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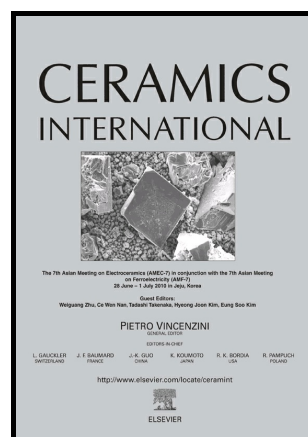
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Crystallochemistry and electrical properties of Al-doped Sr₂SiO₄ electrolytes

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Abstract

A new Sr₂Si_{1-x}Al_xO_{4-x/2}□_{x/2} (x=0, 0.028, 0.06 and 0.075) silicate series has been prepared and the structural and electrical properties were investigated. These samples present three different polymorphs (β, α'_L and α'_H) and the increase of Al-content resulted in a shift of the polymorph content, from a monoclinic form to an orthorhombic one. Aluminium contents higher than x = 0.06 did not lead to any further change in the ratio of polymorphs. High temperature X-ray diffraction studies showed a complete stabilization of α'_L, presenting a superstructure orthorhombic form, at temperatures between RT and 250 °C. At temperatures higher than 500 °C a basic orthorhombic cell is stabilised. Thermogravimetric analysis displayed an increase of the water uptake on doping, due to the generation of oxide vacancies in the lattice. Electrical characterization in dry and wet atmospheres showed an important proton contribution to the conductivity. In addition, measurements in N₂ and O₂ fluxes displayed a significant p-type electronic contribution to the overall conductivity.

Keywords: SOFC, electrolyte, Sr₂SiO₄, Al-doping

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1. Introduction

Ceramic ion conductors have numerous applications in different areas, such as sensors, fuel cells, etc. Yttria stabilized zirconia (YSZ) [1] and gadolinia doped ceria (CGO) [2], with a fluorite structure, are the state-of-the-art materials. However, these materials present several inconveniences, such as a necessity for high operation temperatures and low stability in reducing environments, respectively. Alternative oxide ion conductors with different structures have also been investigated: materials based on a perovskite, such as $\text{La}_{0.9}\text{Sr}_{0.1}\text{Ga}_{0.8}\text{Mg}_{0.2}\text{O}_{2.85}\square_{0.15}$ [3]; brownmillerite-type, such as $\text{Ba}_2\text{In}_2\text{O}_5$ [4]; $\text{La}_2\text{Mo}_2\text{O}_9$ [5]; apatite-based structure, such as $\text{La}_{9.33}\text{Si}_6\text{O}_{26}$ [6]; LaBaGaO₄-like, with β - K_2SO_4 type structure [7]; and, melilite-type, such as $\text{LaSrGa}_3\text{O}_7$ [8].

Proton conductors have received great attention due to their generally better conductivity at reduced operation temperatures. Most of the research on electrolytes for proton conductors solid oxide fuel cells (PC-SOFCs) has been focused on perovskite-type materials based on BaCeO_3 and BaZrO_3 [9-12]. The former are the state-of-the-art, with very high conductivity values ($\sim 10^{-2} \text{ S cm}^{-1}$ at 600 °C), however, these materials are prone to carbonation in the presence of CO_2 . [13,14] For the latter, BaZrO_3 has a much better chemical stability towards CO_2 , however, the high temperatures required for the preparation of dense ceramics lead to high grain boundary resistances and low conductivity values.

In the last few years, our research group have been involved in the identification of new electrolytes for intermediate temperature SOFCs, mainly focusing on systems where “loosely bound” oxide anions were the charge carriers, for instance, $\text{Ca}_3(\text{SiO}_4)\text{O}$ [15,16], $\text{Sr}_3(\text{SiO}_4)\text{O}$ [17] and $\text{Ln}_2(\text{TO}_4)\text{O}$ ($\text{Ln}=\text{La}, \text{Nd}, \text{T}=\text{Si}, \text{Ge}$) [18]. In those systems four oxygen atoms form part of isolated tetrahedral units and the remaining one is “loosely bound” and prone to present a higher mobility. In order to enhance the conductivity, different doping strategies in the tetrahedral position were carried out to increase the

oxygen vacancy concentration. As result, moderate conductivity values ($\sim 10^{-4}$ S cm⁻¹ at 800 °C) were observed, with the presence of a significant proton contribution to the overall conductivity.

Very recently, the structural and electrical properties of aluminium doped dicalcium silicate, Ca₂SiO₄, a major constituent of belite Portland cements, were studied. [19] Dicalcium silicate exhibits several polymorphs on heating (γ , β , α'_L , α'_H and α) and it is worth mentioning that the transition β (monoclinic) \rightarrow γ (orthorhombic) gives rise to a significant change in the cell volume, degrading the material. [20-22] Cuesta *et al.* [19] demonstrated that Al-doping in the silicon position stabilized the β form, preventing the formation of the γ -polymorph and the consequent ceramic failure. In addition, an increase in conductivity was observed on doping, due to the generation of oxide vacancies.

The analogous distrontium silicate exhibits three polymorphs: a monoclinic form similar to β -Ca₂SiO₄, and two orthorhombic polymorphs on heating, the first one displaying a superstructure cell, both related to α'_L and α'_H -Ca₂SiO₄, respectively. [23-31] The structure of Sr₂SiO₄ is analogous to that of Ca₂SiO₄ and based on similar structural units: M²⁺ ions (M = Ca, Sr) and isolated SiO₄ tetrahedra, where there are two different strontium sites, Sr1 and Sr2, seven- and eight-fold coordinated, respectively. The structure of both monoclinic and orthorhombic Sr₂SiO₄ polymorph are displayed in Fig. 1.

Adopting the nomenclature used for Ca₂SiO₄, it is worth mentioning that the transitions $\beta \rightarrow \alpha'_L$ and $\alpha'_L \rightarrow \alpha'_H$ occur at ~ 85 and ~ 500 °C, respectively, much lower temperatures than for Ca₂SiO₄, where the analogous transitions take place at 690 and 1160 °C, respectively. [30] The α' form of Sr₂SiO₄ is expected to be stabilized at room temperature by doping with appropriate cations. In addition, it has been reported that the α' form of Sr₂SiO₄ may be stabilized at room temperature through replacing Sr by small amounts of Ba. [25]

The objective of the present study is to investigate the effect of Al-doping on the structure and electrical properties of Sr₂SiO₄-based materials. An enhancement in the electrical properties is expected

due to the fact that the aliovalent substitution of silicon by aluminium generates oxide vacancies in the lattice.

2. Experimental section

2.1 Synthesis and powder characterization

$\text{Sr}_2\text{Si}_{1-x}\text{Al}_x\text{O}_{4-x/2}\square_{x/2}$ ($x= 0, 0.028, 0.06$ and 0.075) series was synthesized by reactive sintering using high-purity oxides and carbonates as starting materials: SrCO_3 (Alfa, 99.99%), SiO_2 (Alfa, 99.5%) and $\gamma\text{-Al}_2\text{O}_3$ (Alfa, 99.997%). Strontium carbonate was precalcined at $1400\text{ }^\circ\text{C}$ for 15 h in order to achieve its full decarbonation. Raw mixtures were ground and preheated at $1050\text{ }^\circ\text{C}$ for 4 hours (heating rate of $10\text{ }^\circ\text{C min}^{-1}$) in order to prepare 3 g of each sample. Under these conditions only a partial reaction takes place. The resulting powders were reground in a Fritsch ball mill (model Pulverisette 7, 45 cm^3 agate vessel containing 7 agate balls with a diameter of 15 mm) for 30 min at 200 rpm with reverse rotation every 10 min and pelletized (10 mm diameter, ~ 1 mm of thickness, 500 MPa). A second thermal treatment was carried out at $1550\text{ }^\circ\text{C}$ for 6 hours (heating rate of $10\text{ }^\circ\text{C min}^{-1}$) and the sample was slowly cooled down to room temperature ($10\text{ }^\circ\text{C min}^{-1}$). Compaction was calculated taking into account the pellet mass, volume and the crystallographic density. Three pellets were prepared under identical experimental conditions: one (after grinding) for the structural and thermal analysis and two additional ones for impedance and microstructural characterization, respectively. Hereafter, $\text{Sr}_2\text{Si}_{1-x}\text{Al}_x\text{O}_{4-x/2}\square_{x/2}$ series is labelled as Sr_2Al_x (Table 1).

X-ray powder diffraction (XRPD) patterns were recorded using a PANalytical X'Pert Pro diffractometer with $\text{CuK}_{\alpha 1}$ radiation to determine the purity and polymorphism of the materials. Variable temperature experiments were carried out in an Anton Paar HTK1200 under static air. Data were collected at 30, 250, 500, 750 and $900\text{ }^\circ\text{C}$ after a delay time of 15 min at each temperature to ensure thermal stabilization.

Phase identification was performed with X'Pert HighScore Plus [32] and Rietveld analysis with GSAS software [33] using the ICSD database for the starting structural models [34]. The usual profile parameters (scale factors, background coefficients, zero-points, half-width, pseudo-Voigt

and asymmetry parameters for the peak-shape) were refined. The atomic parameters were fixed and not refined.

2.2. Thermal analysis

Thermogravimetric analysis (TGA) data were recorded on a SDT-Q600 analyzer supplied by TA Instruments. The temperature was varied from RT to 900 °C at a heating/cooling rate of 10 °C min⁻¹. Different heating/cooling cycles were carried out to study the thermal reversibility and reproducibility of the measurements on samples in open platinum crucibles under a mixed flow of dry-air (10 mL min⁻¹) and wet-air (15 mL min⁻¹) saturated by bubbling in water at 20 °C.

2.3. Microstructural characterization and conductivity measurements

The microstructure of the ceramics was observed by scanning electron microscopy (SEM) (Jeol JSM-6490LV) combined with energy dispersive spectroscopy (EDS). The grain size of the sintered pellets was estimated from SEM micrographs, using the linear intercept method from at least 30 random lines and three different micrographs with the help of image-analysis software. [35] Platinum electrodes were made by coating opposite pellet faces with METALOR® 6082 platinum ink and heating to 800 °C at a rate of 5 °C min⁻¹ for 15 min in air to decompose the paste and to harden the Pt residue. Impedance spectra were obtained using a frequency response analyser (Solartron 1260) in different dry and wet gases (N₂ and O₂) in the 0.1 Hz to 1 MHz frequency range with an ac perturbation of 100 mV. The spectra were recorded upon cooling from 900 to 500 °C with a stabilization time of 10 minutes between consecutive measurements. Impedance spectra were analyzed using the ZView program. [36] The resistance and capacitance values of the different contributions were obtained by fitting the impedance spectra data with an equivalent circuits formed by (RQ) elements in series, where R is a resistance and Q is a pseudocapacitance in parallel.

Conductivity measurements as a function of oxygen partial pressure $p(\text{O}_2)$, from $\sim 10^{-22}$ atm to 0.21 atm were performed in a closed tube-furnace cell. $p(\text{O}_2)$ values were monitored using a YSZ oxygen sensor positioned close to the sample. The conductivity was continuously recorded as a function of $p(\text{O}_2)$. The process consisted of flushing the system with a dry gas mixture of 5% H₂ in

Ar for 12 h at 800 °C and 900 °C in order to reach a minimum in oxygen activity inside the furnace and ensure that the sample was close to the equilibrium under those conditions. Next, the flushing was switched off, and the oxygen partial pressure was allowed to slowly recover to atmospheric pressure by free diffusion, since the system was not fully airtight. Each isothermal cycle took more than 48 h to complete.

The ion transference numbers were determined by the modified electromotive force method (emf) proposed by Gorelov. [37,38] This modification of the classical emf technique eliminates possible errors in the determination of ion transference numbers, arising due to electrode polarization. These errors are not negligible for electrolyte materials which have relatively low electronic conductivity. [38,39] The ionic transport numbers, t_o , were measured under a $p(\text{O}_2)$ gradient of dry 5% H_2 -Ar/air, using a continuous flux of these gases in the 650-900 °C temperature range. A YSZ tube was used to measure the theoretical emf under these conditions, E_{th} . The emf observed in the sample, E_{obs} , was measured with an external variable resistance R_M , in parallel to the measuring cell, varying from 100 M Ω to 1.5 k Ω . Experimental data were fitted with the equivalent circuit (Eqn (1)) proposed by Gorelov. [37,40] In this case, the ionic resistance R_o , the electronic resistance R_e and the polarization resistance R_η are related to the emf values (E_{th} and E_{obs}) from the relationship:

$$\frac{E_{th}}{E_{obs}} - 1 = (R_o + R_\eta) \left[\frac{1}{R_e} + \frac{1}{R_M} \right] \quad (1)$$

The dependence of $(E_{th}/E_{obs} - 1)$ vs $1/R_M$ is a linear plot with slope $(R_o + R_\eta)$ and the intercept of the $(1/R_M)$ -axis is equal to $(-1/R_e)$. Total resistance R_T was determined independently by impedance spectra data. Oxide ion transport numbers were calculated as: $t_o = 1 - R_T/R_e$.

3. Results and discussion

3.1. Single phase existence ranges

The samples were prepared by reactive sintering, a method where the synthesis and densification process take place in a single heating step. This methodology have been proven useful

with different materials, leading to pure and dense ceramic samples with an adequate microstructure [19,41]. After the heating process the pellet samples were ground and studied by X-ray powder diffraction.

XRPD patterns for $\text{Sr}_2\text{Si}_{1-x}\text{Al}_x\text{O}_{4-x/2}\square_{x/2}$ ($x= 0, 0.028, 0.06$ and 0.075) series are shown in Fig. 2. All samples are composed by two related polymorphs without the presence of any additional secondary phases. The solubility limit of aluminium in the reactive sintering conditions at $1550\text{ }^\circ\text{C}$ is close to $x=0.06$, above this concentration the cell volumes are quite similar, indicating that more aluminium cannot be incorporated into the structure. (see Table 1).

Rietveld refinements were performed by fitting only the overall parameters and weight fractions for all the phases. For the monoclinic phase, the structural description, previously reported by Catti *et al.*, was used [25]. For the orthorhombic structure, the models available in the ICSD database (No. 35666 and 35667), claimed to represent Sr_2SiO_4 above 85°C , did not satisfactory fit the XRPD pattern, where some small reflections were not given any intensity. [25,30]. However, as was reported by Withers *et al.* [30] and Liu *et al.* [31], the polymorph above $85\text{ }^\circ\text{C}$ presents a tripled superstructure, therefore, a structural model from the analogous $\alpha'_L\text{-Ca}_2\text{SiO}_4$ was used, where the a-axis is multiplied by three. Ca atoms were replaced by Sr and cell parameters were appropriately modified to match the basic unit cell for this polymorph. With this phase all small reflections were satisfactory fitted.

Fig. 3 shows an example of the fit for $\text{Sr}_2\text{Al}_{0.06}$. Good agreements factors were achieved (Table 1). The undoped phase, Sr_2SiO_4 , shows at room temperature two different polymorphs, monoclinic (β) and orthorhombic (α'_L), with percentages of 88 and 12 wt.%, respectively (Table 1). However, as the aluminium content increases the orthorhombic phase turns to be the main phase, with values close to 92 wt.%. For Al content $x > 0.06$ no change in the polymorph ratio is observed. It has to be noted that Al-doped Sr_2SiO_4 series present a different behaviour to analogous Ca_2SiO_4 [19]. The undoped calcium sample shows a mixture of β monoclinic and γ orthorhombic

polymorphs at room temperature, with the % monoclinic form increasing on Al-doping, with a 99 and 1 wt.% of β and γ -Ca₂SiO₄, respectively, for Ca₂Al_{0.028}.

As can be seen in Table 1, the unit cell parameters and cell volumes for Sr₂Si_{1-x}Al_xO_{4-x/2}□_{x/2} increase on doping, reaching its maximum at x=0.06, indicating the end of the successful replacement of Si by Al. This behaviour is expected due to the fact that Al³⁺ (0.39 Å) has a larger ionic radius than Si⁴⁺ (0.26 Å) in a tetrahedral environment. This result further confirms the partial substitution of silicon by aluminium.

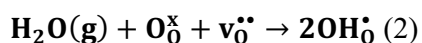
XRPD experiments on heating for Sr₂Al_{0.06} showed a transition between RT and 250 °C from a mixture of monoclinic and orthorhombic phases to a fully orthorhombic Pna2₁ phase (α'_L). Fig. 4 left and Table 2 show the 25-29 °2 θ range, where the reflection assigned to the monoclinic phase (112) disappear. Between 250 and 500 °C the (002) reflection is no longer present (see Fig. 4 right for the 29-32 °2 θ range) where a second polymorphic change was observed, from a superstructure cell to a basic orthorhombic symmetry. In the literature this high temperature orthorhombic phase was reported to be Pmnb, however in this work a Pnma space group, analogous to α'_H -Ca₂SiO₄, was also tested. The latter gave significantly lower R_F values, an improvement of almost 1.5%, thus, Pnma was used in all high temperature refinements. No transition to a hexagonal polymorph, $\alpha'_H \rightarrow \alpha$, as in Ca₂SiO₄, was detected up to 900 °C. It is worth mentioning that Al doping changes the polymorphism at room temperature but no differences were observed on heating.

The synthesis of Sr_{2-x/2}Al_{x/2}Si_{1-x/2}Al_{x/2}O₄ (x= 0.06), where Al is replacing both Sr and Si sites, was also attempted. However, despite the lack of any major segregation of secondary phases, the amount of monoclinic is very high and close to the value for the undoped sample, indicating that Al is not substituting for Sr. A similar behaviour was observed in previous related work for Sr₃SiO₅-based compounds, where Al is only incorporated in Si-site. [17]

3.2. Thermogravimetric analysis

Thermogravimetric measurements were carried out to analyse the effect of aluminium substitution on the oxygen vacancy concentration and water uptake. Measurements were recorded in humidified air and the curves on heating and cooling were completely reproducible. Only the curves acquired on cooling are plotted for better comparison of the data.

As can be seen in Fig. 5, the weight increase on cooling (below 700 °C) on Al-doping due to the higher oxygen vacancy concentration and the consequent incorporation of water with the formation of protonic defects, taking place as follows (Eq. 2), using the Kröger-Vink notation:



This water uptake is clearly seen if the mass gain for the undoped composition, Sr_2Al_0 , 0.01 wt% (0.001 moles of water per unit formula), is compared to that for $\text{Sr}_2\text{Al}_{0.075}$, 0.07% (0.010 moles of water per unit formula). For $\text{Sr}_2\text{Al}_{0.075}$ in dry conditions, the water uptake is significantly lower, 0.01 wt% (0.001 moles of water per unit formula), due to the no incorporation of water into the vacancies.

In the very high temperature region, 700-900 °C, a minor weight change is also observed likely ascribed to the oxygen uptake/desorption into the vacancies of the structure. This phenomenon will be further studied by impedance spectroscopy in the next section.

3.3. Microstructural analysis

The sintering conditions led to dense samples with compactions increasing on doping, for instance, $\text{Sr}_2\text{Al}_{0.0}$ exhibits a relative density of 85% and this value increases up to 95% for the Al-doped compounds. Fig. 6 shows the SEM micrographs of the sintered pellets, which show that the average grain size grows with the Al-doping: 3.6, 5.3, 6.3 and 8.5 μm for $\text{Sr}_2\text{Al}_{0.0}$, $\text{Sr}_2\text{Al}_{0.028}$, $\text{Sr}_2\text{Al}_{0.06}$ and $\text{Sr}_2\text{Al}_{0.075}$, respectively. Regarding porosity, it is significantly high for the undoped composition, as is reflected by the moderate compaction value about 85%. For the Al-doped samples the porosity decreases dramatically, indicating clearly that Al-doping improve the

densification and grain growth rate. No indications of liquid phase formation or phase segregations are found at the grain boundary region. Additionally, no contamination due to the ball-milling process was detected in the sintered pellets.

3.4. Electrochemical characterization

Impedance spectroscopy was carried out to analyse the electrical properties of the samples and establish the existence of mixed ionic oxide-proton conductivity. For this purpose, the measurements were performed under different flowing gases: dry and wet N₂ and O₂ gases. Fig. 7 shows representative impedance plots for Sr₂Al_{0.06} measured in the different gases at 800 °C. It is clear that there is only one contribution to the conductivity, ascribed to the grain interior with capacitance values of 4.7 pF cm⁻¹. Note that the grain boundary contribution is negligible in these samples. In addition, the impedance data were very well fitted to the following equivalent circuit: (R_bQ_b)Q_e, using the Zview program [36], where the subscripts b and e denote grain interior (bulk) and electrode processes, respectively.

Fig. 8 shows the Arrhenius plots of the overall conductivity for Sr₂Al₀ and Sr₂Al_{0.06} samples under different atmospheres. It is quite evident from these impedance spectroscopy measurements that these materials present a significant proton contribution to the overall conductivity, with conductivities always higher in wet atmospheres in the whole temperature range studied, indicating that a minor fraction of water is retained in the structure even at high temperature, which is supported by the thermogravimetric analysis.

It is also very clear that there is an increase of the conductivity on doping, for instance, the total conductivity values in wet N₂ for Sr₂Al₀ and Sr₂Al_{0.06} at 800 °C are 4.3 10⁻⁶ and 2.5 10⁻⁵ S cm⁻¹, respectively, a clear indication of the generation of oxide vacancies derived from the substitution of Si⁴⁺ by Al³⁺ (see Table 3). It should also be commented that the conductivity of these samples is higher than those reported for analogous Ca₂SiO₄ [19]. For instance, the conductivity of Ca₂Al_{0.028} were 2.7 10⁻⁵ and 2.5 10⁻⁶ S cm⁻¹ in wet O₂ and dry N₂, respectively, at 900 °C, and the values for Sr₂Al_{0.028} were 5.5 10⁻⁵ and 1.7 10⁻⁵ S cm⁻¹, respectively, in the same conditions. In addition, all the

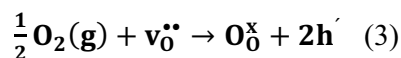
samples measured in dry atmospheres present activation energies ranging between 1.2 and 1.3 eV, typical values for oxide ion conductivity. In the case of wet atmospheres the activation energy decreases to values ranging between 1.0 and 1.1 eV, as expected when a proton is implicated in the conductivity with better mobility and consequently lower activation energy.

The increase in the conductivity on Al doping can be correlated with the increase in the O vacancy concentration, which will lead to an increase in the oxide ion conductivity, along with a protonic contribution to the conductivity in wet atmospheres due to water incorporation into these oxide ion vacancies (Eq. 2). An interesting feature to consider is how these oxide ion vacancies are accommodated in the structure. Since Si is tetrahedral, the introduction of O vacancies on Al doping might at first glance be expected to give some 3 coordinate Si or Al. Such a low coordination situation might, however, be expected to be unfavourable, and rather we can propose that the O vacancies are accommodated by condensation with adjacent Si/AlO₄ units. Thus O vacancies could be accommodated by the formation of either (Si/Al)₂O₇ or (Si/Al)₃O₁₁ similar to the types of defects proposed for related La_{1-x}Ba_{1+x}GaO_{4-x/2} and La_{1-x}Ca_xNbO_{4-x/2} [7,43]. In this way the O vacancy defect is accommodated in the structure while maintaining the tetrahedral coordination. In terms of oxide ion migration, prior modelling studies of structurally related La_{1-x}Ba_{1+x}GaO_{4-x/2} suggested a conduction mechanism via a cooperative “cog-wheel –type” process involving the breaking and reforming of these units [7], and a similar mechanism is likely here. In order to confirm this mechanism, Sr_{1.94}La_{0.06}SiO_{4.03}, a composition without oxygen vacancies was prepared. For that sample it was observed a significant drop in the conductivity, $1.1 \cdot 10^{-7} \text{ S cm}^{-1}$ at 800 °C, and an important increase in the activation energy, 1.90 eV (see Fig. 8), due to the absence of oxygen vacancies and the no possibility of a “cog-wheel type” conduction mechanism.

In terms of proton conduction, the water incorporation would lead to the break up of these units to give (Si/Al)O₄H units. Proton migration would then be a combination of proton migration around the tetrahedral unit (intra-tetrahedron motion) and between adjacent tetrahedral units (inter-tetrahedron motion). Further computer modelling studies would be required to elucidate whether, as

for $\text{La}_{1-x}\text{Ba}_{1+x}\text{GaO}_{4-x/2}$, the former intra-tetrahedron motion is the rate-limiting step for long range proton motion.

It is very important to also note that the conductivities increases under oxidising conditions, likely due to a small p-type electronic contribution to the conductivity, following the next equation, Eqn.3:



Increasing the aluminium content increases the amount of holes which leads to higher p-type electronic conductivity on Al doping. This behaviour was already observed for Ca_2SiO_4 and M_3SiO_5 -based (M= Ca and Sr) materials, which also exhibit a p-type electronic contribution to the conductivity.

Conductivity measurements in varying oxygen partial pressure were also performed, confirming the presence of a p-type electronic contribution at high $p(\text{O}_2)$ values. At intermediate temperatures, a plateau region is observed and ascribed to the pure ion conduction (proton and oxide). Finally, at very low $p(\text{O}_2)$ values, a very significant decrease in the conductivity is observed. This behaviour has been previously reported in related materials [15-17, 19] and it is ascribed to a drop in the water partial pressure at low $p(\text{O}_2)$ values. This results in a sudden decrease of the proton conductivity. It is worth mentioning that at very low oxygen partial pressures, and at 800 °C, a second plateau is observed, where the only contribution to the overall conductivity is ascribed to oxide species. These results confirmed the behaviour observed in Fig. 8, where the samples still present proton conductivity at high temperature, a trend that was already observed for other electrolytes such as $\text{La}_2\text{Zr}_2\text{O}_7$, SrZrO_3 , BaCeO_3 [44-48].and related $(\text{Ca},\text{Sr})_3\text{SiO}_5$ and Ca_2SiO_4 -based materials [15-17, 19].

Oxide transference numbers were estimated by a modified electromotive force to take into consideration the electrode polarization of the samples. The measurements were carried out under ambient 5% H_2 -Ar/air gradient to confirm that the materials are ionic conductors in these conditions.

The ion transference number decreases from 0.94 to 0.81 in the 900-650 °C temperature range (Fig.9 inset). This significant lowering in the values can be assigned to a higher contribution of the proton mobility at lower temperatures. As it can be seen in Fig. 8, dry and wet measurements, for both atmospheres, became more comparable as the temperature increases, therefore, the higher proton contribution is reflected on the ion transference number as the temperature decreases. In addition, there is significant electronic p-type contribution to the conductivity that reduces the ion transference numbers at elevated temperatures (see Fig. 8 and 9).

Conclusions

Aluminium is successfully incorporated into the structure of distortium silicate, leading to the almost complete stabilization of the superstructure orthorhombic polymorph (Pna2₁ space group) at room temperature. XRPD experiments on heating showed a full transition from the monoclinic form to the orthorhombic one between RT and 250 °C. At temperatures beyond 250 °C it was observed a complete transit to a basic orthorhombic cell (Pnma). Thermogravimetric analysis and impedance spectroscopy confirmed the formation of oxide vacancies on doping due to an increase in the water uptake and the enhancement of the electrical conductivity. Measurements in wet atmospheres revealed the presence of a very significant proton contribution to the conductivity in the whole range of temperature. Electrical conductivity measurements as a function of the oxygen partial pressure displayed proton conductivity and a small p-type electronic contribution at reducing and oxidizing conditions, respectively. However, the reported conductivities are at least two orders of magnitude lower than that of YSZ but there is room for optimisation in this new family of conductors.

Acknowledgements

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Table 1. Refined unit cell parameters, polymorph content and Rietveld agreement factors for Sr₂Si_{1-x}Al_xO_{4-x/2}□_{x/2} (x=0, 0.028, 0.06 and 0.075) series.

	Sr ₂ SiO ₄	Sr ₂ Si _{0.972} Al _{0.028} O _{3.986}	Sr ₂ Si _{0.94} Al _{0.06} O _{3.97}	Sr ₂ Si _{0.925} Al _{0.075} O _{3.9625}
Monoclinic polymorph (M)				
a (Å)	5.6614(1)	5.6597(1)	5.6609(6)	5.6602(9)
b (Å)	7.0807(1)	7.0790(2)	7.0786(7)	7.0775(9)
c (Å)	9.7575(1)	9.7678(2)	9.7744(9)	9.7766(9)
β (°)	92.626(1)	92.532(2)	92.479(8)	92.511(9)
V/Z (Å³)	97.68(1)	97.74(1)	97.83(5)	97.82(1)
Amount (wt%)	87.6(1)	23.7(1)	8.0(1)	7.8(2)
Orthorhombic polymorph (O)				
a (Å)	21.2295(8)	21.2227(2)	21.2213(2)	21.2231(2)
b (Å)	9.7417(3)	9.7526(1)	9.7560(1)	9.7572(1)

c (Å)	5.6685(2)	5.6691(1)	5.6698(1)	5.6697(1)
V (Å³)	97.69(5)	97.78(2)	97.82(2)	97.84(1)
Amount (wt%)	12.4(1)	76.3(1)	92.0(1)	92.2(1)
R_{wp} (%)	11.08	11.55	12.91	14.07
R_p (%)	8.38	8.56	9.28	10.09
R_F^M/R_F^O	2.66/3.86	3.64/3.60	4.97/4.18	5.22/4.49

Table 2. Refined unit cell parameters, polymorph content and Rietveld agreement factors for Sr₂Si_{0.94}Al_{0.06}O_{3.97}□_{0.03} at different temperatures.

Temperature	30 °C	250 °C	500 °C	750 °C	900 °C
V/Z (M/O)(Å³)	97.94(5)/97.95(2)	98.77(3)	99.80(1)	100.97(1)	101.69(1)
Amount M/O (wt%)	8.0(1)/92.0(1)	100	100	100	100
R_{wp} (%)	14.25	15.25	13.46	14.45	14.65
R_p (%)	10.40	11.41	9.13	10.02	10.45
R_F^M/R_F^O	6.33/4.41	5.85	4.67	6.11	7.11

M = Monoclinic (P12₁/n1 space group)

O = Orthorhombic (Pna2₁ at T ≤ 250 °C and Pnma space groups)

Table 3. Conductivity values for Sr₂Si_{1-x}Al_xO_{4-x/2}□_{x/2} (x=0, 0.028, 0.06 and 0.075) series at 800 °C in several flowing atmospheres: dry and wet N₂ and dry and wet O₂.

	Dry N₂ (S cm⁻¹)	Wet N₂ (S cm⁻¹)	Dry O₂ (S cm⁻¹)	Wet O₂ (S cm⁻¹)
Sr₂SiO₄	1.5 10 ⁻⁶	4.3 10 ⁻⁶	1.2 10 ⁻⁵	1.6 10 ⁻⁵
Sr₂Si_{0.972}Al_{0.028}O_{3.986}	4.8 10 ⁻⁶	5.3 10 ⁻⁶	1.9 10 ⁻⁵	2.1 10 ⁻⁵
Sr₂Si_{0.94}Al_{0.06}O_{3.97}	5.1 10 ⁻⁶	2.5 10 ⁻⁵	1.3 10 ⁻⁵	3.3 10 ⁻⁵
Sr₂Si_{0.925}Al_{0.075}O_{3.9625}	2.4 10 ⁻⁶	3.0 10 ⁻⁵	3.8 10 ⁻⁵	8.3 10 ⁻⁵

Figure captions

Figure 1. Crystal structure of Sr₂SiO₄, with the SiO₄ groups shown as tetrahedra and strontium 1 and strontium 2 cations displayed as spheres, respectively. (left) Projection close to the cb plane for the monoclinic polymorph. (right) Projection close to the cb plane for the orthorhombic polymorph (Pna2₁).

Figure 2. XRPD patterns for a) Sr_2SiO_4 , b) $\text{Sr}_2\text{Si}_{0.972}\text{Al}_{0.028}\text{O}_{3.986}$, c) $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ and d) $\text{Sr}_2\text{Si}_{0.925}\text{Al}_{0.075}\text{O}_{3.9625}$ after firing at 1550 °C.

Figure 3. Observed (crosses), calculated (full line) and difference (bottom) Rietveld refined XRPD patterns for $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$. The vertical bars are the allowed Bragg reflections.

Figure 4. XRPD patterns from 30 to 900 °C for $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ in the (left) 25-29 and 29-32 °2 θ angular range. Monoclinic and orthorhombic $\text{Pna}2_1$ reflections are labelled as * and ▲, respectively.

Figure 5. TG curves, from the top to the bottom, for Sr_2SiO_4 , $\text{Sr}_2\text{Si}_{0.972}\text{Al}_{0.028}\text{O}_{3.986}$, $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ and $\text{Sr}_2\text{Si}_{0.925}\text{Al}_{0.075}\text{O}_{3.9625}$ on cooling under humidified synthetic air and for $\text{Sr}_2\text{Si}_{0.925}\text{Al}_{0.075}\text{O}_{3.9625}$ on cooling under dry synthetic air.

Figure 6. SEM micrograph of polished and thermally etched surface of sintered pellets for a) Sr_2SiO_4 , b) $\text{Sr}_2\text{Si}_{0.972}\text{Al}_{0.028}\text{O}_{3.986}$, c) $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ and d) $\text{Sr}_2\text{Si}_{0.925}\text{Al}_{0.075}\text{O}_{3.9625}$.

Figure 7. Complex impedance spectra at 800 °C for $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ under several flowing atmospheres.

Figure 8. Arrhenius plots of $\log \sigma_T$ for (left) Sr_2SiO_4 and (right) $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$ (same scale) under several flowing atmospheres and $\text{Sr}_{1.94}\text{La}_{0.06}\text{SiO}_{4.03}$ under a dry N_2 flow.

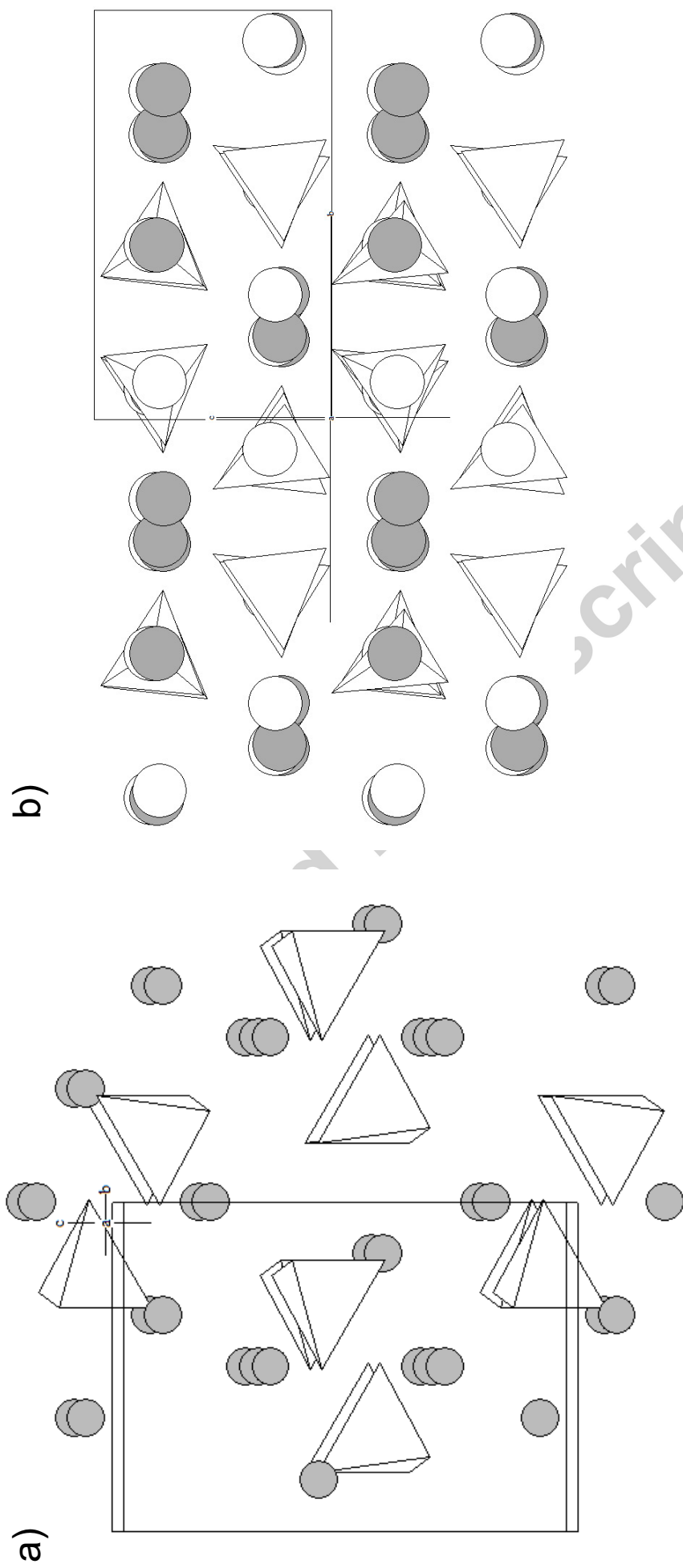
Figure 9. Overall conductivity data at 800 and 900 °C as a function of oxygen partial pressure for $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}$. The inset shows the dependence of the ion oxide transference number, t_o , for the same sample under a dry 5% H_2 -Ar/air gradient.

Table captions

Table 1. Refined unit cell parameters, polymorph content and Rietveld agreement factors for $\text{Sr}_2\text{Si}_{1-x}\text{Al}_x\text{O}_{4-x/2}\square_{x/2}$ ($x=0, 0.028, 0.06$ and 0.075) series.

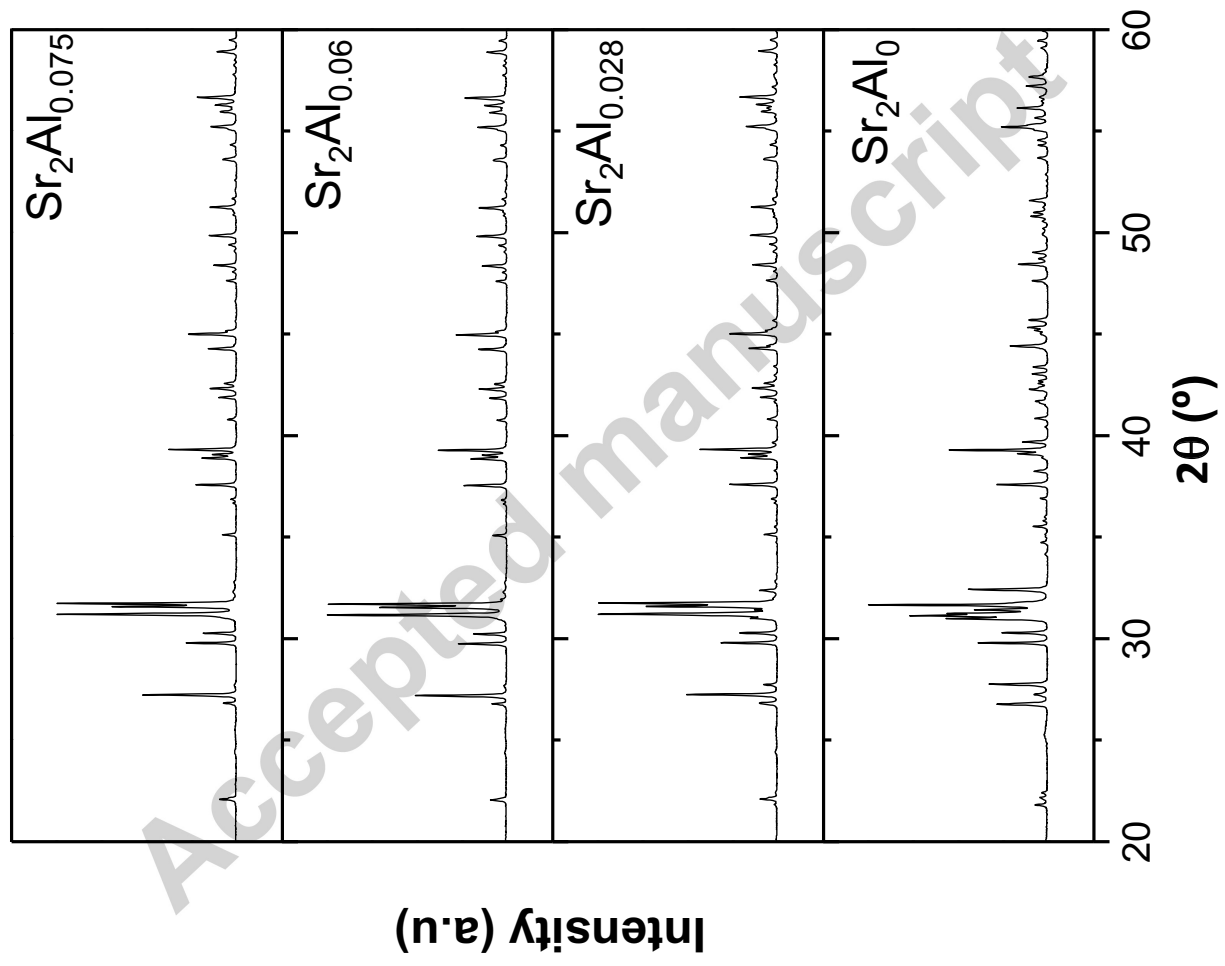
Table 2. Refined unit cell parameters, polymorph content and Rietveld agreement factors for $\text{Sr}_2\text{Si}_{0.94}\text{Al}_{0.06}\text{O}_{3.97}\square_{0.03}$ at different temperatures.

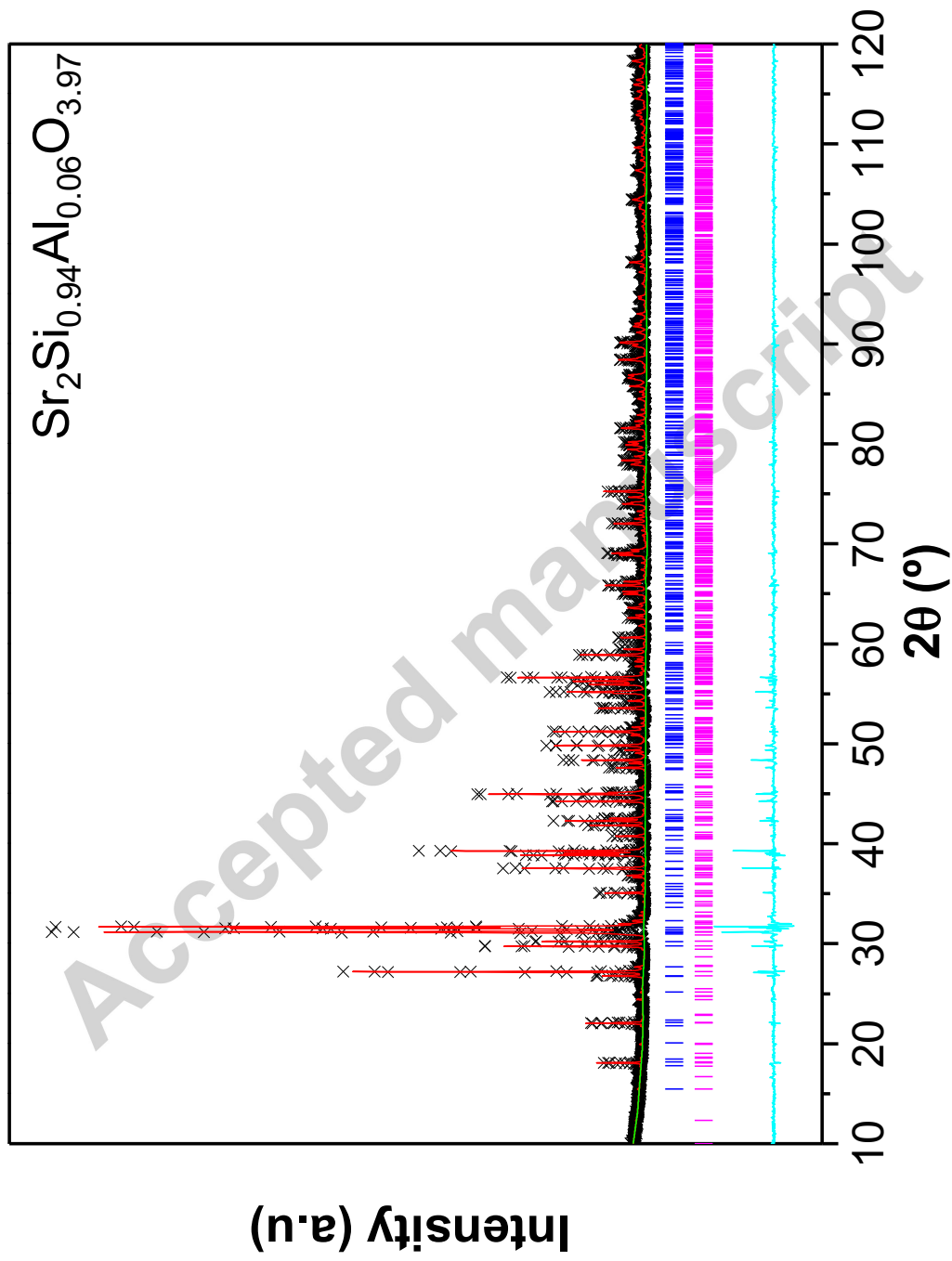
Table 3. Conductivity values for $\text{Sr}_2\text{Si}_{1-x}\text{Al}_x\text{O}_{4-x/2}\square_{x/2}$ ($x=0, 0.028, 0.06$ and 0.075) series at 800 °C in several flowing atmospheres: dry and wet N_2 and dry and wet O_2 .

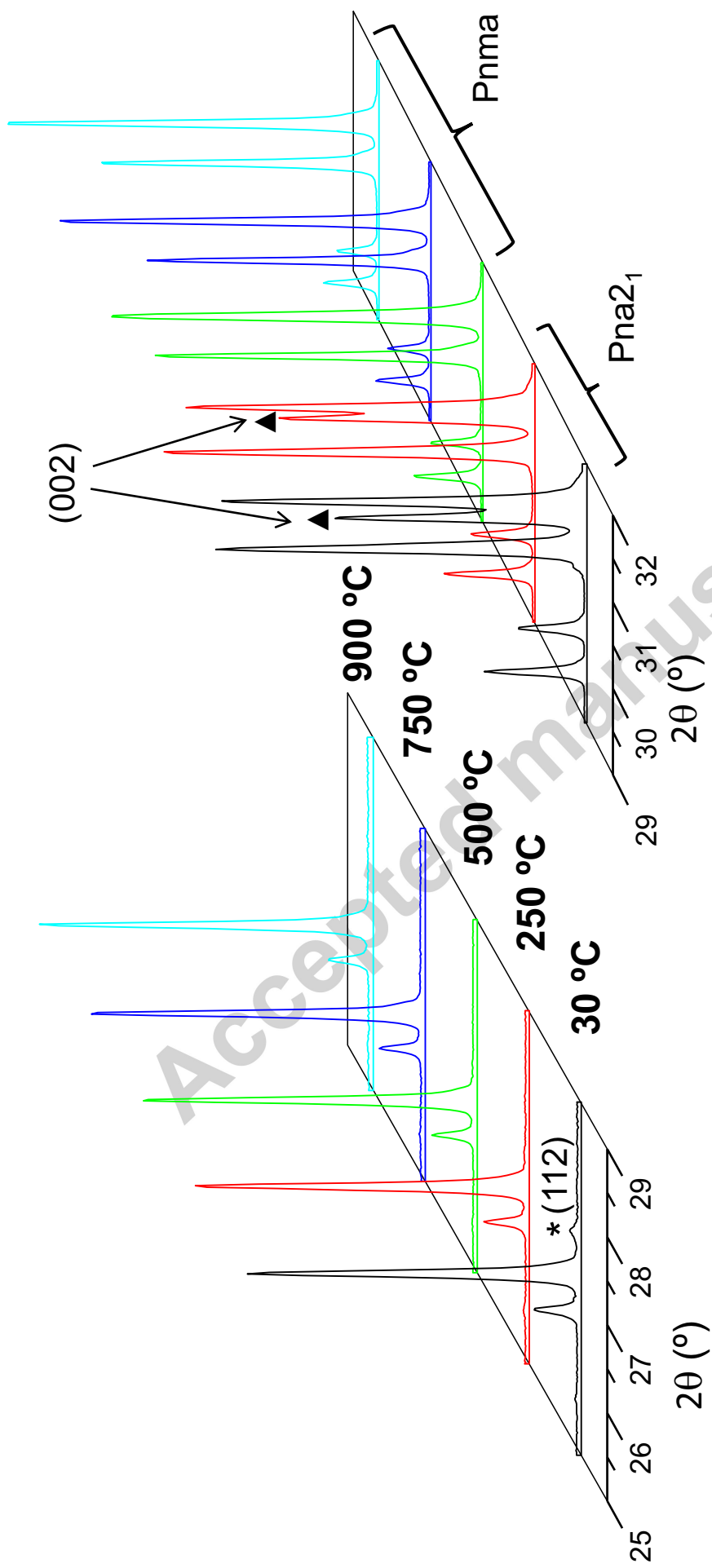


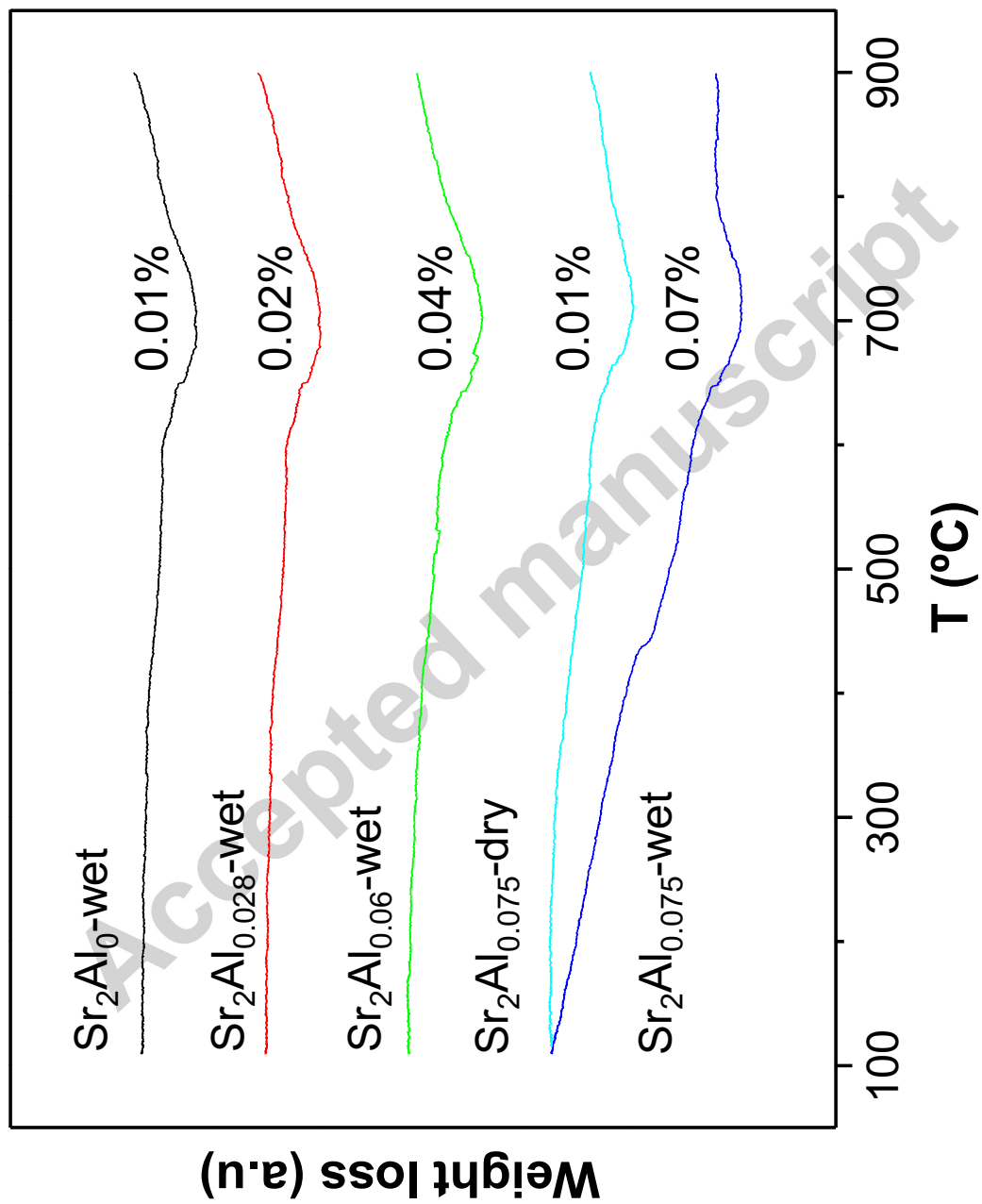
Figure

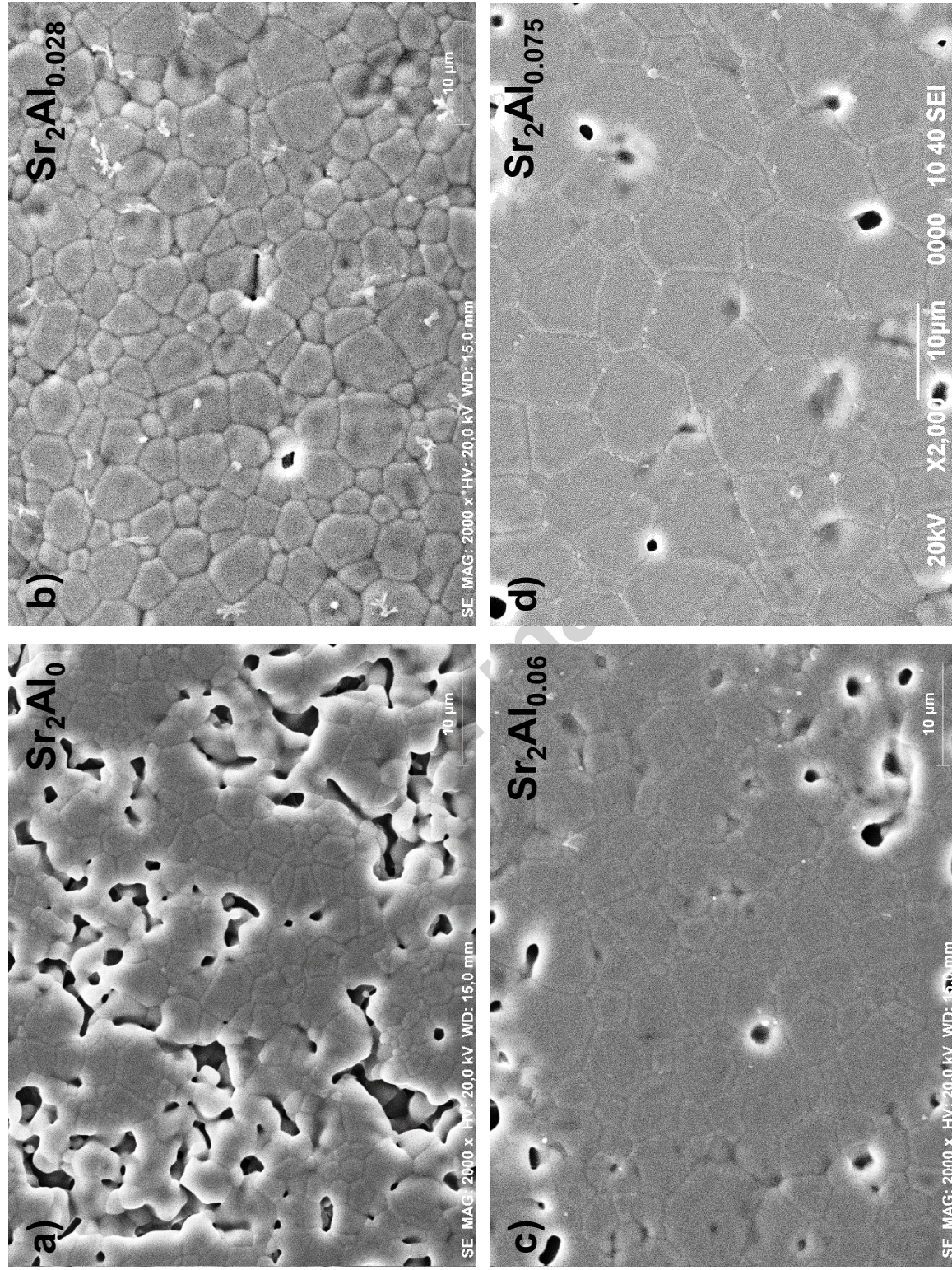
Figure 1

**Figure 2**

**Figure 3**

**Figure 4**



**Figure 6**

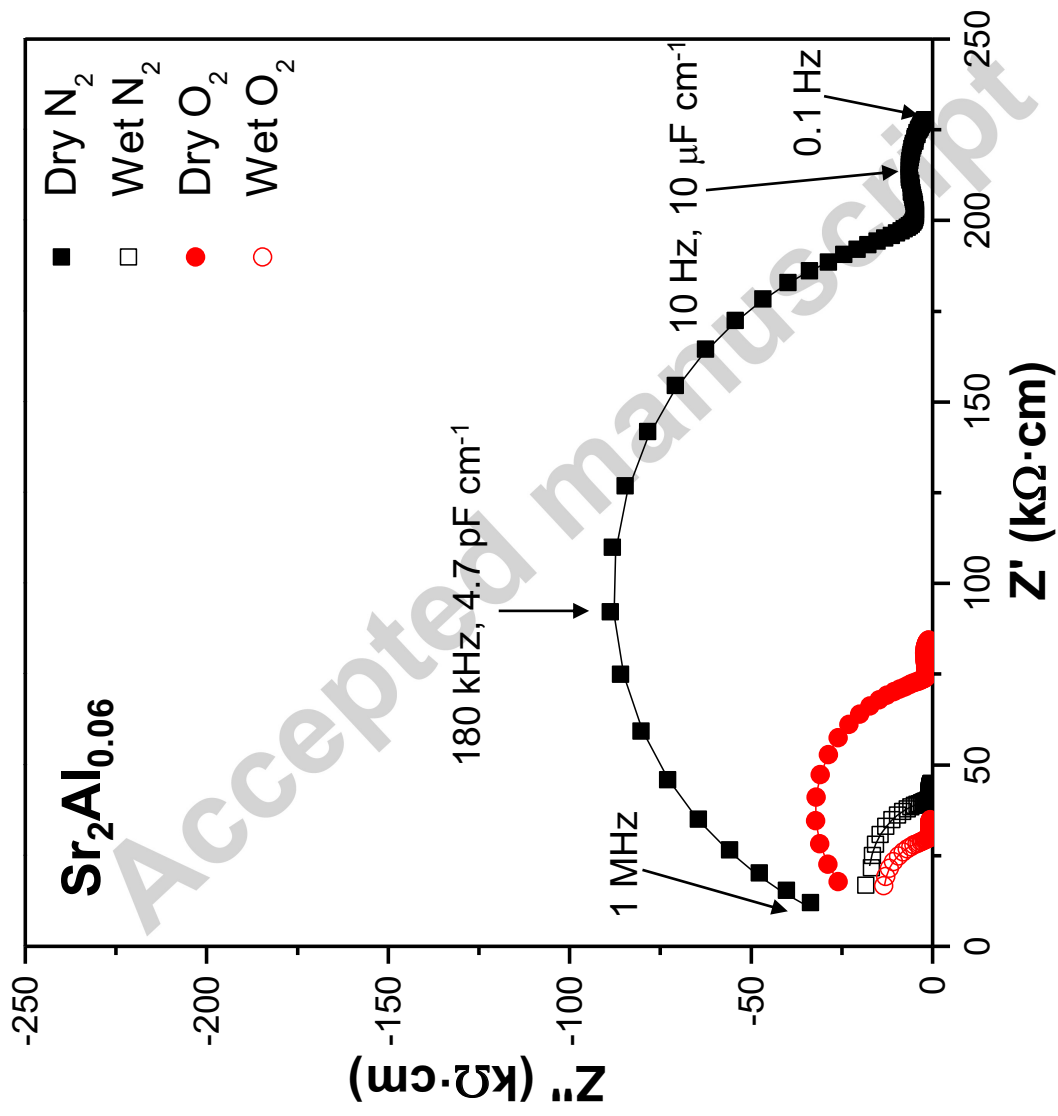


Figure 7

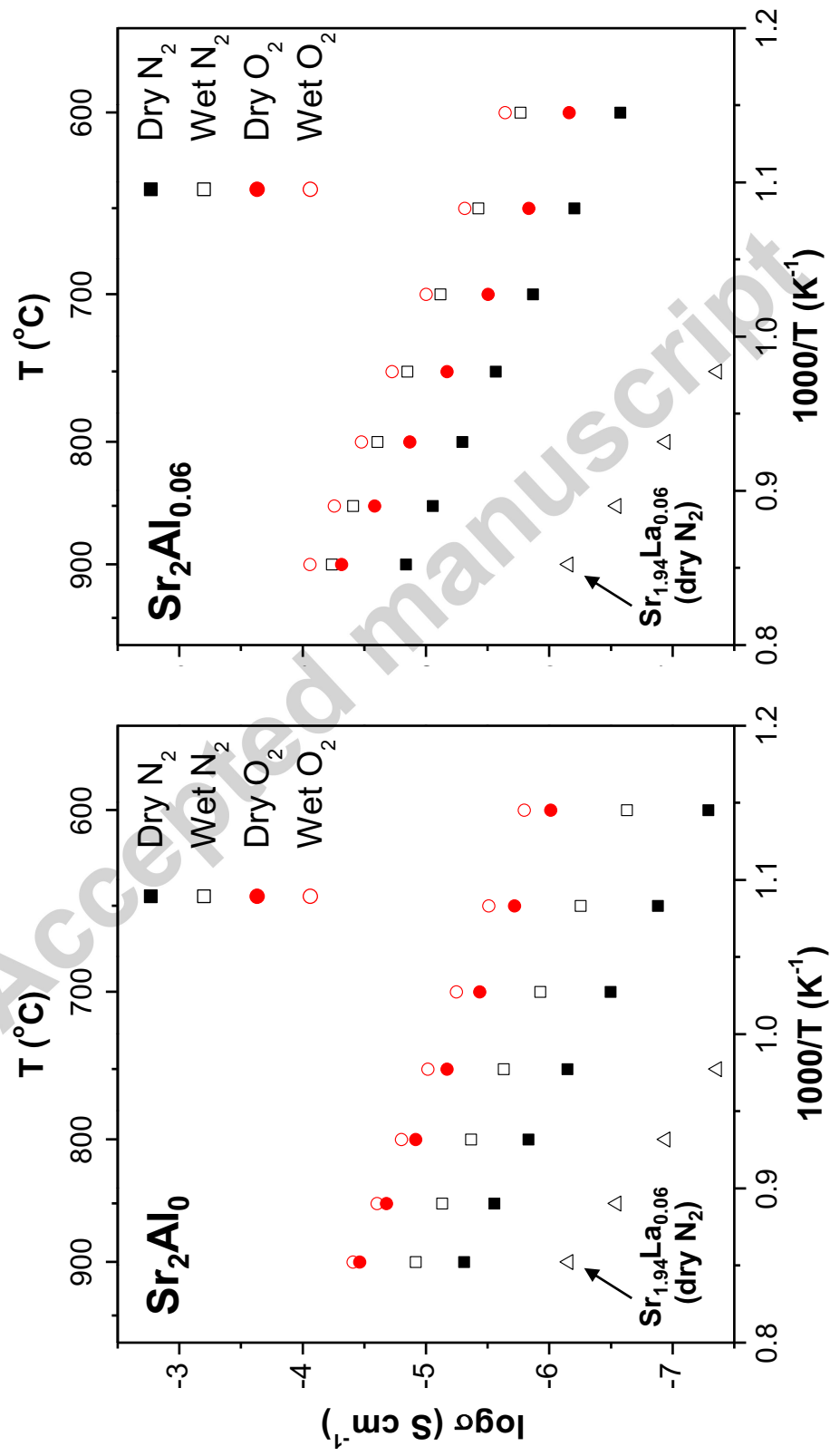
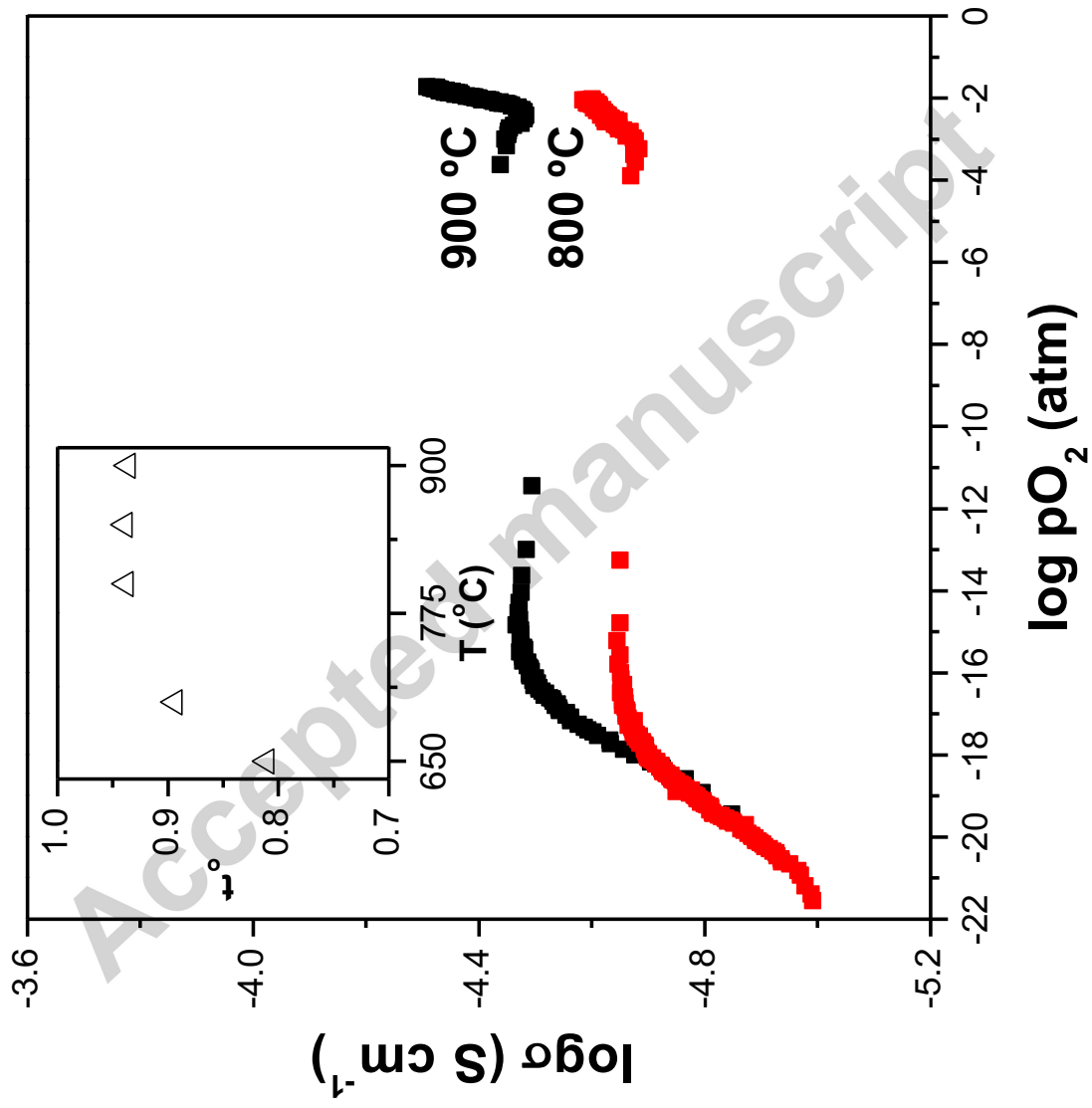


Figure 8

**Figure 9**