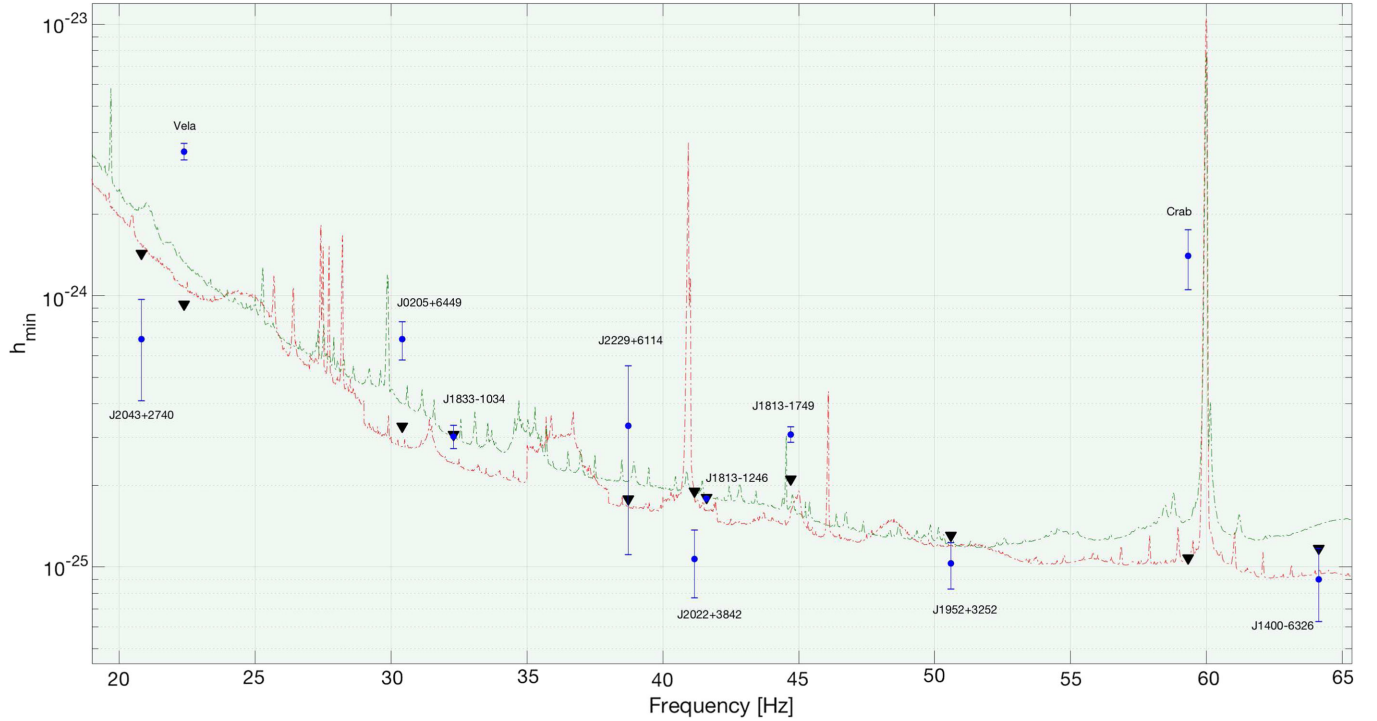


FIG. 2. Blue points: Value of the theoretical spin-down limit computed for the 11 known pulsars in our analysis, corresponding to Table II; error bars correspond to the  $1\sigma$  confidence level. Black triangles: Median over the analyzed frequency band of the upper limits on the GW amplitude, corresponding to Table VI. Red dashed line: Estimated sensitivity at 95% confidence level of a narrow-band search using data from LIGO H. Green dashed line: Estimated sensitivity at 95% confidence level of a narrow-band search using data from LIGO L.

the spin-down grid, of the DS is selected as a *candidate*. The initial *outliers* are identified among the candidates using a threshold nominally corresponding to 1% (taking into account the number of trials [12]) on the p-value of the DS's noise-only distribution<sup>2</sup> and are subject to a follow-up stage in order to understand their nature. The follow-up procedure consists of the following steps: check if the outlier is close to known instrumental noise lines; compute the signal amplitude and check if it is constant throughout the run; compute the time evolution of the SNR (which we expect to increase as the square root of the observation time for stationary noise); and compute the *five-vector coherence*, which is an indicator measuring the degree of consistency between the data and the estimated waveform [6]. For each target, if no outlier is confirmed by the follow-up we set an upper limit on the GW amplitude and NS ellipticity; see Appendix A for more details.

#### IV. SELECTED TARGETS

We have selected pulsars whose spin-down limit could possibly be beaten, or at least approached, based on the average sensitivity of O1 data; see Fig. 2. Pulsar distances

<sup>2</sup>The noise-only distribution is computed from the values of the DS excluded in each frequency subband when selecting the local maxima and then an extrapolation of the long tail of the done.

and spin-down limits are listed in Table II. As distance estimations for the pulsars we have used the best fit value and relative uncertainties given by each independent measure; see pulsars list below and Table II for more details. The uncertainty on the spin-down limit in Table II can be computed using the relation for the variance propagation.<sup>3</sup> For two of these pulsars (Crab and Vela) the spin-down limit has been already beaten in a past narrow-band search using Virgo VSR4 data [11]. The other targets are analyzed in a narrow-band search for the first time. The timing measures for the 11 pulsars were provided by the 76-meter Lovell telescope and the 42-foot radio telescopes at Jodrell Bank (UK), the 26-meter telescope at Hartebeesthoek (South Africa), the 64-meter Parkes radio telescope (Australia) and the Fermi Large Area Telescope (LAT) which is a space satellite. For seven of these pulsars (Crab, Vela, J0205 + 6449, J1813-1246, J1952 + 3252, J2043 + 2740 and J2229 + 6114) updated ephemerides covering the O1 period were available and a targeted search was done in a recent work [7] beating the

<sup>3</sup>If variable  $Y$  is defined from  $x_i$  random variables with variance  $\sigma_{x_i}^2$ , then the variance  $\sigma_Y^2$  can be estimated as

$$\sigma_Y^2 = \sum_i \left( \frac{\partial Y}{\partial x_i} \right)^2 \sigma_{x_i}^2.$$

spin-down limit for all of them, while for the remaining four pulsars we have used older measures extrapolating the rotational parameters to the O1 epoch. A list of the analyzed pulsars follows.

J0205 + 6449: Ephemerides obtained from Jodrell Bank. This pulsar had a glitch on November 11, 2015, which can affect the CW search [23]. For this reason we have performed the narrow-band search only on data before the glitch as done in [7]. The distances are set according to [14].

J0534 + 2200 (Crab): One of the high value targets for CW searches [7] due to its large spin-down value. For this pulsar it was possible to beat the spin-down limit in a narrow-band search using Virgo VSR4 data [11]. Ephemerides have been obtained from the Jodrell Bank telescope.<sup>4</sup> The nominal distance for the Crab pulsar and its nebula is quoted in the literature as  $2.0 \pm 0.5$  kpc [24]; we therefore assume uncertainty corresponding to the  $1\sigma$  confidence level.

J0835-4510 (Vela): Like the Crab pulsar, Vela is one of the traditional targets for CW searches. Although it spins at a relatively low frequency (compared to the others), it is very close to the Earth ( $d \approx 0.28$  kpc), thus making it a potentially interesting source. Ephemerides were obtained from the Hartebeesthoek Radio Astronomy Observatory in South Africa.<sup>5</sup> The distance and its uncertainty are taken according to [25].

J1400-6326: First discovered as an INTEGRAL source and then identified as a pulsar by the Rossi X-ray Timing Explorer (RXTE). This NS is located in the galactic supernova remnant G310.6-1.6 and it is supposed to be quite young; the distance and its uncertainty correspond to the  $1\sigma$  confidence level [15].

J1813-1246: Ephemerides covering the O1 time span have been provided by the Fermi-LAT Collaboration [7]. Only a lower upper limit is present on the distance.

J1813-1749: Located in one of the brightest and most compact TeV sources discovered in the HESS Galactic Plane Survey, HESS J1813-178. It is a young energetic pulsar that is responsible for the extended x rays, and probably the TeV radiation as well. Timing was obtained from Chandra and XMM Newton data [17]; the pulsar's distance and uncertainty are taken from [26] and correspond to the  $1\sigma$  confidence level.

J1833-1034: Located in the supernova remnant G21.5-0.9. This source has been known for a long time as one of the Crab-like remnants. The evidence for a pulsar was found by analyzing Chandra data; the distance and its uncertainty are set according to [18] and correspond to the  $1\sigma$  confidence level.

<sup>4</sup><http://www.jb.man.ac.uk/pulsar/crab.html>.

<sup>5</sup><http://www.hartrao.ac.za/>

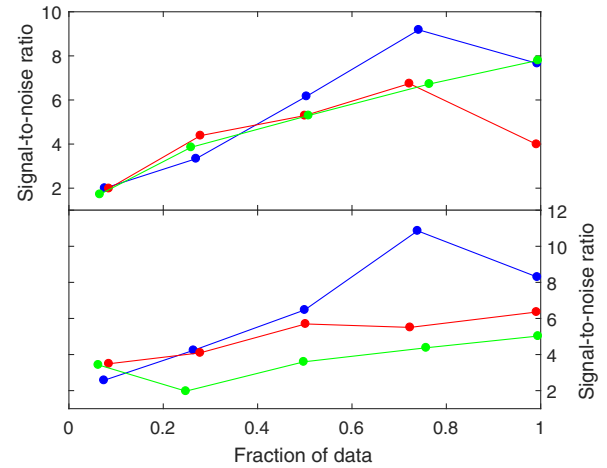


FIG. 3. Top panel: SNR computed with respect to the fraction of data for the J1833-1034 outlier in the Hanford (red line), Livingston (green) and joint (blue) analysis respectively. Bottom panel: SNR computed with respect to the fraction of data for the Vela outlier in the Hanford (red line), Livingston (green) and joint (blue) analysis respectively.

J1952 + 3252: Ephemerides have been obtained from Jodrell Bank [7]. Distance and uncertainty are taken from kinematic measures of [19].

J2022 + 3842: It is a young energetic pulsar that was discovered in Chandra observations of the radio supernova remnant SNR G76.9 + 1.0. Distance and uncertainty are set according to [21].

J2043 + 2740: Ephemerides obtained from the Fermi-LAT Collaboration [7]. The distance is estimated using the dispersion measure by [22] and using the model from [27]. The uncertainty on distance is set according to the model and correspond to the  $1\sigma$  confidence level.

J2229 + 6114: Ephemerides obtained from Jodrell Bank [7]. Distance and uncertainty are estimated by [28] using the model [29].

## V. RESULTS

In this section we discuss the results of the analysis. First, in Sec. VA we briefly describe the initial outliers, for most of which the follow-up described in Sec. III has been enough to exclude a GW origin. Two outliers, belonging respectively to the Vela and J1833-1034 pulsars, needed a deeper study. The studies discussed in detail in the next section disfavor the signal hypothesis and seem to suggest these outliers as marginal noise events. Nevertheless the outliers showed some promising features and for this reason a follow-up using O2 data has been carried out and described in Appendix C. The outliers were no longer present in O2 data and therefore they were inconsistent with the persistent nature of CW signals. Finally, in Sec. VB upper limits on the strain amplitude for the 11 targets are discussed.

TABLE III. The table reports the outliers found in our analysis for each analyzed pulsar. The first column is the name of the pulsar; the second indicates the number of outliers found in the analysis. The third and fourth columns show respectively the outlier frequency and spin-down. The last column reports the corresponding p-value. For the two targets J1813-1749 and J1952 + 3252 the outliers did not undergo the follow-up procedure due to the fact that they can easily be associated with known noise lines; see Appendix B.

Name	Number of candidates	Frequency (Hz)	Spin-down (Hz/s)	p-value
J0205 + 6449	1	30.4046480	$-8.937 \times 10^{-11}$	0.003
J0534 + 2200 (Crab)	1	59.3702101	$-7.3920 \times 10^{-10}$	0.005
J0835-4510 (Vela)	1	22.3884563	$-3.12 \times 10^{-12}$	0.009
J1813-1246	2	41.5779102, 41.5852264	$-1.285 \times 10^{-11}$ , $-1.284 \times 10^{-11}$	0.007, 0.005
J1813-1749	36	Close to 44.705 Hz	...	$<10^{-6}$
J1833-1034	1	32.2807633	$-1.0535 \times 10^{-10}$	0.0004
J1952 + 3252	6	Close to 50.601	...	$<10^{-5}$
J1400-6326	2	64.1089253, 64.1406011	$-8.008 \times 10^{-11}$ , $-8.937 \times 10^{-11}$	0.002, 0.003
J2022 + 3842	1	41.1603319	$-7.297 \times 10^{-11}$	0.007

### A. Outliers outlook

We have found initial outliers for 9 of the 11 analyzed pulsars. More precisely, for most pulsars we have found one or two outliers, with the exception of J1813-1749 (36 outliers) and J1952 + 3252 (6 outliers). For J2043 + 2740 and J2229 + 6114 no outlier has been found. A summary of the outliers found in the analysis is given in Table III. The follow-up has clearly shown that in the case of J1952 + 3252 and J1813-1749 the outliers arise from noise disturbances in LIGO H (for J1813-1749) and in LIGO L (for J1952 + 3252); see Appendix B for more details. Most of the remaining outliers show an inconsistent time evolution of the SNR together with a low coherence between LIGO H and LIGO L and hence have been ruled out. As mentioned before, two outliers, one for J1833-1034 and one for Vela, have shown promising features during the basic follow-up steps: no known noise line is present in their neighborhood, the amplitude estimation is compatible and nearly constant among the LIGO L and LIGO H runs and their SNR appears to increase with respect to the integration time (see Fig. 3). Even if the trend of the SNR does not increase monotonically with time, as expected for real signals, we have decided to follow up on these outliers due to the fact that they show a completely different SNR trend with respect to all the other outliers found in this work. Moreover each outlier's significance increases in the multi-IFOs search, suggesting a possible coherent source.

J1833-1034 and Vela outliers: In order to establish if the outliers were not artifacts created by the narrow-band search, we also looked for the two outliers using two other analysis pipelines for targeted searches, which used a Bayesian approach: one designed for searching for nontensorial modes in CW signals [30], and the other developed for canonical CW target searches<sup>6</sup> and parameter

estimation [32]. Both pipelines produced odds, listed in Table IV, which show a small preference for the presence of a candidate compatible with general relativity. The odds values are not surprising due to the fact that we are using values for the frequency and the spin-down which are fixed to the ones found in the narrow-band search. Hence, a trial factor should be taken into account in order to make a robust estimation on the signal hypothesis preference. Besides the previous considerations, the values in Table IV clearly show that the outliers are not artifacts created by the narrow-band pipeline. We have also compared the estimation of the outlier parameters obtained from the *five-vector*,  $\mathcal{F}$ -statistic and Bayesian [6,8,32] pipelines. The inferred parameters are listed in Table V and seem to be compatible among the three independently developed targeted pipelines, thus suggesting the true presence of these outliers inside the data.

In order to establish each outlier's nature, a complete understanding of the noise background is needed. For this reason the first check was to look at the DS distribution in the narrow-band search. In the presence of a true signal we expect to see a single significant peak in the DS. Figure 4 shows the distribution of the DS (maximized over the spin-down corrections) for J1833-1034 and for Vela over the

TABLE IV. Odds obtained for the two outliers by the Bayesian pipelines [30,31]. The second column shows the odds of any nontensorial signal hypothesis versus the canonical CW signal hypothesis, the third column is the odds ratio of the canonical signal hypothesis vs the Gaussian noise hypothesis, and the last column is the odds ratio between the coherent signal among the two detectors vs the hypothesis that the outliers arise from an incoherent noise between LIGO H and L.

Name	$\log_{10} \mathcal{O}_{GR}^{nGR}$	$\log_{10} \mathcal{O}_N^S$	$\log_{10} \mathcal{O}_I^C$
J0835-4510 (Vela)	-0.55	2.30	1.07
J1833-1034	-0.73	2.73	1.34

<sup>6</sup>Frequency and spin-down value are fixed to the outlier's value found in the narrow-band search.

TABLE V. Estimation of the GW parameters,  $h_0$ ,  $\cos \iota$  and  $\psi$ , from three targeted search pipelines [6,8,32]. The intervals refer to the 95% confidence level.

J0835-4510 (Vela)	$h_0 \times 10^{-25}$	$\cos \iota$	$\psi$ (rad)
Five-vector	$5.7^{+2.3}_{-2.1}$	$-0.09^{+0.27}_{-0.19}$	$0.69^{+0.57}_{-0.58}$
Bayesian	$6.6^{+3.1}_{-3.7}$	$-0.14^{+0.28}_{-0.48}$	$0.57^{+0.31}_{-0.30}$
$\mathcal{F}$ -statistic	7.1	-0.13	0.55
J1833-1034	$h_0 \times 10^{-25}$	$\cos \iota$	$\psi$ (rad)
Five-vector	$1.6^{+0.5}_{-0.6}$	$0.10^{+0.30}_{-0.20}$	$0.58^{+0.35}_{-0.51}$
Bayesian	$1.8^{+0.8}_{-1.7}$	$0.24^{+0.64}_{-0.31}$	$0.58^{+0.56}_{-0.51}$
$\mathcal{F}$ -statistic	2.0	0.22	0.59

frequency band analyzed. We notice that for J1833-1034 the outlier is the only clear peak present in the analysis, surrounded by several lower peaks in the detection statistic which are not above the corresponding p-value threshold. On the other hand, for Vela, several peaks in the DS are present, with significance below but similar to that of the outlier, thus suggesting that the Vela outlier can be due to a non-Gaussian background.

A further test consists of checking the distribution of the DS in a narrow-band search performed using the same frequency/spin-down region but in a sky position shifted by about 0.5 degrees. Using this method we keep the contribution of non-Gaussian noise in the DS while removing a possible signal contribution. Figure 5 shows the distribution of the DS obtained for the J1833-1034 and Vela outliers. In both cases no over-threshold peaks are present; however the analyzed bands seem similarly polluted by non-Gaussian contributions which produce peaks in the DS. We have also studied the significance of

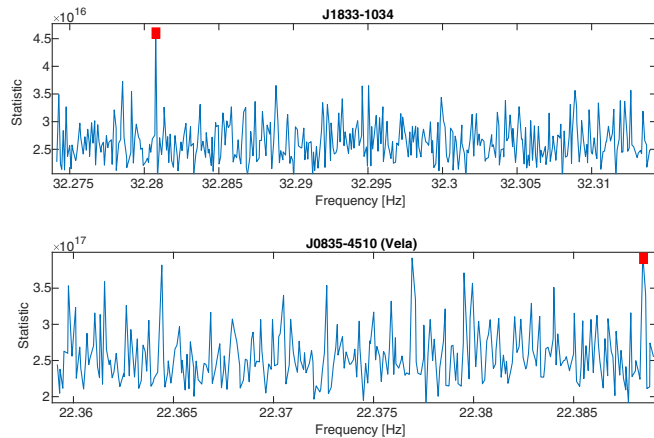


FIG. 4. Values of the local maximum of the DS over the spin-down corrections and the frequency subbands for J1833-1034 (top panel) and Vela (bottom panel). The outliers are highlighted with the red square.

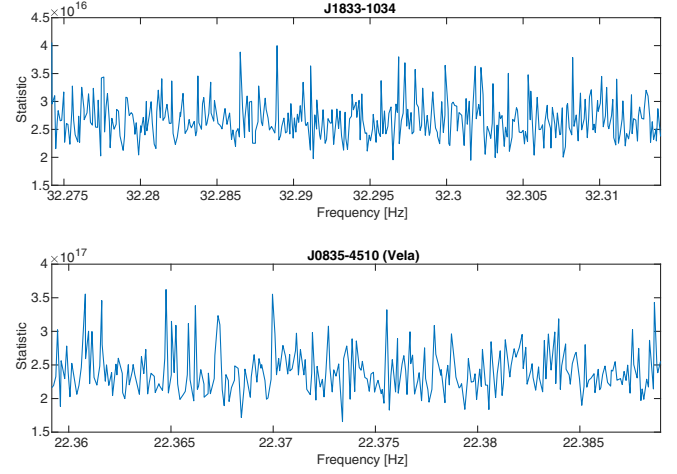


FIG. 5. Values of the local maximum of the DS over the spin-down corrections and the frequency subbands for a close position to J1833-1034 (top panel) and Vela (bottom panel).

the outliers using two of the three targeted search pipelines. As done previously, we have built a noise distribution of the DS, performing the targeted searches in other sky positions in order to compute the outliers' p-value. Using the trials factor from the narrow-band search we have found the outliers to have a higher resulting p-value with respect to the 1% threshold used in the initial outlier selection process during the narrow-band search, increasing the likelihood that these outliers were generated from noise. Some of the previous tests disfavor the signal hypothesis and seem to indicate the presence of a coherent noise disturbance among the interferometers. Previous works such as [7,33] have already pointed out the presence of some nontrivial coherent noise artifacts among the IFOs

TABLE VI. Median over the analyzed frequency band of the upper limits obtained on the GW amplitude for the 11 known pulsars. In the fourth column we report the ratio between the spin-down limit listed in Table II and the median of the upper limit; uncertainties correspond to the  $1\sigma$  confidence level and are due to the uncertainties on the pulsars' distances. The last column reports the median upper limit on the fraction of rotational energy lost due to GW emission.

Name	$h_{ul} \times 10^{-25}$	$\epsilon_{ul} \times 10^{-4}$	$h_{ul}/h_{sd}$	$\dot{E}_{rot}/\dot{E}_{GW}$
J0205 + 6449	3.76	7.7	$0.54 \pm 0.09$	0.29
J0534 + 2200 (Crab)	1.08	0.58	$0.07 \pm 0.02$	0.005
J0835-4510 (Vela)	9.28	5.3	$0.27 \pm 0.02$	0.07
J1400-6326	1.17	2.7	$1.3 \pm 0.4$	...
J1813-1246	1.80	2.5	$> 1.0$	...
J1813-1749	1.9	4.8	$0.64 \pm 0.04$	0.41
J1833-1034	3.08	13	$0.99 \pm 0.09$	...
J1952 + 3252	1.31	1.4	$1.31 \pm 0.22$	...
J2022 + 3842	1.90	11	$1.77 \pm 0.35$	...
J2043 + 2740	14.4	47	$2.07 \pm 0.83$	...
J2229 + 6114	1.78	3.4	$0.54 \pm 0.35$	0.30



which can produce outliers. For this reason, in the spirit of what was done in [33], we have looked at O2 data. If the outliers are really due to a “standard” CW signal, they are expected to be present also in O2 data, due to their persistent nature. We have analyzed the data using the narrow-band pipeline but no evidence for these outliers was found in data. In conclusion the outliers are not true

CW signals. More details on the O2 analysis can be found in Appendix C.

### B. Upper limits

Following the procedure described in Appendix A we have set 95% C.L. upper limits on GW strain amplitude in

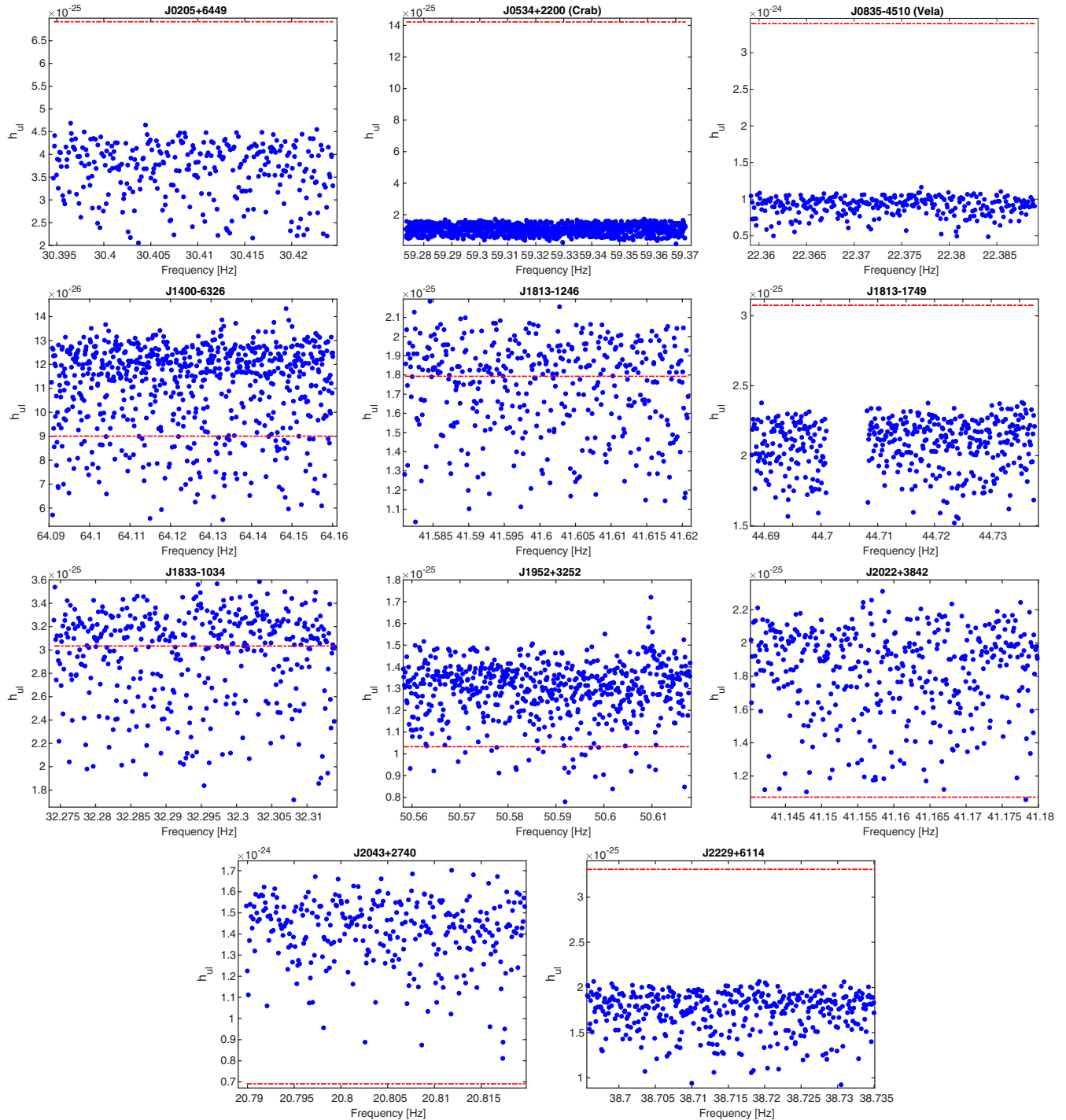


FIG. 6. Plots of the 95% C.L. upper limit on the GW amplitude for the 11 pulsars. The blue dots indicate the amplitude upper limits set with our analysis; the red dashed lines indicate the theoretical spin-down limit in Table II.

every  $10^{-4}$  Hz subband. In each of these bands the upper limit was computed by injecting simulated GW signals with several different amplitudes and finding the amplitude such that 95% of the injected signals with that amplitude produce a value of the DS corresponding to the nominal overall p-value of 1%. Table VI gives an overview of the overall sensitivity reached in our search using the median of the upper limits among the analyzed frequency bands; for graphs of the upper limits see Fig. 6. For J2043 + 2740, J1952 + 3252 and J2022 + 3842 our overall sensitivity is clearly above the spin-down limit. For J1813-1246 and J1833-1034 our overall sensitivity is close to the spin-down limit, producing values of the upper limits both below and above the spin-down limit. For J1400-6326 we have obtained a large fraction of the upper limits in the narrow-band search below the spin-down while for J0205 + 6449 and J2229 + 6114 we have beaten the spin-down limit in a narrow-band search for the very first time. For Crab and Vela pulsars we have obtained upper limits respectively  $\sim 7$  and  $\sim 3.5$  times lower than those computed in a past analysis [11]. This improvement is due to a combination of two factors: the enhanced sensitivity of advanced detectors and the choice to compute upper limits over  $10^{-4}$  Hz subbands instead of the full analysis band, thus reducing the impact of the look-elsewhere effect in each subband [12]. Finally the narrow-band search for J1813-1749 beats the spin-down limit (if we exclude from the search the frequency region around the LIGO H artifact), constraining for the first time their CW emission. Pulsars J1813-1749 and J1400-6326 have not been previously analyzed in targeted searches, due to the lack of ephemeris covering O1 or previous runs. Even if we consider the uncertainties on the pulsars' distances, propagated in Table VI for the spin-down limit and upper-limit ratio, we are still able to beat the spin-down for those five pulsars.

## VI. CONCLUSION

In this paper we have reported the result of the first narrow-band search using Advanced LIGO O1 data for 11 known pulsars. For each pulsar, a total of about  $10^7$  points in the frequency and spin-down space have been explored. For nine pulsars, outliers have been found and analyzed in a follow-up stage. Most of the outliers did not pass the follow-up step and were labeled as noise fluctuations or instrumental noise artifacts. We have found two near-threshold outliers, one for J1833-1034 and another for the Vela pulsar, which needed deeper studies but eventually were rejected. In particular, the outliers have been searched for in the first five months of the LIGO O2 run and were not confirmed. We have computed upper limits on the signal strain, finding for five pulsars values below the spin-down limit in the entire narrow-band search (Crab, J1813-1749, J0205 + 6449, 2229 + 6114 and Vela). For the Crab and Vela pulsars the

upper limits significantly improve with respect to past analyses. For an additional three targets (J1833-1034, J1813-1246 and J1400-6326), the median upper limit across the search bands is below or very close to the spin-down limit. For J1813-1749, which has never been analyzed in a targeted search, we have beaten the spin-down limit for the first time while for J0205 + 6449 and J2229 + 6114 the spin-down limit has been beaten for the first time in a narrow-band search. Seven of the 11 pulsars analyzed in this work were also analyzed using O1 data in a target search [7]. The upper limits found in this work are about 2–3 times higher with respect to targeted searches: the sensitivity loss is due to the fact that we are exploring a large number of templates in the frequency spin-down plane. On the other hand we have put for the first time upper limits in a small frequency spin-down region around the expected values.

The analysis of forthcoming Advanced LIGO and Virgo runs [34], with improved sensitivities and longer durations, could provide the first detection of continuous gravitational signals from spinning neutron stars, which would help to shed light on their structure and properties.

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## APPENDIX A: UPPER LIMIT

Once we have concluded that our data are compatible with noise, upper limits on the GW amplitude can be computed. The upper-limits computation consists of injecting many different signals with fixed amplitude  $H_0$  and parameters  $\eta$ ,  $\psi$  with a uniform distribution into the real data. According to the frequentist paradigm, the 95% confidence level upper limit can be computed by asking that the 95% of the injected signals provide a value of the DS greater than the threshold for the candidates' selection used in the analysis. The signal must be injected at the beginning of the analysis, i.e. before the Doppler corrections, and the entire analysis procedure must be followed in order to compute the DS. This procedure is not suitable for narrow-band searches due to the fact that an injection is needed in every analyzed frequency subband. This problem can be overcome by injecting simultaneously many different signals in many different frequency subbands in just one data set and then performing the narrow-band search. Repeating this step  $N$  times produces  $N$  different data sets, each containing a signal in each analyzed frequency subband. Then for each subband we ask for the 95% DSs produced by the injected signal to be greater than the value used for the candidates' selection, obtaining in this way the value of the upper limit for a given frequency subband. Practically this procedure is done using several tricks in order to speed up the computation, as detailed in the following. First of all we assume that our data are the linear superposition of noise  $n(t)$  and an injected signal  $h_{\text{inj}}(t)$ ,

namely  $s(t) = n(t) + h_{\text{inj}}(t)$ . According to the linearity of the FFT, the five-vector of  $s(t)$  will be the summation of the two independent five-vectors of the noise and the injected signal:

$$\vec{X} = \vec{X}_{\text{noise}} + \vec{X}_{\text{inj}}. \quad (\text{A1})$$

The estimators of the GW polarization, which are the building blocks of the DS, are linear due to the scalar product with respect to the sidereal templates  $A_+(t), A_\times(t)$ . Hence using Eq. (A1) we can write the analysis estimator as

$$\hat{H}^{+/\times} = \hat{H}_{\text{noise}}^{+/\times} + \hat{H}_{\text{inj}}^{+/\times}. \quad (\text{A2})$$

Equation (A2) indicates that before the calculation of the DS we can keep separate the estimators computed from our real data set and the ones arising from an injected signal. This leads to the possibility to change the GW amplitude  $H_0$  of the injected signal, directly rescaling the absolute value of the estimators  $\hat{H}_{\text{inj}}^{+/\times}$  without reperforming all the corrections in the time domain and thus saving computational time. As stressed before, the form of the injected signal  $h_{\text{inj}}(t)$  should be built in such a way as to contain a signal in each analyzed frequency subband. Formally we can write  $h_{\text{inj}}(t)$  as the superposition of  $N$  different signals, each one located in a random-frequency bin of each frequency subband:

$$h_{\text{inj}}(t) = H_0 [H^+ A_+(t) + H^\times A_\times(t)] e^{i\phi_0} \times \sum_{S=1}^N e^{i\phi_{\text{Röm}}^S(t)} e^{i\phi_{\text{rot}}^S(t)}, \quad (\text{A3})$$

where  $\phi_{\text{Röm}}^S(t)$  and  $\phi_{\text{rot}}^S(t)$  are the usual phase evolution due to the Römer and rotational frequency evolution of the signal  $\mathcal{S}$  [12]. Assuming that the  $N$  different signals are injected with a constant frequency step  $\Delta f_{\text{inj}}$  in the frequency grid starting from a frequency  $f_0$ , i.e.  $f_S = f_0 + S\Delta f_{\text{inj}}$ , we can manipulate the Eq. (A3) to obtain

$$h_{\text{inj}}(t) = H_0 [H^+ A_+(t) + H^\times A_\times(t)] \times e^{i\phi_{\text{Röm}}^0(t)} e^{i\phi_{\text{rot}}^0(t)} e^{i\phi_0} \sum_{S=1}^N e^{i2\pi S\Delta f_{\text{inj}}(t+p(t))}, \quad (\text{A4})$$

where  $p(t)$  is the Römer correction and the superscript 0 refers to the phase evolution of a signal injected at the frequency  $f_0$ . By defining  $k = 2\pi i\Delta f_{\text{inj}}(t + p(t))$ , we can now exploit the geometrical series present in Eq. (A4) to write

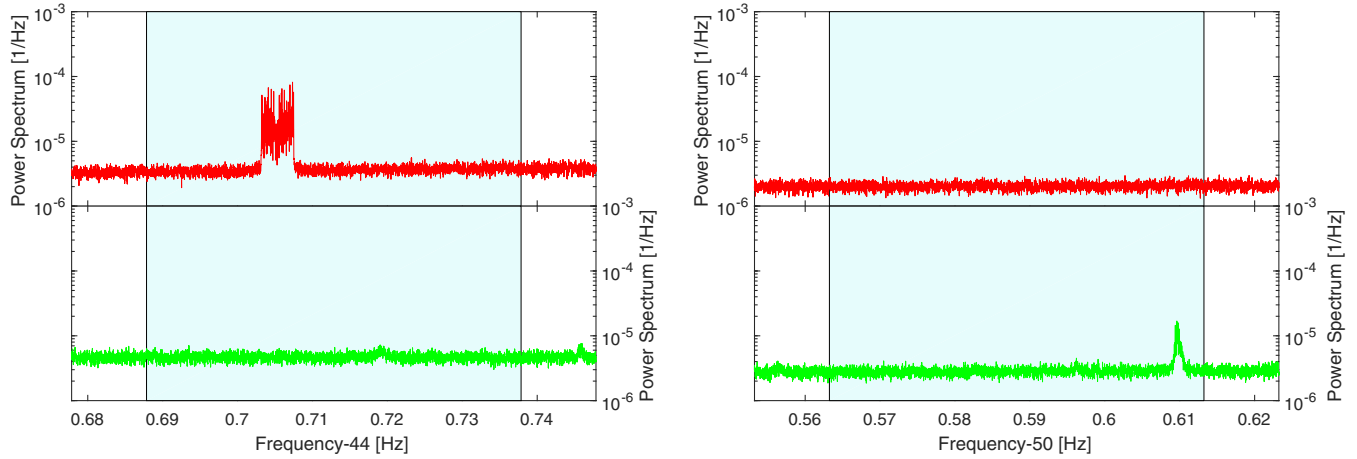


FIG. 7. Left: Power spectrum of Hanford (red line) and Livingston (green line) data inside the frequency region explored by the narrow-band search (blue box) around J1813-1749. Right: Power spectrum of Hanford (red line) and Livingston (green line) data inside the frequency region explored by the narrow-band search (blue box) around J1953 + 3252.

$$h_{\text{inj}}(t) = H_0 [H^+ A_+(t) + H^\times A_\times(t)] e^{i\phi_{\text{Rdm}}^0(t)} e^{i\phi_{\text{rot}}^0(t) \frac{1-e^{(N+1)k}}{1-e^k}}. \quad (\text{A5})$$

Practically in our analysis, for each data set, we select a random frequency bin in the first analyzed frequency subband and then we replicate it on the frequency grid using Eq. (A5) and setting  $\Delta f_{\text{inj}}$  equal to the width of the subbands. This procedure together with the linearity of the FFT allows us to strongly reduce the computational time while obtaining the same results.

## APPENDIX B: KNOWN INSTRUMENTAL NOISE LINES

The data from the gravitational waves interferometer are polluted by several instrumental noise lines. Many of these disturbances have been identified during the run. Their presence can produce in the analysis a large number of outliers. We have found that the 36 outliers J1813-1749 are due to a noise line associated with the magnetometer channels in Hanford at 44.7029 Hz. The presence of the noise line can also be seen in the left panel of Fig. 7, where we show the power spectrum around the region explored by the narrow-band search. Concerning the six outliers from J1952 + 3252, we know that they are due to an artifact that is part of a 1.9464 Hz comb in the Livingston data. This disturbance is shown in the power spectrum in the right plot of Fig. 7.

## APPENDIX C: O2 FOLLOW-UP OF THE OUTLIERS

We have used these data in a narrow-band search in order to check if the outliers found for J1833-1034 and Vela in O1 were still present. The parameters of the narrow-band searches have been set in such a way to cover the expected

frequency and spin-down of the outlier during the O2 epoch. The Vela pulsar glitched on December 12, 2016, between 11:31 and 11:46 UT.<sup>7</sup> The glitch has been classified as a canonical Vela glitch [35]. In order to prevent the glitch from affecting our analysis we have started to analyze data from January 12, 2017, when the spin-down variation is supposed to be recovered. Moreover we have also increased the spin-down range by a factor of 3.7 with respect to the O1 analysis. A summary of the narrow-band search parameters is given in Table VII.

Our analysis has produced no significant outlier for either J1833-1034 or Vela. Figure 8 shows the histograms of the DS obtained in the narrow-band search with respect to the threshold for the outliers' selection, for J1833-1034 and Vela respectively. In order to estimate our sensitivity in this search and compare the results with the sensitivity reached in O1, we have also computed the upper limit on the GW amplitude  $h_0$  for J1833-1034 and Vela over the narrow-frequency region explored. The procedure that we have used is the same used for O2, and the values of the upper limits are shown in Fig. 9 for J1833-1034 and Vela respectively. The median value of the amplitude upper limit for J1833-1034 is  $1.25 \times 10^{-25}$  which is nearly a factor 2 lower than the one obtained for the O1 analysis in Table VI, thus indicating that if the outlier found in O1 were a true persistent CW signal, it would have appeared in the O2 analysis with a higher significance. Similarly, for Vela we have obtained a median value of the amplitude upper limit of  $3.41 \times 10^{-25}$  which is about three times better than the one obtained in the O1 analysis; see Table VI. We then conclude that both outliers are not confirmed in O2.

<sup>7</sup><http://www.astronomerstelegam.org/?read=9847>

TABLE VII. This table reports the explored range for the rotational parameters of each pulsar. The columns are the central frequency of the search ( $f_0$ ), explored frequency band ( $\Delta f$ ), central spin-down value of the search ( $\dot{f}_0$ ), explored spin-down band ( $\Delta \dot{f}_0$ ), the frequency ( $f_{O2}$ ) and spin-down ( $\dot{f}_{O2}$ ) of the outliers at the O2 epoch reference time on November 30, 2016.

Name	$f_0$ (Hz)	$\Delta f$ (Hz)	$\dot{f}_0$ (Hz/s)	$\Delta \dot{f}$ (Hz/s)	$f_{O2}$ (Hz)	$\dot{f}_{O2}$ (Hz/s)
J0835-4510 (Vela)	22.37289950	0.05	$-3.1159 \times 10^{-11}$	$2.4024 \times 10^{-13}$	22.38712428	$-3.1128 \times 10^{-11}$
J1833-1034	32.29004216	0.05	$-1.0542 \times 10^{-10}$	$1.7266 \times 10^{-13}$	32.27625775	$-1.0534 \times 10^{-10}$

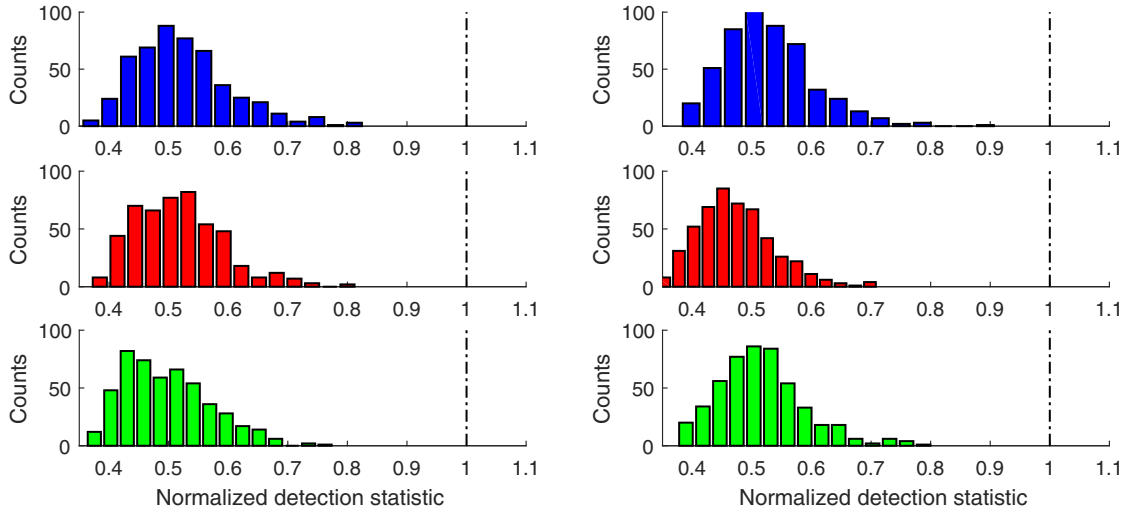


FIG. 8. Left: Histograms of the DS obtained in the J1833-1034 O2 narrow-band search for the (top) joint search, (middle) Handford search, and (bottom) Livingston search; the x-axis is normalized to the DS threshold in each search. Right: Histograms of the DS obtained in the Vela O2 narrow-band search for the (top) joint search, (middle) Handford search, and (bottom) Livingston search; the x-axis is normalized to the DS threshold in each search.

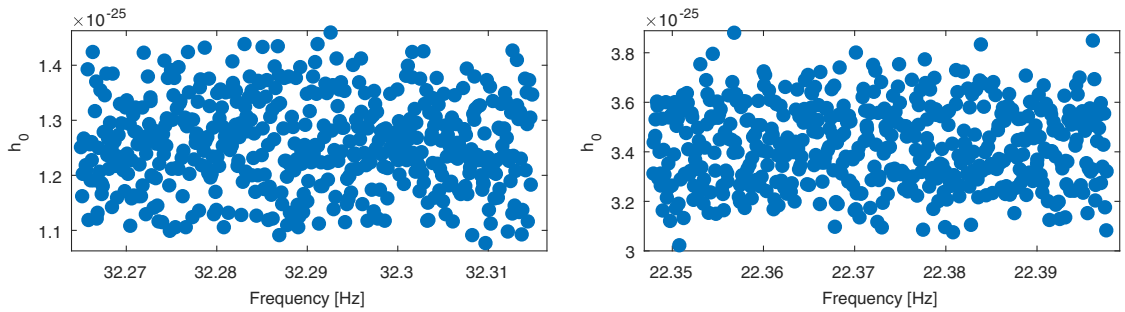


FIG. 9. Left: Upper limits on the GW amplitude  $h_0$  over the narrow-frequency region analyzed in O2 for J1833-1034. Right: Upper limits on the GW amplitude  $h_0$  over the narrow-frequency region analyzed in O2 for Vela.

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