

Search for a light charged Higgs boson in $t \rightarrow H^\pm b$ decays, with $H^\pm \rightarrow cb$, in the lepton+jets final state in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector

ATLAS Collaboration

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Search for a light charged Higgs boson in $t \rightarrow H^\pm b$ decays, with $H^\pm \rightarrow cb$, in the lepton+jets final state in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector



The ATLAS collaboration

E-mail: atlas.publications@cern.ch

ABSTRACT: A search for a charged Higgs boson, H^\pm , produced in top-quark decays, $t \rightarrow H^\pm b$, is presented. The search targets H^\pm decays into a bottom and a charm quark, $H^\pm \rightarrow cb$. The analysis focuses on a selection enriched in top-quark pair production, where one top quark decays into a leptonically decaying W boson and a bottom quark, and the other top quark decays into a charged Higgs boson and a bottom quark. This topology leads to a lepton-plus-jets final state, characterised by an isolated electron or muon and at least four jets. The search exploits the high multiplicity of jets containing b -hadrons, and deploys a neural network classifier that uses the kinematic differences between the signal and the background. The search uses a dataset of proton-proton collisions collected at a centre-of-mass energy $\sqrt{s} = 13$ TeV between 2015 and 2018 with the ATLAS detector at CERN's Large Hadron Collider, amounting to an integrated luminosity of 139 fb^{-1} . Observed (expected) 95% confidence-level upper limits between 0.15% (0.09%) and 0.42% (0.25%) are derived for the product of branching fractions $\mathcal{B}(t \rightarrow H^\pm b) \times \mathcal{B}(H^\pm \rightarrow cb)$ for charged Higgs boson masses between 60 and 160 GeV, assuming the SM production of the top-quark pairs.

KEYWORDS: Hadron-Hadron Scattering

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1 Introduction

The Higgs boson’s discovery by the ATLAS and CMS Collaborations at the Large Hadron Collider (LHC) [1, 2] and subsequent campaigns to precisely measure its properties [3, 4] confirmed that the Standard Model (SM) of particle physics is an effective description of nature up to the TeV energy scale.

Within the SM framework, the Brout-Englert-Higgs mechanism [5–10] is responsible for generating the masses of the gauge bosons via electroweak symmetry breaking (EWSB). The Higgs boson emerges from the EWSB as the only physical spin-0 CP-even particle of the SM, while the remaining components of the complex Higgs field doublet are absorbed into the longitudinal components of the gauge bosons.

In scenarios beyond the Standard Model, the Higgs sector is typically extended to incorporate new degrees of freedom. A popular and minimal extension of the SM paradigm is provided by two-Higgs-doublet models (2HDM) [11], where the Higgs sector consists of two complex doublets, a mixture of the two doublets fulfils the same role as the SM Higgs field and generates a Higgs boson (h) similar to that in the SM, and the other mixture gives rise to a neutral CP-even Higgs boson (H), a neutral CP-odd Higgs boson (A), and a charged Higgs boson (H^\pm). Unlike the SM, a general 2HDM allows flavour changing neutral current (FCNC) interactions at tree level, which need to be suppressed. This is normally achieved by requiring that all fermions with the same electric charge couple to one Higgs doublet only [12, 13], a condition referred to as “Natural Flavour Conservation” (NFC). Depending on the assignment of up/down-type quark and lepton couplings to each Higgs doublet, 2HDMs are categorised into four different types: type-I, type-II, type-III (lepton-specific), and type-IV (flipped, also known as type-Y). A concrete realisation of a type-II 2HDM with NFC includes the Minimal Supersymmetric Standard Model (MSSM) [14, 15]. Similar Higgs sectors also arise in axion models [16].

Direct searches for new scalar particles, constraints from flavour observables, and precision measurements of the discovered Higgs boson suggest that the natural mass scale for additional Higgs bosons from 2HDMs with NFC lies above several hundred GeV, depending on the assumed model parameters (see e.g. refs. [17–20]). On the other hand, such constraints can be evaded or mitigated in other non-minimal extensions of the SM Higgs sector. A particularly rich phenomenology is expected in models with three Higgs doublets (3HDM) [21], which feature three CP-even and two CP-odd neutral Higgs bosons, as well as two charged Higgs bosons. In 3HDMs the lightest charged Higgs boson can be lighter than the top quark and can decay mainly into either a τ -lepton and a neutrino, a charm quark and a bottom quark, or a strange quark and a charm quark [22].

A search for $H^\pm \rightarrow cb$ decays¹ in top-quark decays was performed by the CMS Collaboration [23] using 19.7 fb^{-1} of proton-proton (pp) collision data collected at $\sqrt{s} = 8 \text{ TeV}$; it reported upper limits at 95% confidence level (CL) on the branching fraction $\mathcal{B}(t \rightarrow H^\pm b)$ of (0.8–0.5)%, assuming $\mathcal{B}(H^\pm \rightarrow cb) = 1.0$, for a charged Higgs boson mass (m_{H^\pm}) between 90 and 150 GeV. Related searches for $H^\pm \rightarrow cs$ in top-quark decays were performed by the ATLAS [24] and CMS [25] Collaborations, based respectively on 4.7 fb^{-1} and 19.7 fb^{-1} of pp collision data collected at the centre-of-mass energies $\sqrt{s} = 7 \text{ TeV}$ and $\sqrt{s} = 8 \text{ TeV}$; these searches reported upper limits at 95% CL on the branching fraction $\mathcal{B}(t \rightarrow H^\pm b)$ varying between 1% and 5%, assuming $\mathcal{B}(H^\pm \rightarrow cs) = 1.0$, for m_{H^\pm} between 90 and 160 GeV.

This paper presents a search for $H^\pm \rightarrow cb$ in top-quark decays which is based on a dataset of pp collisions collected at a centre-of-mass energy $\sqrt{s} = 13 \text{ TeV}$ between 2015 and 2018 with the ATLAS detector, amounting to an integrated luminosity of 139 fb^{-1} . The analysis focuses on a data sample enriched in top-quark pair production ($t\bar{t}$), where one top quark decays into a leptonically decaying W boson and a bottom quark, and the other top quark decays into a H^\pm boson and a bottom quark, as illustrated in figure 1. Compared to searches for $H^\pm \rightarrow cs$ and $H^\pm \rightarrow \tau\nu$ in $t\bar{t}$ events, searches for $H^\pm \rightarrow cb$ take

¹Charge conjugation is implied throughout the paper; the notation cb is used in place of $c\bar{b}/\bar{c}b$.

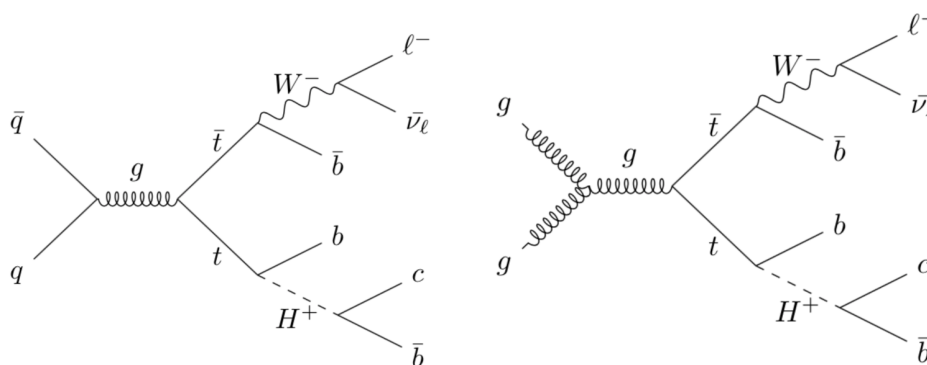


Figure 1. Illustrative leading-order Feynman diagrams for the signal considered. Charge-conjugated Feynman diagrams are also assumed.

advantage of a significantly smaller yield of the irreducible SM background originating from $t\bar{t}$ production with a W boson decaying into the Cabibbo-Kobayashi-Maskawa suppressed mode $W \rightarrow cb$. The search exploits the high multiplicity of jets containing b -hadrons (b -jets), as expected from signal events, and deploys a neural network classifier that uses the kinematic differences between the signal and the background. Thanks to a sizeably larger dataset and improved analysis techniques, this search improves the expected sensitivity to $H^\pm \rightarrow cb$ in top-quark decays by a factor of five compared to the previous publication [23] and explores an extended m_{H^\pm} range, between 60 and 160 GeV.

2 ATLAS detector

The ATLAS experiment [26] at the LHC is a multipurpose particle detector with a forward-backward symmetric cylindrical geometry and a near 4π coverage in solid angle.² It consists of an inner tracking detector surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field, electromagnetic and hadron calorimeters, and a muon spectrometer. The inner tracking detector covers the pseudorapidity range $|\eta| < 2.5$. It consists of silicon pixel, silicon microstrip, and transition radiation tracking detectors. Lead/liquid-argon (LAr) sampling calorimeters provide electromagnetic (EM) energy measurements with high granularity. A steel/scintillator-tile hadron calorimeter covers the central pseudorapidity range ($|\eta| < 1.7$). The endcap and forward regions are instrumented with LAr calorimeters for both the EM and hadronic energy measurements up to $|\eta| = 4.9$. The muon spectrometer surrounds the calorimeters and is based on three large superconducting air-core toroidal magnets with eight coils each. The field integral of the toroids ranges between 2.0 and 6.0 T m across most of the detector. The muon spectrometer includes a system of precision tracking chambers and fast detectors for triggering. A two-level trigger

²ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upwards. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the z -axis. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$. Angular distance is measured in units of $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$.

system is used to select events. The first-level trigger is implemented in hardware and uses a subset of the detector information to accept events at a rate below 100 kHz. This is followed by a software-based trigger that reduces the accepted event rate to 1 kHz on average depending on the data-taking conditions. An extensive software suite [27] is used for real and simulated data reconstruction and analysis, for operation and in the trigger and data acquisition systems of the experiment.

3 Objects definition and event selection

Data from pp collisions at $\sqrt{s} = 13$ TeV were recorded by the ATLAS detector between 2015 and 2018. Only data for which all detector subsystems were operational are used [28]; this dataset corresponds to an integrated luminosity of 139 fb^{-1} [29, 30].

Events were recorded with a single-electron or single-muon trigger, with minimum thresholds for the transverse momentum (p_T) varying from 20 to 26 GeV depending on the lepton flavour and the data-taking period. The triggers with the lowest p_T thresholds included isolation requirements based on the inner tracking detector or electromagnetic calorimeter measurements [31–34].

In each event, the primary vertex is defined as the reconstructed vertex having the highest scalar sum of the squared p_T of associated tracks with $p_T \geq 0.5$ GeV.

Electrons are reconstructed from energy clusters in the electromagnetic calorimeter that are geometrically matched to a track in the inner tracking detector. Electrons are required to satisfy $|\eta| < 2.47$ and a “Tight” identification requirement using a likelihood-based method [35]; electrons are rejected if their calorimeter clusters lie within the transition region between the barrel and endcaps of the electromagnetic calorimeter, $1.37 < |\eta| < 1.52$. Muons are reconstructed from muon spectrometer tracks matched to tracks in the inner tracking detector in the pseudorapidity range $|\eta| < 2.5$. Muon identification is based on “Medium” requirements [36]. Selected electrons and muons are required to have $p_T \geq 27$ GeV.

Isolation criteria are applied to the selected electrons and muons. For electrons, the scalar sum of the transverse energy in calorimeter energy clusters within a cone of size $\Delta R = 0.2$ around the electron is required to be less than 6% of the electron p_T , excluding clusters originating from the electron itself. For muons, the scalar sum of the p_T of tracks within a fixed-size cone around the muon (excluding its associated track) must be less than 6% of the muon p_T . The track isolation cone size is $\Delta R = 0.3$ for muon $p_T < 50$ GeV and $\Delta R = 0.2$ for muon $p_T > 50$ GeV.

Jets are reconstructed from topological energy clusters in the calorimeter [37] using the particle-flow method [38], based on the anti- k_t clustering algorithm [39, 40] with a radius parameter of 0.4. The jet energy is calibrated at particle level [41], and jets are required to have $|\eta| < 2.5$ and a minimum p_T of 25 GeV. For jets with $|\eta| < 2.4$ and $p_T < 60$ GeV, those originating from additional pp collisions in the same or neighbouring bunch crossings (pile-up) are suppressed by the use of the “jet-vertex-tagger” [42].

Jets containing b -hadrons are identified with the DL1r b -tagging algorithm [43]. A jet is b -tagged if the DL1r score is above a certain threshold, referred to as an operating point

(OP). Four OPs are defined with average expected efficiencies for b -jets of 60%, 70%, 77% and 85%, as determined in simulated $t\bar{t}$ events. The DL1r b -tagging score is divided into five exclusive bins according to the OPs. The distribution obtained by ordering these five bins from higher to lower b -jet efficiency is referred to as the “pseudo-continuous” b -tagging score.

The missing transverse momentum, with magnitude E_T^{miss} , is defined as the negative vector sum of the transverse momenta of all selected and calibrated physics objects. Low-momentum tracks from the primary vertex that cannot be associated with any of the reconstructed physics objects described before are also included in the E_T^{miss} calculation [44].

A sequential overlap removal procedure is applied to ensure that the same calorimeter energy deposit or the same track is not associated with two or more different reconstructed objects, following the prescription described in ref. [45].

The events are required to have exactly one selected electron or muon that matches the lepton that fired the trigger, and at least four jets. At least two of the jets are required to pass the 60% OP b -tagging requirement and an additional jet is required to pass the 70% b -tagging OP. In order to suppress background from multijet production, additional requirements are made on E_T^{miss} as well as on the transverse mass of the lepton and E_T^{miss} system (m_T^W):³ $E_T^{\text{miss}} > 20$ GeV and $E_T^{\text{miss}} + m_T^W > 60$ GeV. The above requirements are referred to as the “preselection”.

4 Monte Carlo samples

Monte Carlo (MC) simulation samples are used to model all backgrounds as well as the H^\pm signal, and evaluate related modelling uncertainties. The main background for this search originates from $t\bar{t}$ production in association with jets, followed by smaller contributions from single-top-quark, V +jets, $t\bar{t}V$, $t\bar{t}H$, diboson and other rare processes involving the production of a top quark. Background due to non-prompt leptons is expected to be negligible, based on studies of data using multiple lepton isolation criteria [46] and analysis of low- E_T^{miss} events.

The matrix-element calculations for all generated samples use the NNPDF3.0NLO [47] set of parton distribution functions (PDFs), unless stated otherwise. In all samples interfaced to PYTHIA 8 [48] or HERWIG 7 [49, 50], the decays of bottom and charm hadrons were simulated using the EVTGEN 1.2.0 program [51]. PYTHIA 8.230 [52] modelled the parton shower, hadronisation, and underlying event, with parameters set according to the A14 tune [53] and using the NNPDF2.3LO [47] PDF set; HERWIG 7.04 used the H7UE tune and the MMHT2014LO [54] PDF set; and HERWIG 7.13 used the HERWIG 7.1 default set of tuned parameters and the same PDF set. The effect of pile-up was modelled by overlaying each simulated hard-scatter event with inelastic pp events generated with PYTHIA 8.186 using the NNPDF2.3LO [55] PDF set and the A3 tune [56].

³ $m_T^W = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos \Delta\phi)}$, where p_T^ℓ is the transverse momentum (energy) of the muon (electron) and $\Delta\phi$ is the azimuthal angle separation between the lepton and the direction of the missing transverse momentum.

The generated events were processed through either a simulation [57] of the ATLAS detector geometry and response using GEANT4 [58] or a faster simulation, where the full GEANT4 simulation of the calorimeter response is replaced by a detailed parameterisation of the shower shapes [59]. Simulated events were processed through the same reconstruction software as the data, and corrections were applied so that the object identification efficiencies, energy scales and energy resolutions matched those determined from data control samples.

4.1 $t\bar{t}$ background simulation

The production of $t\bar{t} \rightarrow WbWb$ events (denoted simply $t\bar{t}$ in the following) was modelled using the POWHEG BOX v2 [60–63] generator at next-to-leading order (NLO) with the h_{damp} parameter⁴ set to 1.5 times the top-quark mass, m_{top} [64]. The matrix-element calculation includes diagrams with a b -quark in the initial state using the five-flavour scheme [65]. The events were interfaced to PYTHIA 8.230 [52] to model the parton shower, hadronisation, and underlying event.

The impact of using a different parton shower and hadronisation model was evaluated by comparing the nominal $t\bar{t}$ sample with another sample produced with the POWHEG BOX v2 generator interfaced to HERWIG 7.04 instead of PYTHIA 8.230.

To assess the uncertainty associated with the NLO generator, the POWHEG BOX v2 sample was compared with a sample of events generated with MADGRAPH5_AMC@NLO 2.6.0 interfaced to PYTHIA 8.230.

Analogously to similar searches performed previously in ATLAS [46, 66], the simulated $t\bar{t}$ events are categorised according to the flavour content of additional jets not originating from the decay of the $t\bar{t}$ system. Events that have at least one b -jet, excluding heavy-flavour jets from top-quark or W -boson decays, are labelled as $t\bar{t}+\geq 1b$; those with no additional b -jets but at least one charm-jet (c -jet) are labelled as $t\bar{t}+\geq 1c$; finally, events not containing any additional heavy-flavour jets are labelled as $t\bar{t}$ +light.

An additional sample to evaluate systematic uncertainties in the modelling of the $t\bar{t}+\geq 1b$ process was produced with the POWHEG BOX RES [67] generator and OPENLOOPS [68–70], using a pre-release of the implementation of this process in POWHEG BOX RES provided by its authors [71]. It was interfaced to PYTHIA 8.240. The four-flavour scheme was used with the b -quark mass set to 4.95 GeV. The factorisation scale was set to $0.5 \times \sum_{i=t,\bar{t},b,\bar{b},j} m_{T,i}$, where $m_{T,i} = (m_i^2 + p_{T,i}^2)^{1/2}$ and j denotes extra light quarks or gluons, the renormalisation scale was set to $\prod_{i=t,\bar{t},b,\bar{b}} m_{T,i}^{1/4}$, and the h_{damp} parameter was set to $0.5 \times \sum_{i=t,\bar{t},b,\bar{b}} m_{T,i}$.

All generated $t\bar{t}$ samples assume a diagonal Cabibbo-Kobayashi-Maskawa matrix, thus neglecting $t\bar{t}$ events with rare $W \rightarrow cb$ decays. Since such decays mimic the expected topology of the $H^\pm \rightarrow cb$ signal, dedicated $t\bar{t}$ samples were produced, setting $V_{cb} = 0.041$. A nominal sample was generated by using POWHEG BOX v2 at NLO with the h_{damp} parameter set to $1.5 \times m_{\text{top}}$, and using MADSPIN [72, 73] to generate $W \rightarrow cb$ decays. The

⁴The h_{damp} parameter is a resummation damping factor and one of the parameters that controls the matching of POWHEG matrix elements to the parton shower and thus effectively regulates the high- p_T radiation against which the $t\bar{t}$ system recoils.

events were interfaced to PYTHIA 8.230 to model the parton shower, hadronisation, and underlying event. Additional samples to assess the uncertainty in the parton shower and hadronisation model and the NLO generator were produced with POWHEG BOX v2 interfaced to HERWIG 7.13 and MADGRAPH5_AMC@NLO 2.6.0 interfaced to PYTHIA 8.230.

The $t\bar{t}$ sample was normalised to the cross-section prediction at next-to-next-to-leading order (NNLO) in QCD including the resummation of next-to-next-to-leading logarithmic (NNLL) soft-gluon terms calculated using TOP++ 2.0 [74–80]. For pp collisions at a centre-of-mass energy of $\sqrt{s} = 13$ TeV, this cross-section corresponds to $\sigma(t\bar{t}) = 832 \pm 51$ pb using a top-quark mass of $m_{\text{top}} = 172.5$ GeV.

4.2 Signal simulation

Samples of $t\bar{t} \rightarrow H^\pm b W^\mp b$ events were generated using the POWHEG BOX v2 generator at NLO with the h_{damp} parameter set to $1.5 \times m_{\text{top}}$, and using MADSPIN and PYTHIA 8.230 to perform the top quark and $H^\pm \rightarrow cb$ decays and to model the parton shower, hadronisation, and underlying event. The W boson was forced to decay leptonically to all three lepton flavours. A total of 11 signal MC samples were generated with m_{H^\pm} ranging from 60 to 160 GeV with 10 GeV spacing; the H^\pm boson’s total width is assumed to be 1 GeV, more than ten times smaller than the expected mass resolution. The signal samples were normalised to the same cross-section as used for the $t\bar{t}$ background sample, and assuming an arbitrary product of branching fractions $\mathcal{B}_{\text{ref}} = \mathcal{B}(t \rightarrow H^\pm b) \times \mathcal{B}(H^\pm \rightarrow cb) = 1\%$.

4.3 Other samples

The associated production of top quarks with W bosons (tW) and single-top-quark production in the t -channel and s -channel were modelled by the POWHEG BOX v2 generator at NLO in QCD using the five-flavour scheme. The tW process was modelled using the diagram removal scheme [81, 82] to handle interference and overlap with $t\bar{t}$ production. A related uncertainty was estimated by comparison with an alternative sample generated using the diagram subtraction scheme [81, 82]. The events were interfaced to PYTHIA 8.230. The uncertainty due to the parton shower and hadronisation model was evaluated by comparing the nominal samples of events with samples where events generated with POWHEG BOX v2 were interfaced to HERWIG 7.04 instead of PYTHIA 8.230. To assess the uncertainty associated with the NLO generator, the nominal samples were compared with samples generated with MADGRAPH5_AMC@NLO 2.6.2 at NLO in QCD using the five-flavour scheme and the NNPDF2.3NLO PDF set. These events were interfaced to PYTHIA 8.230.

For the tW channel single-top-quark process, the inclusive cross-section was normalised to the theory prediction calculated at NLO in QCD with NNLL soft-gluon corrections [83, 84]. The inclusive cross-section for t -channel and s -channel single-top-quark production was calculated at NLO in QCD with HATHOR 2.1 [85, 86].

The V +jets ($V = W, Z$) and the diboson (WW, WZ, ZZ) production was simulated with the SHERPA 2.2.1 or 2.2.2 generator [87] depending on the process. The simulation of V +jets used the NLO matrix elements for up to two partons, and leading-order (LO) matrix elements for up to four partons, calculated with the Comix [88] and OPENLOOPS libraries. The diboson samples, including fully leptonic final states and semileptonic final

states, where one boson decays leptonically and the other hadronically, were generated using matrix elements at NLO accuracy in QCD for up to one additional parton and at LO accuracy for up to three additional parton emissions; off-shell effects and Higgs boson contributions are accounted for, where appropriate. The calculations were matched with the SHERPA parton shower [89] using the MEPS@NLO prescription [90–93], The NNPDF3.0NNLO set of PDFs was used for the matrix-element calculation, along with the dedicated set of tuned parton-shower parameters developed by the SHERPA authors. The V +jets samples were normalised to a NNLO prediction [94].

The production of $t\bar{t}H$ events was modelled using the POWHEG BOX v2 generator at NLO. The events were interfaced to PYTHIA 8.230. The impact of using a different parton shower and hadronisation model was evaluated by showering the nominal hard-scatter events with HERWIG 7.04. To assess the uncertainty associated with the NLO generator, the nominal samples were compared with samples generated with MADGRAPH5_AMC@NLO 2.6.2 at NLO in QCD using the five-flavour scheme and the NNPDF2.3NLO PDF set; these events were interfaced to PYTHIA 8.230. The $t\bar{t}H$ cross-section was calculated at NLO QCD and NLO electroweak accuracies using MADGRAPH5_AMC@NLO as reported in ref. [95].

The production of $t\bar{t}V$ and tHq events was modelled at NLO using the MADGRAPH5_AMC@NLO 2.3.3 generator interfaced to PYTHIA 8.210.

The production of tZq events was modelled using the MADGRAPH5_AMC@NLO 2.3.3 generator at LO. The events were interfaced to PYTHIA 8.210. The tZq total cross-section was calculated at NLO using MADGRAPH5_AMC@NLO 2.3.3 with the NNPDF3.0NLO PDF set.

5 Analysis strategy

This section presents an overview of the analysis strategy developed for the $H^\pm \rightarrow cb$ search; it closely follows that of similar searches performed previously by ATLAS [66, 96].

5.1 Event categorisation

This search targets the production of a charged Higgs boson via top-quark-pair decay, $t\bar{t} \rightarrow H^\pm b W^\mp b$, followed by the decays $H^\pm \rightarrow cb$ and $W^\mp \rightarrow \ell\nu$, where ℓ denotes an electron or muon.⁵ The resulting signal event topology is characterised by four jets in the final state, three of them originating from b -quarks and one from a c -quark, which can be effectively exploited to suppress the background. Additional jets can also be present because of initial- or final-state radiation.

In order to maximise the sensitivity of the search, the preselected events are categorised into different analysis regions depending on the number of jets (4, 5 and 6) and on the number of b -tagged jets (3 and ≥ 4); they are summarised in figure 2. Events with two b -tagged jets are retained only if they have an additional b -tagged jet selected with

⁵The small additional signal acceptance from leptonic τ -lepton decays arising from the W -boson decay is also taken into account.

a looser requirement: they satisfy the 70% OP but fail the nominal (i.e. 60%) OP (denoted 2b+1bl). Therefore, a total of nine analysis regions are considered: (4j, 2b+1bl), (5j, 2b+1bl), (6j, 2b+1bl), (4j, 3b), (5j, 3b), (6j, 3b), (4j, 4b), (5j, $\geq 4b$) and (6j, $\geq 4b$), where ($n_j, m_b + k_{bl}$) indicates n selected jets, among which m and k are b -tagged with the nominal and loose b -tagging requirements, respectively; event categories with at least three b -tagged jets passing the 60% OP are used in the statistical analysis and thus are referred to as “fit regions”.

As explained in section 5.2, the 2b+1bl regions are used to derive data-based corrections to improve the modelling of the $t\bar{t}$ background. The definition of these analysis regions ensures a background composition as close as possible to what is expected in fit regions with three b -tagged jets. The main signal regions are (4j, 3b) and (5j, 3b); for $\mathcal{B} = 0.1\%$, the signal purity in these regions can reach a maximum value over the explored m_{H^\pm} range of 2.8% for the former and 1.9% for the latter. The (6j, 3b) region has lower signal purity, below 1.3% for $\mathcal{B} = 0.1\%$, and it is used mainly to constrain the background modelling uncertainties via a profile-likelihood fit (see section 7). Finally, the (5j, $\geq 4b$) and (6j, $\geq 4b$) regions offer a clean control sample to calibrate the $t\bar{t} + \geq 1b$ background, while data belonging to the (4j, 4b) region are used to recover acceptance for signal events with a c -quark misidentified as a b -jet.

Figure 3 shows the fractions of the different background components in the analysis regions. In all regions, $t\bar{t}$ production is the main source of SM background; it accounts for more than 80% of the expected background. The $t\bar{t}$ background composition depends on the jet and b -jet multiplicities. Analysis regions that include a 2b+1bl requirement have a background composition consisting of up to 62% $t\bar{t}$ +light and up to 28% $t\bar{t} + \geq 1b$, with an increasing fraction of $t\bar{t} + \geq 1b$ at higher jet multiplicity. The fit regions with a 3b requirement have higher fraction of $t\bar{t} + \geq 1b$ background, up to 61%. Most of the $t\bar{t}$ +light background events in these regions have a b -tagged c -jet from the hadronic W -boson decay, in addition to the two b -jets from the top-quark decays. The background in regions with the $\geq 4b$ requirement is mostly $t\bar{t} + \geq 1b$. The $t\bar{t} + \geq 1c$ and other backgrounds are small in all analysis regions; the $t\bar{t} + \geq 1c$ background mostly populates the 2b+1bl regions, accounting for up to 22% of the expected SM background contribution. Non- $t\bar{t}$ production contributes up to 13% in any analysis region.

5.2 $t\bar{t}$ modelling

The main background for this search originates from $t\bar{t}$ production in association with jets. It was observed that the $t\bar{t}$ simulation does not provide a fully satisfactory description of the jet multiplicity and transverse energy distributions in data; this motivates the introduction of a data-based approach to correct the $t\bar{t}$ simulation, similar to that developed in recent ATLAS searches [46, 97].

The data and the SM prediction are compared in the 2b+1bl regions separately for events with four, five or six jets. In these three analysis regions, independent corrections for the $t\bar{t}$ simulation are derived as a function of H_T^{all} , which is defined as the scalar sum of the transverse momenta of all selected objects in the event and E_T^{miss} .

N_j ↓	N_b →	2b + 1bl: exactly two b -tagged jets (60% OP) plus one loose b -tagged jet (70% OP)	3b: exactly three b -tagged jets (60% OP)	$\geq 4b$: at least four b -tagged jets (60% OP)
4j: exactly four jets		4j, 2b + 1bl (data-based $t\bar{t}$ corrections, 10 bins)	4j, 3b (signal region, 10 bins)	4j, 4b ($t\bar{t} + \geq 1b$ background control region and large S/B region, 1 bin)
5j: exactly five jets		5j, 2b + 1bl (data-based $t\bar{t}$ corrections, 10 bins)	5j, 3b (signal region, 10 bins)	5j, $\geq 4b$ ($t\bar{t} + \geq 1b$ background control region and large S/B region, 1 bin)
6j: exactly six jets		6j, 2b + 1bl (data-based $t\bar{t}$ corrections, 10 bins)	6j, 3b (signal region, shape correction for the NN discriminant in low S/B bins, 10 bins)	6j, $\geq 4b$ ($t\bar{t} + \geq 1b$ background control region, 1 bin)
		Regions used to derive $t\bar{t}$ correction	Regions for NN training and fit	

Figure 2. Summary of the analysis regions along with information about their usage, as described in the text. The b -tagging selection is based on the DL1r algorithm’s 60% efficiency OP.

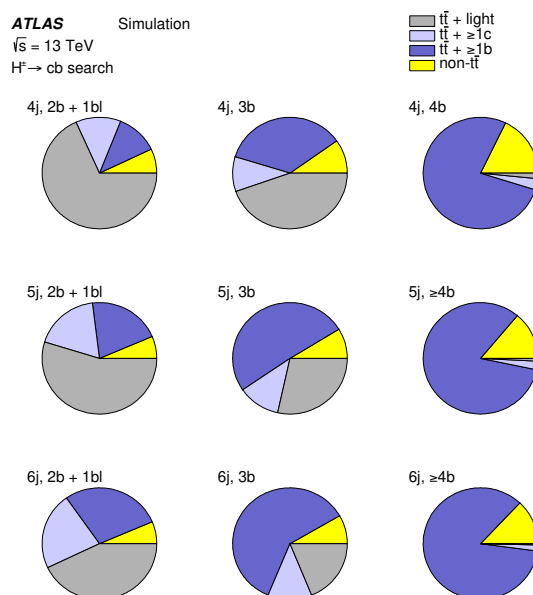


Figure 3. Fractional contributions of the various processes to the total background prediction in each analysis region. The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, W/Z +jets, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The predictions for the various background contributions are obtained through the simulation as described in section 4.

The correction factor in a given H_T^{all} bin ($H_T^{\text{all},i}$) and jet multiplicity region (j^i) is defined as

$$C(H_T^{\text{all},i}, j^i) = \frac{N^{\text{data}}(H_T^{\text{all},i}, j^i) - N^{\text{non-}t\bar{t}}(H_T^{\text{all},i}, j^i)}{N^{t\bar{t}}(H_T^{\text{all},i}, j^i)}$$

where $N^{\text{data}}(H_T^{\text{all},i}, j^i)$, $N^{\text{non-}t\bar{t}}(H_T^{\text{all},i}, j^i)$ and $N^{t\bar{t}}(H_T^{\text{all},i}, j^i)$ represent respectively the yields observed in data, and the predicted non- $t\bar{t}$ and $t\bar{t}$ yields in the H_T^{all} bin and jet multiplicity under consideration. The non- $t\bar{t}$ yields include the small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, W/Z +jets, diboson, tHq and tZq backgrounds. The fraction of subtracted non- $t\bar{t}$ background is about 7.1% in the (4j, 2b+1bl) region, 6.4% in the (5j, 2b+1bl) region, and 6.4% in the (6j, 2b+1bl) region. In all jet multiplicities, the derived corrections are close to unity for H_T^{all} above 800 GeV, and increase monotonically towards lower H_T^{all} values, reaching 1.2 for $H_T^{\text{all}} = 200$ GeV. The corrections are parameterised as a function of H_T^{all} in each jet multiplicity bin using rational functions of varying degree. A possible signal contamination in the 2b+1bl regions would mostly result in a normalisation offset for the $t\bar{t}$ prediction, smaller than 1.5%, due to the similarity of the $t\bar{t}$ background and signal H_T^{all} distribution shapes. Such an effect is fully absorbed by the systematic uncertainties of the $t\bar{t}$ prediction and has a numerically negligible impact on the signal extraction.

It was verified that after the inclusion of the data-based $t\bar{t}$ corrections there is a consistently better agreement between data and the SM prediction in all analysis regions and for a wide range of observables. Figure 4 compares the H_T^{all} distribution of data events with that of the background prediction before and after applying the data-based correction to the $t\bar{t}$ background simulation in analysis regions with three b -tagged jets that contain events not used to derive the $t\bar{t}$ corrections. An improved agreement between the background prediction and data is observed, despite the fact that the correction was derived in a region with substantially lower $t\bar{t}+\geq 1c$ and $t\bar{t}+\geq 1b$ fractions (see figure 3). This demonstrates the appropriateness of applying these corrections to all $t\bar{t}$ components ($t\bar{t}$ +light, $t\bar{t}+\geq 1c$, and $t\bar{t}+\geq 1b$) in the simulation. To account for residual differences in their modelling, the associated systematic uncertainties are treated as uncorrelated between the three $t\bar{t}$ components (see section 6.3).

5.3 Neural network discriminant

A feed-forward neural network is used to separate the H^\pm signal from the large SM background. The neural network input layer receives kinematic information about reconstructed objects, invariant masses of jet-pair permutations, and pseudo-continuous b -tagging scores, for a total of 29 input variables, summarised in table 1. The selected jets are initially sorted according to their pseudo-continuous b -tagging scores. For jets with degenerate b -tagging scores a p_T ordering is applied. After the jet ordering, the fourth jet in signal events is expected to originate from the charm-quark produced in the H^\pm decays. The neural network uses the p_T , η and ϕ of the first six sorted jets along with the b -tagging score for the fourth, fifth and sixth jets.⁶ These kinematic distributions and their correlations

⁶The input variables corresponding to the fifth (sixth) jet are set to zero in the case of events with only four (five) jets.

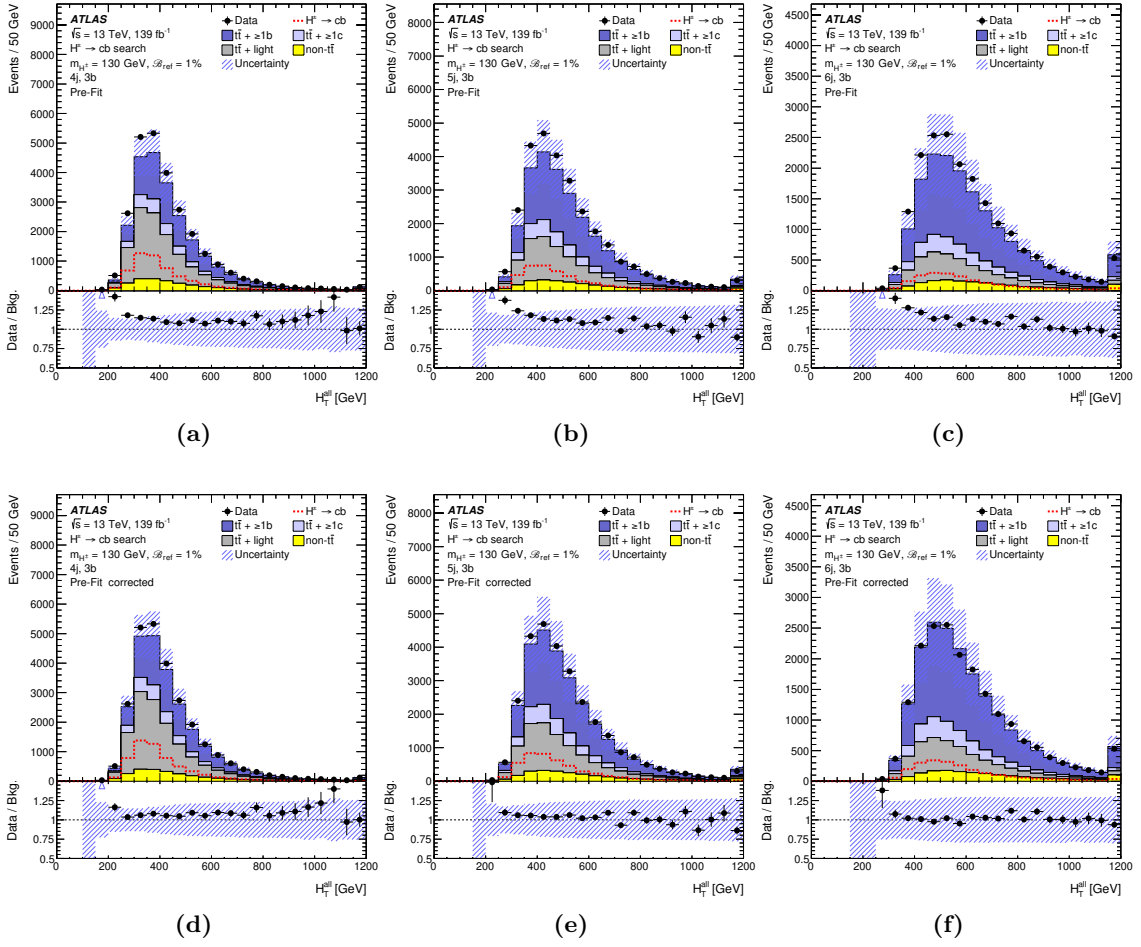


Figure 4. Comparison between the data and prediction for the H_T^{all} distribution (a-c) before and (d-f) after the inclusion of the data-based $t\bar{t}$ correction (“corrected”), in regions with three b -tagged jets prior to the likelihood fit to data (“Pre-Fit”), see section 7. The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, W/Z +jets, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The expected H^\pm signal for $m_{H^\pm} = 130$ GeV is displayed as a dashed red line normalised to $\mathcal{B}_{\text{ref}} = 1\%$. The bottom panels display the ratios of data to the SM background prediction (“Bkg.”) before the likelihood fit. The hashed area represents the total uncertainty of the background. The last bin in all figures contains the overflow.

allow the signal to be distinguished from SM background by considering the resonances produced in the event, while the b -tagging scores are effective in distinguishing the three $t\bar{t}$ background components ($t\bar{t}$ +light, $t\bar{t}+\geq 1b$ and $t\bar{t}+\geq 1c$) from the signal. The p_T , η , and ϕ of the lepton, as well as E_T^{miss} and its ϕ angle, are also included to fully characterise the event kinematics. Finally, three dijet invariant masses correlated with m_{H^\pm} are obtained by calculating the invariant mass of each of the three leading jets and the fourth jet; the inclusion of these invariant mass distributions improves the neural network’s separation of signal from background by about 20%.

Input variables	Number of variables
p_T , η , and ϕ of the first six leading jets	18
b -tagging score of the fourth, fifth, and sixth jets	3
Lepton p_T , η , and ϕ	3
Missing transverse energy and its ϕ angle	2
Invariant mass between each of the three leading jets and the fourth jet	3
Total	29

Table 1. List of input variables used in the neural network training. Jets are sorted according to their pseudo-continuous b -tagging scores, followed by their p_T , in case of having the same score.

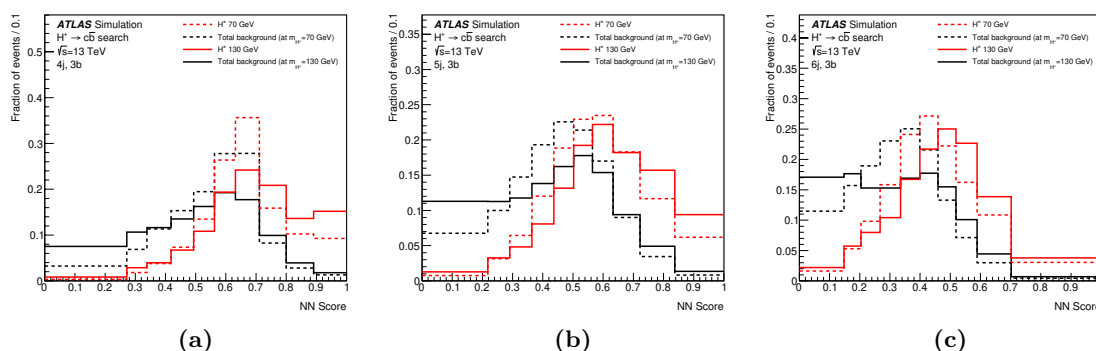


Figure 5. Comparison of the distributions of the NN score of the H^\pm signal with $m_{H^\pm} = 70$ GeV (red dashed) and $m_{H^\pm} = 130$ GeV (red solid), and the total SM background evaluated at the same masses (black dashed and black solid, respectively) in the analysis regions (a) (4j, 3b), (b) (5j, 3b), and (c) (6j, 3b). Each NN score distribution is normalised to unit area.

The neural network input layer is followed by fully connected hidden layers using rectified-linear-unit activation functions; a sigmoid function is then used in the output layer. Batch normalisation [98] is applied before each hidden layer. The training uses the Adam optimiser [99] in combination with four-fold cross-training [100], which provides orthogonal data samples for the optimisation of the neural network hyperparameters. The best hyperparameter configuration includes 2 layers and 190 neurons. The impact of overtraining is mitigated with the “dropout” [101] method in combination with the “MaxNorm” constraint.

The neural network is trained by using events with exactly four, five or six jets, and at least three b -tagged jets, corresponding to the analysis regions with highest signal purity. The training also includes the value of the m_{H^\pm} parameter, which for signal events is defined to be the true mass of the signal sample. In the case of background events, a random value of the H^\pm mass, taken from the fraction of signal masses in the input dataset, is assigned to each event [102]. In addition to increasing the size of the training sample, the use of a mass-parameterised neural network allows the different signals to be differentiated. Figure 5 compares the distributions of the neural network output score (“NN score”) of the H^\pm signal, for $m_{H^\pm} = 70$ GeV and $m_{H^\pm} = 130$ GeV, and the total SM background in the three most sensitive regions, (4j, 3b), (5j, 3b) and (6j, 3b).

6 Systematic uncertainties

Several sources of systematic uncertainty can affect the normalisation of the signal and backgrounds, as well as the shape of their corresponding NN scores distributions. Each source of systematic uncertainty is considered to be uncorrelated with the other sources. Correlations for a given systematic uncertainty are maintained across processes and event categories, unless explicitly otherwise stated. The following sections describe the considered systematic uncertainties.

6.1 Luminosity and pile-up

The uncertainty in the measurement of the integrated luminosity of the used dataset is 1.7% [29]; it affects the overall normalisation of all processes estimated from the simulation. The uncertainty is derived using the LUCID-2 detector for the baseline luminosity measurements [30], from a calibration of the luminosity scale using x - y beam-separation scans.

An uncertainty is assigned to the modelling of pile-up in simulation to account for differences between the predicted and measured inelastic cross-sections in a given fiducial volume [103].

The uncertainties in the luminosity measurement and pile-up modelling are treated as being correlated across the analysis regions and all physics processes.

6.2 Object reconstruction

Uncertainties associated with electrons and muons arise from the trigger, reconstruction, identification and isolation efficiencies, as well as the momentum scale and resolution [31, 32, 35, 36]. They are measured using data enriched in $Z \rightarrow \ell^+\ell^-$ and $J/\psi \rightarrow \ell^+\ell^-$ events ($\ell = e, \mu$). Uncertainties in jet measurements arise from the jet energy scale and resolution, and the efficiency to pass the jet-vertex-tagger requirements [42, 104]. The largest contribution comes from the jet energy scale, the uncertainty of which is split into 29 uncorrelated components, and depends on jet p_T and η , jet flavour, pile-up treatment, and simulation of the hadronic shower shape. The jet energy scale is calibrated with a series of simulation-based corrections and measurements in data samples enriched in photon or Z boson production in association with jets or in multijet production.

Uncertainties associated with energy scales and resolutions of leptons and jets are propagated to the E_T^{miss} reconstruction. Additional uncertainties affecting the reconstruction of low-energy particles present in the event, not associated with any leptons or jets, are measured in $Z \rightarrow \ell^+\ell^-$ data by studying the recoil of the Z boson [44].

Efficiencies to tag jets in the simulation are corrected to match the efficiencies measured in data by applying p_T -dependent factors. The b -jet efficiencies and c -jet mis-tagging rates are measured in a data sample enriched in $t\bar{t}$ events [105, 106], while the light-jet mis-tagging rates are measured in a multijet data sample enriched in light-jets [107]. Uncertainties affecting b -, c -, and light-jet efficiencies or mis-tagging rates are decomposed into 45, 15 and 20 uncorrelated components, respectively.

All uncertainties affecting object reconstruction are treated as being correlated across the analysis regions and all physics processes, including signal.

6.3 $t\bar{t}$ background modelling

The systematic uncertainties assigned to the $t\bar{t}$ background are designed to cover potential mismodelling of this background as a function of jet and b -jet multiplicities. Since the diagrams that contribute to $t\bar{t}+\geq 1b$, $t\bar{t}+\geq 1c$, and $t\bar{t}$ +light processes are different, their associated uncertainties are treated as being uncorrelated from each other, unless otherwise stated.

Uncertainties associated with the choice of matrix-element generator and parton shower and hadronisation models are obtained by comparing the nominal $t\bar{t}$ sample with alternative samples described in section 4. These uncertainties are evaluated in a consistent way by first putting the alternative $t\bar{t}$ samples through the same data-based correction procedure utilised to correct the nominal $t\bar{t}$ sample (section 5.2). The $t\bar{t}$ modelling uncertainties are then evaluated by comparing the alternative $t\bar{t}$ samples with the nominal samples, both sets having had the data-based correction applied. These uncertainties are also decorrelated between different jet multiplicity regions.

The uncertainty due to initial- and final-state radiation (ISR/FSR) was estimated by varying the parameters of the A14 parton shower tune [108] as described in ref. [97]. Uncertainties accounting for missing higher-order QCD corrections in the matrix-element calculation are estimated by varying the renormalisation and factorisation scales in POWHEG BOX v2 independently by factors of 2 and 0.5 relative to the nominal scales choice. Uncertainties due to higher-order QCD corrections and ISR/FSR modelling are treated as being correlated between different jet multiplicity regions.

A normalisation uncertainty of 50% is assumed separately for $t\bar{t}+\geq 1b$ and $t\bar{t}+\geq 1c$. For $t\bar{t}+\geq 1c$, the uncertainty choice is conservative given the ability to determine this background from data and the very limited sensitivity of the final results to this choice. In the case of $t\bar{t}+\geq 1b$, this choice is motivated by the observed level of disagreement between data and prediction for this background [109].

Systematic uncertainties in the data-based $t\bar{t}$ corrections arise from the statistical uncertainty in the parameterisation of the correction factors and subtraction of the non- $t\bar{t}$ backgrounds. These uncertainties are uncorrelated between each jet multiplicity but correlated across $t\bar{t}+\geq 1b$, $t\bar{t}+\geq 1c$ and $t\bar{t}$ +light background components.

The background originating from $t\bar{t}$ events with a W boson decaying into a charm quark and bottom quark was modelled with dedicated samples of simulated events. Uncertainties from the NLO generator choice, as well as from the parton shower and hadronisation models, for this subset of $t\bar{t}$ events are estimated by comparing the nominal prediction with alternative events generated as discussed in section 4. An additional cross-section uncertainty for this process is assigned by combining in quadrature a 6% uncertainty in the inclusive $t\bar{t}$ production cross-section [80] with a 3% uncertainty in the V_{cb} measurements [110].

6.4 Signal modelling

Several normalisation and shape uncertainties are taken into account for the H^\pm signal. Since the signal samples were produced with the same generator and settings as the $t\bar{t}$ background, no alternative signal samples have been generated. Instead, the uncertainties

for the $t\bar{t}$ +light background associated with the choice of matrix-element generator, parton shower and hadronisation models, as well as due to ISR/FSR, are also assigned to the signal. In addition, the uncertainty in the $t\bar{t}$ inclusive cross section is taken into account. These uncertainties are taken to be correlated with the $t\bar{t}$ +light background and uncorrelated across jet multiplicity regions. Signal modelling uncertainties have a negligible impact on the final result.

6.5 Modelling of other backgrounds

Uncertainties affecting the modelling of the single-top-quark background include a +5%/−4% uncertainty in the total cross-section, estimated as a weighted average of the theoretical uncertainties in t -, Wt - and s -channel production [83, 111, 112]. Uncertainties associated with the choice of NLO generator and parton shower and hadronisation model are evaluated by using alternative samples introduced in section 4. The uncertainty in the ISR and FSR modelling was estimated with the same procedure deployed to evaluate the corresponding source for the $t\bar{t}$ background.

Uncertainties affecting the normalisation of the W +jets and Z +jets backgrounds are estimated for the sum of both contributions (denoted V +jets). Agreement between data and the total background prediction is found to be within approximately 40% [113], which is taken to be the total normalisation uncertainty correlated across all V +jets subprocesses. An additional 25% uncertainty is added in quadrature to the inclusive 40% uncertainty for each additional jet beyond the fourth [114]; this procedure results in 47% and 52% uncertainties in regions with five or six jets, respectively.

Uncertainties in the diboson background normalisation include 5% from the NLO theory cross-sections [115]. Similarly to the V +jets background, an additional 25% normalisation uncertainty is added in quadrature for each additional jet, assuming that at leading order the diboson background contributes two jets from the decay of one of the vector bosons. Therefore, the total normalisation uncertainty is 36%, 44%, and 50% for events with four jets, five jets or six jets, respectively. These uncertainties are comparable to the level of disagreement found between the measured differential cross-section for WZ production as a function of jet multiplicity and that predicted by the simulation [116]. For both diboson and V +jets backgrounds, additional shape uncertainties are neglected compared to the large assigned normalisation uncertainties, which in turn have a negligible impact on the final result.

Modelling uncertainties for $t\bar{t}H$ production were evaluated by comparing the nominal sample with alternative samples introduced in section 4. The cross-section uncertainty for $t\bar{t}H$ production was estimated to be +9%/−12% [95].

The uncertainty in the $t\bar{t}V$ and tZq cross-sections is estimated to be 60% based on the observed level of disagreement between data and predictions [117, 118]. The uncertainty in the tHq cross-sections is conservatively assumed to be 50%. These backgrounds have negligible impact on the results.

7 Results

To test for the presence of a signal, a joint analysis of the NN score distributions in regions with three b -jets and the total yields in regions with four or more b -jets is performed. The NN score is binned in ten bins in all analysis regions with three b -jets. The statistical analysis uses a binned likelihood function $\mathcal{L}(\mu, \boldsymbol{\theta})$ constructed as a product of Poisson probability terms over all bins considered in the search. This function depends on the signal-strength parameter μ , defined as a factor multiplying the expected yield of $H^\pm \rightarrow cb$ signal events for $\mathcal{B}_{\text{ref}} = 1\%$, and $\boldsymbol{\theta}$, a set of nuisance parameters that encode the effect of systematic uncertainties on the signal and background expectations. All nuisance parameters are subject to Gaussian or log-normal constraints in the likelihood. Therefore, the expected total number of events in a given bin depends on μ and $\boldsymbol{\theta}$.

For a given value of μ , the nuisance parameters $\boldsymbol{\theta}$ allow variations of the expectations for signal and background according to the corresponding systematic uncertainties, and their fitted values result in the deviations from the nominal expectations that globally provide the best-fit to the data. This procedure allows the impact of systematic uncertainties on the search sensitivity to be reduced by taking advantage of the highly populated background-dominated bins included in the likelihood fit. Statistical uncertainties in each bin of the predicted NN score distributions are taken into account by dedicated parameters in the fit. The best-fit branching fraction is obtained by performing a binned likelihood fit to the data under the signal-plus-background hypothesis, maximising the likelihood function $\mathcal{L}(\mu, \boldsymbol{\theta})$ over μ and $\boldsymbol{\theta}$.

The fitting procedure was initially validated through extensive studies using fits to real data where bins of the NN score distributions with signal contamination above 5% or 10% were excluded (referred to as ‘blinding’ or ‘loose blinding’ requirements). In both cases, the robustness of the model for systematic uncertainties was established by verifying the stability of the fitted background when varying assumptions about some of the leading sources of uncertainty. After this, the data blinding requirements are removed and a fit under the signal-plus-background hypothesis is performed. Further checks involve a comparison of the fitted nuisance parameters before and after removal of the blinding requirements, and their values are found to be consistent. The fit results are also validated by comparing the NN score distributions in data with post-fit background predictions in the 2b+1bl regions (not directly used in the fit), and by performing an extensive comparison between data and post-fit background predictions for several kinematic variables in all considered analysis regions. The fit is found to consistently improve the modelling of all inspected observables.

The test statistic q_μ is defined as the profile likelihood ratio, $q_\mu = -2 \ln(\mathcal{L}(\mu, \hat{\boldsymbol{\theta}}_\mu) / \mathcal{L}(\hat{\mu}, \hat{\boldsymbol{\theta}}))$, where $\hat{\mu}$ and $\hat{\boldsymbol{\theta}}$ are the values of the parameters that maximise the likelihood function (subject to the constraint $0 \leq \hat{\mu} \leq \mu$), and $\hat{\boldsymbol{\theta}}_\mu$ are the values of the nuisance parameters that maximise the likelihood function for a given value of μ . The test statistic q_μ is evaluated with the RooFit package [119, 120]. A related test statistic is used to determine whether the observed data is compatible with the background-only hypothesis (the so-called discovery test) by setting $\mu = 0$ in the profile likelihood ratio and leaving $\hat{\mu}$ unconstrained: $q_0 = -2 \ln(\mathcal{L}(0, \hat{\boldsymbol{\theta}}_0) / \mathcal{L}(\hat{\mu}, \hat{\boldsymbol{\theta}}))$. The p -value (referred to as p_0), representing

the level of agreement between the data and the background-only hypothesis, is estimated by integrating the distribution of q_0 , based on the asymptotic formulae in ref. [121], above the observed value of q_0 in the data. Upper limits on μ , and thus on the branching ratio \mathcal{B} , are derived by using q_μ in the CL_s method [122, 123]. For a given signal scenario, values of \mathcal{B} yielding CL_s < 0.05, where CL_s is computed using the asymptotic approximation [121], are excluded at $\geq 95\%$ CL.

A binned likelihood fit under the signal-plus-background hypothesis is performed on the NN score distributions in the six fit regions considered. In the regions with exactly three b -tagged jets, which have the highest sensitivity, the NN score is distributed over ten bins; in the regions with at least four b -tagged jets, which have a limited number of data events or small relative signal contributions, only one bin is used. The only unconstrained parameter of the fit is the signal strength. Figures 6 and 7 show a comparison of the NN score distributions for data and prediction in the regions with exactly three or at least four b -tagged jets, respectively, before and after performing the fit to data. Tables 2 and 3 display the yields before and after performing the fit to data for all the analysis regions.

The large number of events in the fit regions, together with their different background compositions, allows the fit to place constraints on the combined effect of several sources of systematic uncertainty. As a result, an improved background prediction is obtained with a significantly reduced uncertainty.

The regions with three b -tagged jets are used to constrain the leading shape uncertainties affecting the $t\bar{t}$ +light and $t\bar{t}+\geq 1b$ background predictions, while regions with four b -tagged jets contribute by adjusting the normalisation of the $t\bar{t}+\geq 1b$ background. One of the main corrections applied by the fit is an increase of the $t\bar{t}+\geq 1b$ normalisation by about a factor of 1.2 ± 0.2 relative to the nominal prediction. This correction is in agreement with those found in previous similar ATLAS searches [46, 97] as well as in dedicated measurements of $t\bar{t}+\geq 1b$ production [124, 125]. Additionally, a few other nuisance parameters are adjusted by the fit (typically by 0.5 standard deviations or less), with the largest effect across all performed fits being on the $t\bar{t}+\geq 1c$ background normalisation, which is increased by a factor of 1.5 ± 0.5 . Some of the largest $t\bar{t}+\geq 1b$ modelling uncertainties (different parton shower and hadronisation models, and comparison to a NLO prediction in four-flavour scheme) are also significantly constrained by the fit, depending on the m_{H^\pm} hypothesis being tested. These uncertainties are reduced to about one third of their original values, owing to the large statistics of $t\bar{t}+\geq 1b$ events available in the analysis regions. Beyond the constraints on a few individual uncertainties, the significant reduction of the total background uncertainty primarily derives from the anti-correlations among systematic uncertainties from different sources resulting from the fit.

The leading uncertainties affecting the signal extraction by the fit are found to be related to the c -tagging calibration (up to $\Delta\mathcal{B} \sim 0.03 \times 10^{-2}$), the calibration of light-jet mis-tagging rate (up to $\Delta\mathcal{B} \sim 0.03 \times 10^{-2}$) and the choice of $t\bar{t}$ NLO generator in fit regions with four jets (up to $\Delta\mathcal{B} \sim 0.05 \times 10^{-2}$). Other uncertainties with a sizeable impact on the signal-strength measurement include $t\bar{t}+\geq 1b$ and $t\bar{t}+\geq 1c$ normalisation uncertainties.

The best-fit \mathcal{B} varies between 0.06% and 0.19% with an absolute uncertainty ranging from about 0.13% to 0.05% depending on m_{H^\pm} . The total uncertainties of the measured

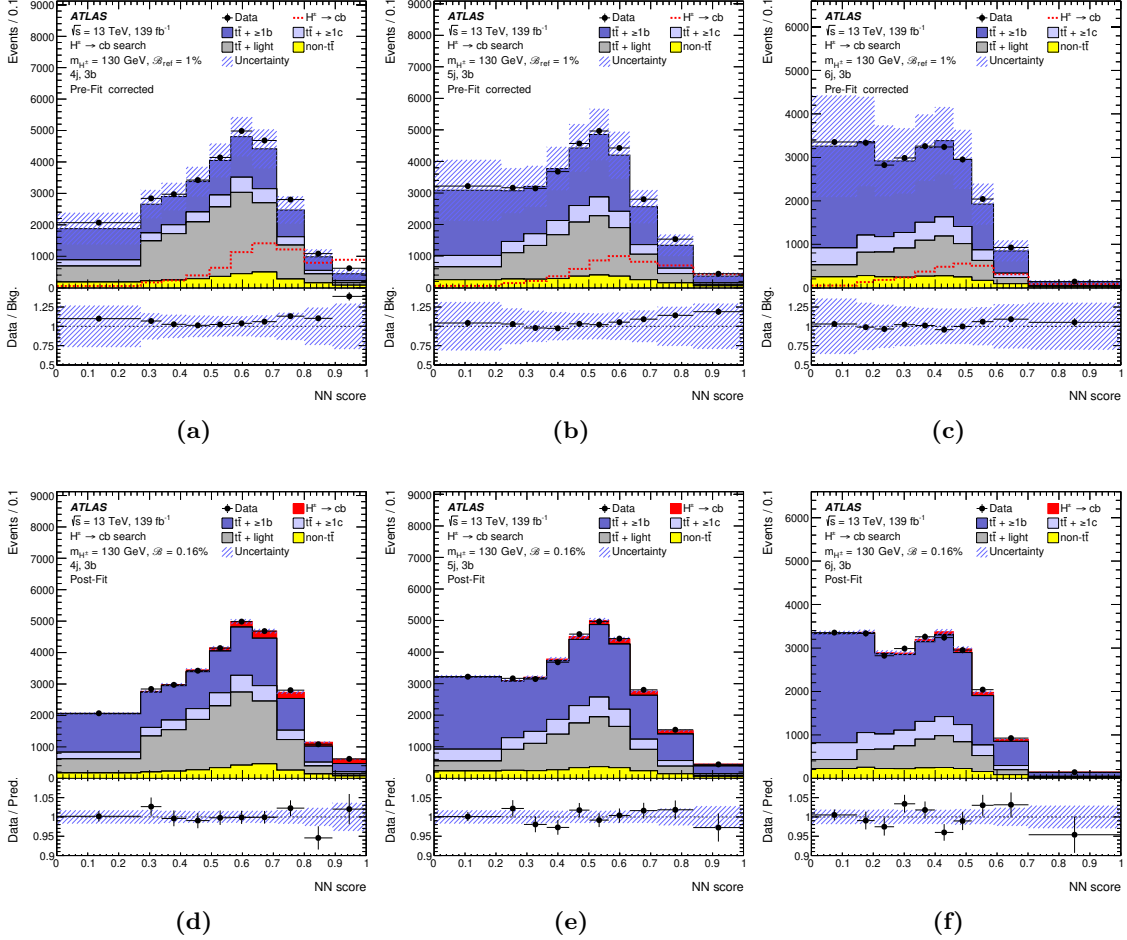


Figure 6. Comparison between the data and prediction for the NN score distributions in the fit regions (4j, 3b), (5j, 3b) and (6j, 3b) (a-c) before the likelihood fit to data (“Pre-Fit”) and (d-f) after the likelihood fit to data (“Post-Fit”). The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, W/Z +jets, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The pre-fit H^\pm signal for $m_{H^\pm} = 130$ GeV is displayed as a dashed red line normalised to $\mathcal{B}_{\text{ref}} = 1\%$. The post-fit H^\pm signal for $m_{H^\pm} = 130$ GeV is displayed as a red histogram normalised to the best-fit branching fraction of 0.16%, added on top of the background prediction. The bottom panels display the ratios of data to either the SM background prediction before the fit (“Bkg”) or the total signal-plus-background prediction after the fit (“Pred”). The hashed area represents the total uncertainty of the background. The $t\bar{t}$ background prediction is corrected according to the procedure described in section 5.2.

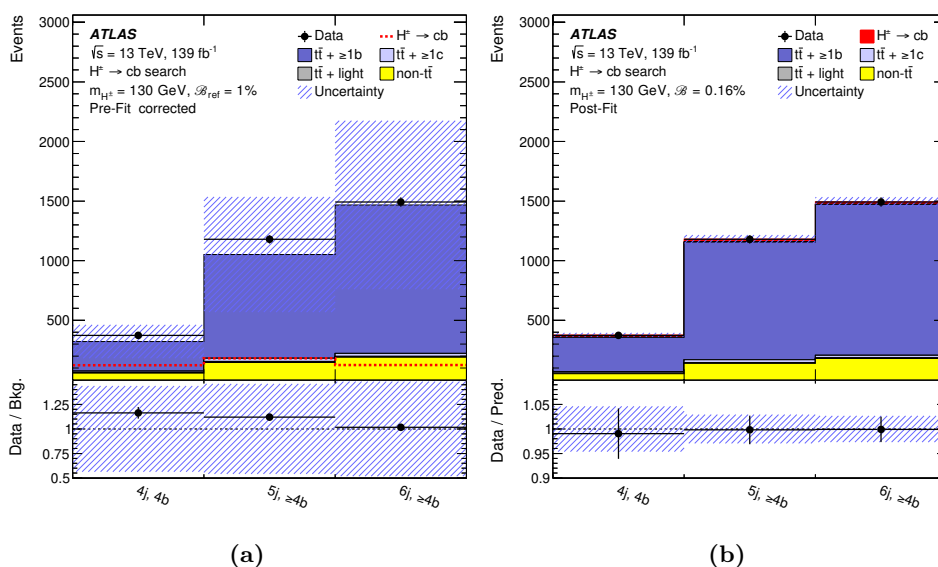


Figure 7. Comparison between the data and prediction for the event yields in the fit regions $(4_j, 4_b)$, $(5_j, \geq 4_b)$ and $(6_j, \geq 4_b)$ (a) before the likelihood fit to data (“Pre-Fit”), and (b) after the likelihood fit to data (“Post-Fit”). The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, W/Z +jets, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The pre-fit H^\pm signal for $m_{H^\pm} = 130$ GeV is displayed as a dashed red line normalised to $\mathcal{B}_{\text{ref}} = 1\%$. The post-fit H^\pm signal for $m_{H^\pm} = 130$ GeV is displayed as a red histogram normalised to the best-fit branching fraction of 0.16% , added on top of the background prediction. The bottom panels display the ratios of data to either the SM background prediction before the fit (“Bkg”) or the total signal-plus-background prediction after the fit (“Pred”). The hashed area represents the total uncertainty of the background. The $t\bar{t}$ background prediction is corrected according to the procedure described in section 5.2.

signal strengths are dominated by systematic uncertainties. The stability of the best-fit \mathcal{B} values was confirmed by separately fitting events containing either one electron or a muon, by varying the nuisance parameters most correlated with \mathcal{B} , decorrelating \mathcal{B} across jet multiplicities and by separately fitting data collected in different data-taking years.

	4j, 2b + 1bl	5j, 2b + 1bl	6j, 2b + 1bl
$H^\pm \rightarrow cb, m_{H^\pm} = 70 \text{ GeV}$	$3\,770 \pm 270$	$2\,700 \pm 220$	$1\,270 \pm 190$
$H^\pm \rightarrow cb, m_{H^\pm} = 130 \text{ GeV}$	$3\,550 \pm 270$	$2\,620 \pm 250$	$1\,310 \pm 180$
$t\bar{t} + \geq 1b$	$4\,900 \pm 2\,700$	$7\,400 \pm 3\,800$	$6\,020 \pm 3\,200$
$t\bar{t} + \geq 1c$	$5\,200 \pm 2\,600$	$6\,600 \pm 3\,400$	$4\,700 \pm 2\,500$
$t\bar{t} + \text{light}$	$27\,900 \pm 3\,600$	$19\,700 \pm 3\,700$	$9\,100 \pm 2\,400$
Non- $t\bar{t}$ backgrounds	$2\,890 \pm 610$	$2\,300 \pm 460$	$1\,370 \pm 310$
Total background (Pre-Fit corrected)	$40\,900 \pm 3\,300$	$36\,000 \pm 3\,500$	$21\,200 \pm 3\,000$
Data	40 889	35 995	21 210
	4j, 3b	5j, 3b	6j, 3b
$H^\pm \rightarrow cb, m_{H^\pm} = 70 \text{ GeV}$	$6\,600 \pm 600$	$4\,650 \pm 400$	$2\,220 \pm 280$
$H^\pm \rightarrow cb, m_{H^\pm} = 130 \text{ GeV}$	$5\,840 \pm 560$	$4\,500 \pm 410$	$2\,220 \pm 290$
$t\bar{t} + \geq 1b$	$8\,800 \pm 4\,800$	$13\,900 \pm 7\,200$	$11\,500 \pm 6\,100$
$t\bar{t} + \geq 1c$	$2\,500 \pm 1\,300$	$3\,300 \pm 1\,700$	$2\,400 \pm 1\,300$
$t\bar{t} + \text{light}$	$11\,100 \pm 1\,600$	$7\,800 \pm 1\,500$	$3\,590 \pm 940$
Non- $t\bar{t}$ backgrounds	$2\,450 \pm 470$	$2\,380 \pm 420$	$1\,590 \pm 310$
Total background (Pre-Fit corrected)	$24\,900 \pm 4\,500$	$27\,300 \pm 6\,200$	$19\,100 \pm 5\,400$
Data	26 614	28 394	19 302
	4j, 4b	5j, $\geq 4b$	6j, $\geq 4b$
$H^\pm \rightarrow cb, m_{H^\pm} = 70 \text{ GeV}$	138 ± 39	168 ± 43	119 ± 56
$H^\pm \rightarrow cb, m_{H^\pm} = 130 \text{ GeV}$	123 ± 35	181 ± 47	124 ± 59
$t\bar{t} + \geq 1b$	250 ± 140	870 ± 490	$1\,250 \pm 710$
$t\bar{t} + \geq 1c$	10 ± 7	26 ± 14	25 ± 14
$t\bar{t} + \text{light}$	5 ± 3	8 ± 6	6 ± 5
Non- $t\bar{t}$ backgrounds	57 ± 15	146 ± 26	191 ± 41
Total background (Pre-Fit corrected)	320 ± 140	$1\,050 \pm 480$	$1\,470 \pm 710$
Data	374	1 179	1 492

Table 2. Pre-fit yields in each of the analysis regions considered. The $t\bar{t}$ background yields are corrected with the procedure described in section 5.2. The contribution from $t\bar{t}$ production with rare $W \rightarrow cb$ decays is included in $t\bar{t} + \text{light}$. The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, $W/Z + \text{jets}$, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The H^\pm signal yields for $m_{H^\pm} = 70 \text{ GeV}$ and $m_{H^\pm} = 130 \text{ GeV}$ are normalised to $\mathcal{B}_{\text{ref}} = 1\%$. The quoted uncertainties are the sum in quadrature of statistical and systematic uncertainties of the yields, computed taking into account correlations among processes resulting from the data-based $t\bar{t}$ correction procedure (see section 5.2).

	4j, 3b	5j, 3b	6j, 3b
$H^\pm \rightarrow cb, m_{H^\pm} = 70 \text{ GeV}$	450 ± 380	330 ± 280	150 ± 130
$H^\pm \rightarrow cb, m_{H^\pm} = 130 \text{ GeV}$	920 ± 350	680 ± 260	330 ± 130
$t\bar{t} + \geq 1b$	10 680 ± 730	15 750 ± 910	12 340 ± 660
$t\bar{t} + \geq 1c$	2 800 ± 1 100	3 400 ± 1 400	2 350 ± 950
$t\bar{t} + \text{light}$	9 970 ± 690	6 380 ± 700	2 810 ± 410
Non- $t\bar{t}$ backgrounds	2 290 ± 410	2 200 ± 390	1 460 ± 290
Total prediction	26 620 ± 250	28 400 ± 310	19 300 ± 250
Data	26 614	28 394	19 302

	4j, 4b	5j, $\geq 4b$	6j, $\geq 4b$
$H^\pm \rightarrow cb, m_{H^\pm} = 70 \text{ GeV}$	8 ± 7	12 ± 11	7 ± 6
$H^\pm \rightarrow cb, m_{H^\pm} = 130 \text{ GeV}$	20 ± 8	25 ± 11	21 ± 11
$t\bar{t} + \geq 1b$	291 ± 23	987 ± 46	1 266 ± 57
$t\bar{t} + \geq 1c$	11 ± 7	26 ± 11	24 ± 11
$t\bar{t} + \text{light}$	3 ± 2	4 ± 4	4 ± 3
Non- $t\bar{t}$ backgrounds	52 ± 13	138 ± 26	179 ± 40
Total prediction	378 ± 17	1 182 ± 34	1 490 ± 38
Data	374	1 179	1 492

Table 3. Post-fit yields in each of the fit regions considered; the analysis regions used to derive the data-based $t\bar{t}$ corrections are not used in the fit, so their yields are not displayed. The contribution from $t\bar{t}$ production with rare $W \rightarrow cb$ decays is included in $t\bar{t} + \text{light}$. The total prediction is shown after the fit to data under the signal-plus-background hypothesis assuming H^\pm signal with $m_{H^\pm} = 130 \text{ GeV}$. The predicted yields for the H^\pm signal with $m_{H^\pm} = 70 \text{ GeV}$ are also shown for reference. The best fit-values of \mathcal{B} for H^\pm signal with $m_{H^\pm} = 130 \text{ GeV}$ and $m_{H^\pm} = 70 \text{ GeV}$ are 0.16% and 0.07% respectively. The small contributions from $t\bar{t}V$, $t\bar{t}H$, single-top-quark, $W/Z + \text{jets}$, diboson, tHq and tZq backgrounds are combined into a single background source referred to as “non- $t\bar{t}$ ”. The quoted uncertainties are the sum in quadrature of statistical and systematic uncertainties of the yields, computed taking into account correlations among nuisance parameters and among processes.

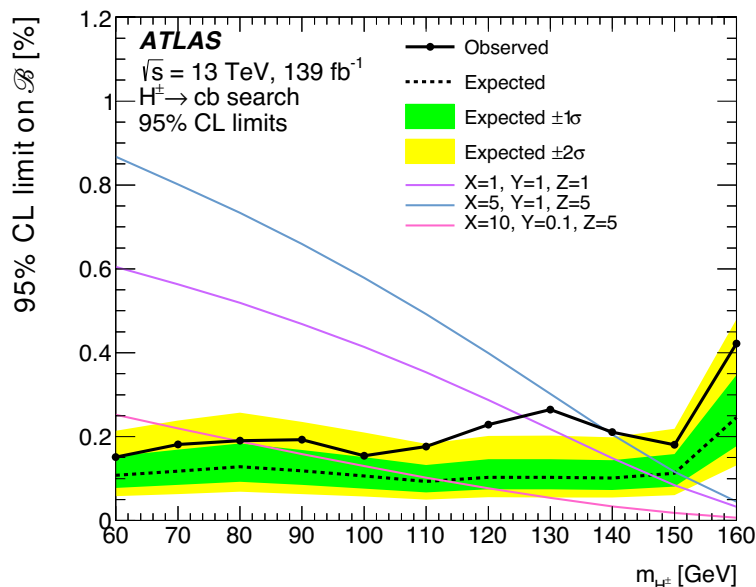


Figure 8. The observed (solid) 95% CL upper limits on $\mathcal{B} = \mathcal{B}(t \rightarrow H^\pm b) \times \mathcal{B}(H^\pm \rightarrow cb)$ as a function of m_{H^\pm} and the expectation (dashed) under the background-only hypothesis. The inner green and outer yellow shaded bands show the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties of the expected limits. The exclusion limits are presented for m_{H^\pm} between 60 and 160 GeV with 10 GeV m_{H^\pm} spacing and linear interpolation between adjacent mass points. Superimposed on the upper limits, the predictions from the 3HDM [21, 22] are shown, corresponding to three benchmark values for the parameters X , Y , and Z described in the text.

There is no significant excess of data events above the background expectation, and 95% CL limits are set on the product of branching fractions \mathcal{B} . Figure 8 shows the observed (expected) 95% CL upper limits on the branching fraction \mathcal{B} as a function of m_{H^\pm} ; they range from 0.15% (0.09%) up to 0.42% (0.25%) depending on m_{H^\pm} . The acceptance loss for the b -jet produced from the decay $t \rightarrow H^\pm b$ increases for m_{H^\pm} close to the top-quark mass, resulting in weaker exclusion limits. Superimposed on the upper limits, the predictions from the 3HDM [21, 22] are shown, corresponding to three benchmark values for the parameters X , Y , and Z , which are functions of the Higgs-doublet vacuum expectation values and the mixing angle between the charged Higgs bosons.

The observed exclusion limits are consistently weaker than the expectation. The largest excess in data has a local significance of about 3σ for $m_{H^\pm} = 130$ GeV. The corresponding best-fit \mathcal{B} is measured to be $(0.16 \pm 0.06)\%$.

Pseudo-experiments from simulated events are generated to estimate the global significance of this excess. In the first step, the correlation across all bins of the NN score distributions for all m_{H^\pm} hypotheses is estimated using the MC samples. Next, the pseudo-data are generated by fluctuating all NN score distribution bins corresponding to the pre-fit background expectation within their statistical uncertainty while preserving the correlations across all m_{H^\pm} values. Each generated pseudo-experiment is analysed with the test statistic described in section 7 built for each of the considered m_{H^\pm} hypothesis. The global p -value is then calculated as the fraction of times the maximum local significance in each

pseudo-experiment exceeds the maximum local significance observed in data. This procedure leads to a global significance of $(2.46 \pm 0.05)\sigma$ in the considered m_{H^\pm} range, with the uncertainty originating from the pseudo-experiments' sample size.

Studies using pseudo-data, defined as the sum of all predicted post-fit backgrounds plus an injected signal of variable strength, were performed to validate the behaviour of the exclusion limits as a function of m_{H^\pm} . It was concluded that a broad excess in the considered m_{H^\pm} range is consistent with the m_{H^\pm} resolution, which is significantly worse than the expected dijet mass resolution ($\sim 15\%$) due to the ambiguity in choosing the correct b -jet to pair with the fourth jet to reconstruct the m_{H^\pm} mass.

8 Conclusion

A search for the $H^\pm \rightarrow cb$ decay mode in top-quark decays is presented. The search uses a dataset of pp collisions collected at a centre-of-mass energy $\sqrt{s} = 13$ TeV between 2015 and 2018 with the ATLAS detector at CERN's Large Hadron Collider, amounting to an integrated luminosity of 139 fb^{-1} . The analysis focuses on a data sample enriched in top-quark pair production, where one top quark decays into a leptonically decaying W boson and a bottom quark, and the other top quark may decay into a H^\pm boson and a bottom quark. The search exploits the high multiplicity of b -jets, as expected from signal events, and deploys a neural network classifier that uses the kinematic differences between the signal and the background. In the absence of a significant excess of data events above the background expectation, model-independent exclusion limits at 95% confidence level on the product of branching fractions $\mathcal{B} = \mathcal{B}(t \rightarrow H^\pm b) \times \mathcal{B}(H^\pm \rightarrow cb)$ are reported as a function of m_{H^\pm} . The observed (expected) limits vary between 0.15% (0.09%) and 0.42% (0.25%) for m_{H^\pm} between 60 and 160 GeV. The largest excess in data has a significance of about 3σ for $m_{H^\pm} = 130$ GeV. In the considered m_{H^\pm} range the global significance is estimated to be about 2.5σ .

Thanks to a much larger dataset and improved analysis techniques, this search has an expected sensitivity to $H^\pm \rightarrow cb$ in top-quark decays that is a factor of five higher than achieved in a previous search by the CMS Collaboration and explores an extended m_{H^\pm} range.

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M. Battaglia [ID](#)¹³³, M. Bauce [ID](#)^{72a,72b}, F. Bauer [ID](#)^{132,*}, P. Bauer [ID](#)²², H.S. Bawa²⁹, A. Bayirli [ID](#)^{11c},
 J.B. Beacham [ID](#)⁴⁹, T. Beau [ID](#)¹²⁴, P.H. Beauchemin [ID](#)¹⁵⁵, F. Becherer [ID](#)⁵², P. Bechtle [ID](#)²²,
 H.P. Beck [ID](#)^{18,q}, K. Becker [ID](#)¹⁶³, C. Becot [ID](#)⁴⁶, A.J. Beddall [ID](#)^{11a}, V.A. Bednyakov [ID](#)³⁶,
 C.P. Bee [ID](#)¹⁴², T.A. Beermann [ID](#)³⁴, M. Begalli [ID](#)^{79b}, M. Begel [ID](#)²⁷, A. Behera [ID](#)¹⁴², J.K. Behr [ID](#)⁴⁶,
 C. Beirao Da Cruz E Silva [ID](#)³⁴, J.F. Beirer [ID](#)^{53,34}, F. Beisiegel [ID](#)²², M. Belfkir [ID](#)⁴, G. Bella [ID](#)¹⁴⁸,
 L. Bellagamba [ID](#)^{21b}, A. Bellerive [ID](#)³², P. Bellos [ID](#)¹⁹, K. Beloborodov [ID](#)³⁵, K. Belotskiy [ID](#)³⁵,
 N.L. Belyaev [ID](#)³⁵, D. Benchekroun [ID](#)^{33a}, Y. Benhammou [ID](#)¹⁴⁸, D.P. Benjamin [ID](#)²⁷, M. Benoit [ID](#)²⁷,
 J.R. Bensinger [ID](#)²⁴, S. Bentvelsen [ID](#)¹¹¹, L. Beresford [ID](#)³⁴, M. Beretta [ID](#)⁵¹, D. Berge [ID](#)¹⁷,
 E. Bergeas Kuutmann [ID](#)¹⁵⁷, N. Berger [ID](#)⁴, B. Bergmann [ID](#)¹²⁹, L.J. Bergsten [ID](#)²⁴,
 J. Beringer [ID](#)^{16a}, S. Berlendis [ID](#)⁶, G. Bernardi [ID](#)¹²⁴, C. Bernius [ID](#)¹⁴⁰, F.U. Bernlochner [ID](#)²²,
 T. Berry [ID](#)⁹², P. Berta [ID](#)¹³⁰, A. Berthold [ID](#)⁴⁸, I.A. Bertram [ID](#)⁸⁸, O. Bessidskaia Bylund [ID](#)¹⁶⁷,
 S. Bethke [ID](#)¹⁰⁷, A. Betti [ID](#)⁴², A.J. Bevan [ID](#)⁹¹, S. Bhatta [ID](#)¹⁴², D.S. Bhattacharya [ID](#)¹⁶²,
 P. Bhattarai [ID](#)²⁴, V.S. Bhopatkar [ID](#)⁵, R. Bi¹²⁶, R.M. Bianchi [ID](#)¹²⁶, O. Biebel [ID](#)¹⁰⁶, R. Bielski [ID](#)¹²⁰,
 N.V. Biesuz [ID](#)^{71a,71b}, M. Biglietti [ID](#)^{74a}, T.R.V. Billoud [ID](#)¹²⁹, M. Bindi [ID](#)⁵³, A. Bingul [ID](#)^{11d},
 C. Bini [ID](#)^{72a,72b}, S. Biondi [ID](#)^{21b,21a}, A. Biondini [ID](#)⁸⁹, C.J. Birch-sykes [ID](#)⁹⁸, G.A. Bird [ID](#)^{19,131},
 M. Birman [ID](#)¹⁶⁵, T. Bisanz [ID](#)³⁴, D. Biswas [ID](#)^{166,k}, A. Bitadze [ID](#)⁹⁸, C. Bittrich [ID](#)⁴⁸, K. Björke [ID](#)¹²²,
 I. Bloch [ID](#)⁴⁶, C. Blocker [ID](#)²⁴, A. Blue [ID](#)⁵⁷, U. Blumenschein [ID](#)⁹¹, J. Blumenthal [ID](#)⁹⁷,
 G.J. Bobbink [ID](#)¹¹¹, V.S. Bobrovnikov [ID](#)³⁵, M. Boehler [ID](#)⁵², D. Bogavac [ID](#)¹²,
 A.G. Bogdanchikov [ID](#)³⁵, C. Boehm [ID](#)^{45a}, V. Boisvert [ID](#)⁹², P. Bokan [ID](#)⁴⁶, T. Bold [ID](#)^{82a},
 M. Bomben [ID](#)¹²⁴, M. Bona [ID](#)⁹¹, M. Boonekamp [ID](#)¹³², C.D. Booth [ID](#)⁹², A.G. Borbély [ID](#)⁵⁷,
 H.M. Borecka-Bielska [ID](#)¹⁰⁵, L.S. Borgna [ID](#)⁹³, G. Borissov [ID](#)⁸⁸, D. Bortoletto [ID](#)¹²³,
 D. Boscherini [ID](#)^{21b}, M. Bosman [ID](#)¹², J.D. Bossio Sola [ID](#)³⁴, K. Bouaouda [ID](#)^{33a}, J. Boudreau [ID](#)¹²⁶,
 E.V. Bouhova-Thacker [ID](#)⁸⁸, D. Boumediene [ID](#)³⁸, R. Bouquet [ID](#)¹²⁴, A. Boveia [ID](#)¹¹⁶, J. Boyd [ID](#)³⁴,
 D. Boye [ID](#)²⁷, I.R. Boyko [ID](#)³⁶, A.J. Bozson [ID](#)⁹², J. Bracinik [ID](#)¹⁹, N. Brahimi [ID](#)^{60d,60c},
 G. Brandt [ID](#)¹⁶⁷, O. Brandt [ID](#)³⁰, F. Braren [ID](#)⁴⁶, B. Brau [ID](#)¹⁰⁰, J.E. Brau [ID](#)¹²⁰,
 W.D. Breaden Madden [ID](#)⁵⁷, K. Brendlinger [ID](#)⁴⁶, R. Brenner [ID](#)¹⁶⁵, L. Brenner [ID](#)³⁴, R. Brenner [ID](#)¹⁵⁷,
 S. Bressler [ID](#)¹⁶⁵, B. Brickwedde [ID](#)⁹⁷, D.L. Briglin [ID](#)¹⁹, D. Britton [ID](#)⁵⁷, D. Britzger [ID](#)¹⁰⁷,
 I. Brock [ID](#)²², R. Brock [ID](#)¹⁰⁴, G. Brooijmans [ID](#)³⁹, W.K. Brooks [ID](#)^{134f}, E. Brost [ID](#)²⁷,
 P.A. Bruckman de Renstrom [ID](#)⁸³, B. Brüers [ID](#)⁴⁶, D. Bruncko [ID](#)^{26b,*}, A. Bruni [ID](#)^{21b}, G. Bruni [ID](#)^{21b},
 M. Bruschi [ID](#)^{21b}, N. Bruscinò [ID](#)^{72a,72b}, L. Bryngemark [ID](#)¹⁴⁰, T. Buanes [ID](#)¹⁵, Q. Buat [ID](#)¹⁴²,
 P. Buchholz [ID](#)¹³⁸, A.G. Buckley [ID](#)⁵⁷, I.A. Budagov [ID](#)^{36,*}, M.K. Bugge [ID](#)¹²², O. Bulekov [ID](#)³⁵,
 B.A. Bullard [ID](#)⁵⁹, S. Burdin [ID](#)⁸⁹, C.D. Burgard [ID](#)⁴⁶, A.M. Burger [ID](#)¹¹⁸, B. Burghgrave [ID](#)⁷,
 J.T.P. Burr [ID](#)⁴⁶, C.D. Burton [ID](#)¹⁰, J.C. Burzynski [ID](#)¹³⁹, E.L. Busch [ID](#)³⁹, V. Büscher [ID](#)⁹⁷,
 P.J. Bussey [ID](#)⁵⁷, J.M. Butler [ID](#)²³, C.M. Buttar [ID](#)⁵⁷, J.M. Butterworth [ID](#)⁹³, W. Buttinger [ID](#)¹³¹,
 C.J. Buxo Vazquez¹⁰⁴, A.R. Buzykaev [ID](#)³⁵, G. Cabras [ID](#)^{21b}, S. Cabrera Urbán [ID](#)¹⁵⁹,
 D. Caforio [ID](#)⁵⁶, H. Cai [ID](#)¹²⁶, V.M.M. Cairo [ID](#)¹⁴⁰, O. Cakir [ID](#)^{3a}, N. Calace [ID](#)³⁴, P. Calafiura [ID](#)^{16a},
 G. Calderini [ID](#)¹²⁴, P. Calfayan [ID](#)⁶⁵, G. Callea [ID](#)⁵⁷, L.P. Caloba^{79b}, D. Calvet [ID](#)³⁸, S. Calvet [ID](#)³⁸,
 T.P. Calvet [ID](#)⁹⁹, M. Calvetti [ID](#)^{71a,71b}, R. Camacho Toro [ID](#)¹²⁴, S. Camarda [ID](#)³⁴,
 D. Camarero Munoz [ID](#)⁹⁶, P. Camarri [ID](#)^{73a,73b}, M.T. Camerlingo [ID](#)^{74a,74b}, D. Cameron [ID](#)¹²²,
 C. Camincher [ID](#)¹⁶¹, M. Campanelli [ID](#)⁹³, A. Camplani [ID](#)⁴⁰, V. Canale [ID](#)^{69a,69b}, A. Canesse [ID](#)¹⁰¹,
 M. Cano Bret [ID](#)⁷⁷, J. Cantero [ID](#)¹¹⁸, Y. Cao [ID](#)¹⁵⁸, F. Capocasa [ID](#)²⁴, M. Capua [ID](#)^{41b,41a},
 A. Carbone [ID](#)^{68a,68b}, R. Cardarelli [ID](#)^{73a}, J.C.J. Cardenas [ID](#)⁷, F. Cardillo [ID](#)¹⁵⁹, T. Carli [ID](#)³⁴,
 G. Carlino [ID](#)^{69a}, B.T. Carlson [ID](#)¹²⁶, E.M. Carlson [ID](#)^{161,153a}, L. Carminati [ID](#)^{68a,68b},

M. Carnesale [ID](#)^{72a,72b}, R.M.D. Carney [ID](#)¹⁴⁰, S. Caron [ID](#)¹¹⁰, E. Carquin [ID](#)^{134f}, S. Carrá [ID](#)⁴⁶,
 G. Carratta [ID](#)^{21b,21a}, J.W.S. Carter [ID](#)¹⁵², T.M. Carter [ID](#)⁵⁰, D. Casadei [ID](#)^{31c}, M.P. Casado [ID](#)^{12,h},
 A.F. Casha [ID](#)¹⁵², E.G. Castiglia [ID](#)¹⁶⁸, F.L. Castillo [ID](#)^{61a}, L. Castillo Garcia [ID](#)¹²,
 V. Castillo Gimenez [ID](#)¹⁵⁹, N.F. Castro [ID](#)^{127a,127e}, A. Catinaccio [ID](#)³⁴, J.R. Catmore [ID](#)¹²²,
 A. Cattai [ID](#)³⁴, V. Cavaliere [ID](#)²⁷, N. Cavalli [ID](#)^{21b,21a}, V. Cavasinni [ID](#)^{71a,71b}, E. Celebi [ID](#)^{11b},
 F. Celli [ID](#)¹²³, M.S. Centonze [ID](#)^{67a,67b}, K. Cerny [ID](#)¹¹⁹, A.S. Cerqueira [ID](#)^{79a}, A. Cerri [ID](#)¹⁴³,
 L. Cerrito [ID](#)^{73a,73b}, F. Cerutti [ID](#)^{16a}, A. Cervelli [ID](#)^{21b}, S.A. Cetin [ID](#)^{11b}, Z. Chadi [ID](#)^{33a},
 D. Chakraborty [ID](#)¹¹², M. Chala [ID](#)^{127f}, J. Chan [ID](#)¹⁶⁶, W.S. Chan [ID](#)¹¹¹, W.Y. Chan [ID](#)⁸⁹,
 J.D. Chapman [ID](#)³⁰, B. Chargeishvili [ID](#)^{146b}, D.G. Charlton [ID](#)¹⁹, T.P. Charman [ID](#)⁹¹,
 M. Chatterjee [ID](#)¹⁸, S. Chekanov [ID](#)⁵, S.V. Chekulaev [ID](#)^{153a}, G.A. Chelkov [ID](#)^{36,a}, A. Chen [ID](#)¹⁰³,
 B. Chen [ID](#)¹⁴⁸, B. Chen [ID](#)¹⁶¹, C. Chen [ID](#)^{60a}, C.H. Chen [ID](#)⁷⁸, H. Chen [ID](#)^{13c}, H. Chen [ID](#)²⁷,
 J. Chen [ID](#)^{60c}, J. Chen [ID](#)²⁴, S. Chen [ID](#)¹²⁵, S.J. Chen [ID](#)^{13c}, X. Chen [ID](#)^{60c}, X. Chen [ID](#)^{13b,ae},
 Y. Chen [ID](#)^{60a}, Y-H. Chen [ID](#)⁴⁶, C.L. Cheng [ID](#)¹⁶⁶, H.C. Cheng [ID](#)^{62a}, A. Cheplakov [ID](#)³⁶,
 E. Cheremushkina [ID](#)⁴⁶, E. Cherepanova [ID](#)³⁶, R. Cherkaoui El Moursli [ID](#)^{33e}, E. Cheu [ID](#)⁶,
 K. Cheung [ID](#)⁶³, L. Chevalier [ID](#)¹³², V. Chiarella [ID](#)⁵¹, G. Chiarelli [ID](#)^{71a}, G. Chiodini [ID](#)^{67a},
 A.S. Chisholm [ID](#)¹⁹, A. Chitan [ID](#)^{25b}, Y.H. Chiu [ID](#)¹⁶¹, M.V. Chizhov [ID](#)^{36,r}, K. Choi [ID](#)¹⁰,
 A.R. Chomont [ID](#)^{72a,72b}, Y. Chou [ID](#)¹⁰⁰, E.Y.S. Chow [ID](#)¹¹¹, T. Chowdhury [ID](#)^{31f},
 L.D. Christopher [ID](#)^{31f}, M.C. Chu [ID](#)^{62a}, X. Chu [ID](#)^{13a,13d}, J. Chudoba [ID](#)¹²⁸, J.J. Chwastowski [ID](#)⁸³,
 D. Cieri [ID](#)¹⁰⁷, K.M. Ciesla [ID](#)⁸³, V. Cindro [ID](#)⁹⁰, I.A. Cioară [ID](#)^{25b}, A. Ciocio [ID](#)^{16a},
 F. Ciroto [ID](#)^{69a,69b}, Z.H. Citron [ID](#)^{165,l}, M. Citterio [ID](#)^{68a}, D.A. Ciubotaru [ID](#)^{25b}, B.M. Ciungu [ID](#)¹⁵²,
 A. Clark [ID](#)⁵⁴, P.J. Clark [ID](#)⁵⁰, J.M. Clavijo Columbie [ID](#)⁴⁶, S.E. Clawson [ID](#)⁹⁸, C. Clement [ID](#)^{45a,45b},
 L. Clissa [ID](#)^{21b,21a}, Y. Coadou [ID](#)⁹⁹, M. Cobal [ID](#)^{66a,66c}, A. Coccaro [ID](#)^{55b}, J. Cochran [ID](#)⁷⁸,
 R.F. Coelho Barrue [ID](#)^{127a}, R. Coelho Lopes De Sa [ID](#)¹⁰⁰, S. Coelli [ID](#)^{68a}, H. Cohen [ID](#)¹⁴⁸,
 A.E.C. Coimbra [ID](#)³⁴, B. Cole [ID](#)³⁹, J. Collot [ID](#)⁵⁸, P. Conde Muño [ID](#)^{127a,127g}, S.H. Connell [ID](#)^{31c},
 I.A. Connelly [ID](#)⁵⁷, E.I. Conroy [ID](#)¹²³, F. Conventi [ID](#)^{69a,ag}, H.G. Cooke [ID](#)¹⁹,
 A.M. Cooper-Sarkar [ID](#)¹²³, F. Cormier [ID](#)¹⁶⁰, L.D. Corpe [ID](#)³⁴, M. Corradi [ID](#)^{72a,72b},
 E.E. Corrigan [ID](#)⁹⁵, F. Corriveau [ID](#)^{101,w}, M.J. Costa [ID](#)¹⁵⁹, F. Costanza [ID](#)⁴, D. Costanzo [ID](#)¹³⁶,
 B.M. Cote [ID](#)¹¹⁶, G. Cowan [ID](#)⁹², J.W. Cowley [ID](#)³⁰, K. Cranmer [ID](#)¹¹⁴, S. Crépe-Renaudin [ID](#)⁵⁸,
 F. Crescioli [ID](#)¹²⁴, M. Cristinziani [ID](#)¹³⁸, M. Cristoforetti [ID](#)^{75a,75b,c}, V. Croft [ID](#)¹⁵⁵,
 G. Crosetti [ID](#)^{41b,41a}, A. Cueto [ID](#)³⁴, T. Cuhadar Donszelmann [ID](#)¹⁵⁶, H. Cui [ID](#)^{13a,13d},
 A.R. Cukierman [ID](#)¹⁴⁰, W.R. Cunningham [ID](#)⁵⁷, F. Curcio [ID](#)^{41b,41a}, P. Czodrowski [ID](#)³⁴,
 M.M. Czurylo [ID](#)^{61b}, M.J. Da Cunha Sargedas De Sousa [ID](#)^{60a}, J.V. Da Fonseca Pinto [ID](#)^{79b},
 C. Da Via [ID](#)⁹⁸, W. Dabrowski [ID](#)^{82a}, T. Dado [ID](#)⁴⁷, S. Dahbi [ID](#)^{31f}, T. Dai [ID](#)¹⁰³,
 C. Dallapiccola [ID](#)¹⁰⁰, M. Dam [ID](#)⁴⁰, G. D'amen [ID](#)²⁷, V. D'Amico [ID](#)^{74a,74b}, J. Damp [ID](#)⁹⁷,
 J.R. Dandoy [ID](#)¹²⁵, M.F. Daneri [ID](#)²⁸, M. Danninger [ID](#)¹³⁹, V. Dao [ID](#)³⁴, G. Darbo [ID](#)^{55b},
 S. Darmora [ID](#)⁵, A. Dattagupta [ID](#)¹²⁰, S. D'Auria [ID](#)^{68a,68b}, C. David [ID](#)^{153b}, T. Davidek [ID](#)¹³⁰,
 D.R. Davis [ID](#)⁴⁹, B. Davis-Purcell [ID](#)³², I. Dawson [ID](#)⁹¹, K. De [ID](#)⁷, R. De Asmundis [ID](#)^{69a},
 M. De Beurs [ID](#)¹¹¹, S. De Castro [ID](#)^{21b,21a}, N. De Groot [ID](#)¹¹⁰, P. de Jong [ID](#)¹¹¹, H. De la Torre [ID](#)¹⁰⁴,
 A. De Maria [ID](#)^{13c}, D. De Pedis [ID](#)^{72a}, A. De Salvo [ID](#)^{72a}, U. De Sanctis [ID](#)^{73a,73b},
 M. De Santis [ID](#)^{73a,73b}, A. De Santo [ID](#)¹⁴³, J.B. De Vivie De Regie [ID](#)⁵⁸, D.V. Dedovich [ID](#)³⁶,
 J. Degens [ID](#)¹¹¹, A.M. Deiana [ID](#)⁴², J. Del Peso [ID](#)⁹⁶, Y. Delabat Diaz [ID](#)⁴⁶, F. Deliot [ID](#)¹³²,
 C.M. Delitzsch [ID](#)⁶, M. Della Pietra [ID](#)^{69a,69b}, D. Della Volpe [ID](#)⁵⁴, A. Dell'Acqua [ID](#)³⁴,
 L. Dell'Asta [ID](#)^{68a,68b}, M. Delmastro [ID](#)⁴, P.A. Delsart [ID](#)⁵⁸, S. Demers [ID](#)¹⁶⁸, M. Demichev [ID](#)³⁶,

S.P. Denisov [ID](#)³⁵, L. D’Eramo [ID](#)¹¹², D. Derendarz [ID](#)⁸³, J.E. Derkaoui [ID](#)^{33d}, F. Derue [ID](#)¹²⁴,
P. Dervan [ID](#)⁸⁹, K. Desch [ID](#)²², K. Dette [ID](#)¹⁵², C. Deutsch [ID](#)²², P.O. Deviveiros [ID](#)³⁴,
F.A. Di Bello [ID](#)^{72a,72b}, A. Di Ciaccio [ID](#)^{73a,73b}, L. Di Ciaccio [ID](#)⁴, A. Di Domenico [ID](#)^{72a,72b},
C. Di Donato [ID](#)^{69a,69b}, A. Di Girolamo [ID](#)³⁴, G. Di Gregorio [ID](#)^{71a,71b}, A. Di Luca [ID](#)^{75a,75b},
B. Di Micco [ID](#)^{74a,74b}, R. Di Nardo [ID](#)^{74a,74b}, C. Diaconu [ID](#)⁹⁹, F.A. Dias [ID](#)¹¹¹, T. Dias Do Vale [ID](#)^{127a},
M.A. Diaz [ID](#)^{134a,134b}, F.G. Diaz Capriles [ID](#)²², J. Dickinson [ID](#)^{16a}, M. Didenko [ID](#)¹⁵⁹,
E.B. Diehl [ID](#)¹⁰³, J. Dietrich [ID](#)¹⁷, S. Díez Cornell [ID](#)⁴⁶, C. Díez Pardos [ID](#)¹³⁸, A. Dimitrievska [ID](#)^{16a},
W. Ding [ID](#)^{13b}, J. Dingfelder [ID](#)²², I-M. Dinu [ID](#)^{25b}, S.J. Dittmeier [ID](#)^{61b}, F. Dittus [ID](#)³⁴,
F. Djama [ID](#)⁹⁹, T. Djobava [ID](#)^{146b}, J.I. Djuvsland [ID](#)¹⁵, M.A.B. Do Vale [ID](#)^{79c}, D. Dodsworth [ID](#)²⁴,
C. Doglioni [ID](#)⁹⁵, J. Dolejsi [ID](#)¹³⁰, Z. Dolezal [ID](#)¹³⁰, M. Donadelli [ID](#)^{79d}, B. Dong [ID](#)^{60c}, J. Donini [ID](#)³⁸,
A. D’Onofrio [ID](#)^{13c}, M. D’Onofrio [ID](#)⁸⁹, J. Dopke [ID](#)¹³¹, A. Doria [ID](#)^{69a}, M.T. Dova [ID](#)⁸⁷,
A.T. Doyle [ID](#)⁵⁷, E. Drechsler [ID](#)¹³⁹, E. Dreyer [ID](#)¹³⁹, T. Dreyer [ID](#)⁵³, A.S. Drobac [ID](#)¹⁵⁵, D. Du [ID](#)^{60a},
T.A. du Pree [ID](#)¹¹¹, F. Dubinin [ID](#)³⁵, M. Dubovsky [ID](#)^{26a}, A. Dubreuil [ID](#)⁵⁴, E. Duchovni [ID](#)¹⁶⁵,
G. Duckeck [ID](#)¹⁰⁶, O.A. Ducu [ID](#)^{34,25b}, D. Duda [ID](#)¹⁰⁷, A. Dudarev [ID](#)³⁴, M. D’uffizi [ID](#)⁹⁸,
L. Dufлот [ID](#)⁶⁴, M. Dührssen [ID](#)³⁴, C. Dülsen [ID](#)¹⁶⁷, A.E. Dumitriu [ID](#)^{25b}, M. Dunford [ID](#)^{61a},
S. Dungs [ID](#)⁴⁷, K. Dunne [ID](#)^{45a,45b}, A. Duperrin [ID](#)⁹⁹, H. Duran Yildiz [ID](#)^{3a}, M. Düren [ID](#)⁵⁶,
A. Durglishvili [ID](#)^{146b}, B. Dutta [ID](#)⁴⁶, B.L. Dwyer [ID](#)¹¹², G.I. Dyckes [ID](#)^{16a}, M. Dyndal [ID](#)^{82a},
S. Dysch [ID](#)⁹⁸, B.S. Dziedzic [ID](#)⁸³, B. Eckerova [ID](#)^{26a}, M.G. Eggleston [ID](#)⁴⁹,
E. Egidio Purcino De Souza [ID](#)^{79b}, L.F. Ehrke [ID](#)⁵⁴, T. Eifert [ID](#)⁷, G. Eigen [ID](#)¹⁵, K. Einsweiler [ID](#)^{16a},
T. Ekelof [ID](#)¹⁵⁷, Y. El Ghazali [ID](#)^{33b}, H. El Jarrari [ID](#)^{33e}, A. El Moussaouy [ID](#)^{33a}, V. Ellajosyula [ID](#)¹⁵⁷,
M. Ellert [ID](#)¹⁵⁷, F. Ellinghaus [ID](#)¹⁶⁷, A.A. Elliot [ID](#)⁹¹, N. Ellis [ID](#)³⁴, J. Elmsheuser [ID](#)²⁷, M. Elsing [ID](#)³⁴,
D. Emelianov [ID](#)¹³¹, A. Emerman [ID](#)³⁹, Y. Enari [ID](#)¹⁵⁰, J. Erdmann [ID](#)⁴⁷, A. Ereditato [ID](#)¹⁸,
P.A. Erland [ID](#)⁸³, M. Errenst [ID](#)¹⁶⁷, M. Escalier [ID](#)⁶⁴, C. Escobar [ID](#)¹⁵⁹, O. Estrada Pastor [ID](#)¹⁵⁹,
E. Etzion [ID](#)¹⁴⁸, G. Evans [ID](#)^{127a}, H. Evans [ID](#)⁶⁵, M.O. Evans [ID](#)¹⁴³, A. Ezhilov [ID](#)³⁵, F. Fabbri [ID](#)⁵⁷,
L. Fabbri [ID](#)^{21b,21a}, G. Facini [ID](#)¹⁶³, V. Fadeyev [ID](#)¹³³, R.M. Fakhruddinov [ID](#)³⁵, S. Falciano [ID](#)^{72a},
P.J. Falke [ID](#)²², S. Falke [ID](#)³⁴, J. Faltova [ID](#)¹³⁰, Y. Fan [ID](#)^{13a}, Y. Fang [ID](#)^{13a,13d}, G. Fanourakis [ID](#)⁴⁴,
M. Fanti [ID](#)^{68a,68b}, M. Faraj [ID](#)^{60c}, A. Farbin [ID](#)⁷, A. Farilla [ID](#)^{74a}, E.M. Farina [ID](#)^{70a,70b},
T. Farooque [ID](#)¹⁰⁴, S.M. Farrington [ID](#)⁵⁰, P. Farthouat [ID](#)³⁴, F. Fassi [ID](#)^{33e}, D. Fassouliotis [ID](#)⁸,
M. Fauci Giannelli [ID](#)^{73a,73b}, W.J. Fawcett [ID](#)³⁰, L. Fayard [ID](#)⁶⁴, O.L. Fedin [ID](#)^{35,a}, G. Fedotov [ID](#)³⁵,
M. Feickert [ID](#)¹⁵⁸, L. Feligioni [ID](#)⁹⁹, A. Fell [ID](#)¹³⁶, C. Feng [ID](#)^{60b}, M. Feng [ID](#)^{13b}, M.J. Fenton [ID](#)¹⁵⁶,
A.B. Fenyuk [ID](#)³⁵, S.W. Ferguson [ID](#)⁴³, J. Ferrando [ID](#)⁴⁶, A. Ferrari [ID](#)¹⁵⁷, P. Ferrari [ID](#)¹¹¹,
R. Ferrari [ID](#)^{70a}, D. Ferrere [ID](#)⁵⁴, C. Ferretti [ID](#)¹⁰³, F. Fiedler [ID](#)⁹⁷, A. Filipčič [ID](#)⁹⁰, F. Filthaut [ID](#)¹¹⁰,
M.C.N. Fiolhais [ID](#)^{127a,127c,b}, L. Fiorini [ID](#)¹⁵⁹, F. Fischer [ID](#)¹³⁸, W.C. Fisher [ID](#)¹⁰⁴, T. Fitschen [ID](#)¹⁹,
I. Fleck [ID](#)¹³⁸, P. Fleischmann [ID](#)¹⁰³, T. Flick [ID](#)¹⁶⁷, B.M. Flierl [ID](#)¹⁰⁶, L. Flores [ID](#)¹²⁵, M. Flores [ID](#)^{31d},
L.R. Flores Castillo [ID](#)^{62a}, F.M. Follega [ID](#)^{75a,75b}, N. Fomin [ID](#)¹⁵, J.H. Foo [ID](#)¹⁵², B.C. Forland [ID](#)⁶⁵,
A. Formica [ID](#)¹³², F.A. Förster [ID](#)¹², A.C. Forti [ID](#)⁹⁸, E. Fortin [ID](#)⁹⁹, M.G. Foti [ID](#)¹²³, L. Fountas [ID](#)^{8,i},
D. Fournier [ID](#)⁶⁴, H. Fox [ID](#)⁸⁸, P. Francavilla [ID](#)^{71a,71b}, S. Francescato [ID](#)⁵⁹, M. Franchini [ID](#)^{21b,21a},
S. Franchino [ID](#)^{61a}, D. Francis [ID](#)³⁴, L. Franco [ID](#)⁴, L. Franconi [ID](#)¹⁸, M. Franklin [ID](#)⁵⁹,
G. Frattari [ID](#)^{72a,72b}, A.C. Freegard [ID](#)⁹¹, P.M. Freeman [ID](#)¹⁹, W.S. Freund [ID](#)^{79b}, E.M. Freundlich [ID](#)⁴⁷,
D. Froidevaux [ID](#)³⁴, J.A. Frost [ID](#)¹²³, Y. Fu [ID](#)^{60a}, M. Fujimoto [ID](#)¹¹⁵, E. Fullana Torregrosa [ID](#)^{159,*},
J. Fuster [ID](#)¹⁵⁹, A. Gabrielli [ID](#)^{21b,21a}, A. Gabrielli [ID](#)³⁴, P. Gadow [ID](#)⁴⁶, G. Gagliardi [ID](#)^{55b,55a},
L.G. Gagnon [ID](#)^{16a}, G.E. Gallardo [ID](#)¹²³, E.J. Gallas [ID](#)¹²³, B.J. Gallop [ID](#)¹³¹, R. Gamboa Goni [ID](#)⁹¹,
K.K. Gan [ID](#)¹¹⁶, S. Ganguly [ID](#)¹⁵⁰, J. Gao [ID](#)^{60a}, Y. Gao [ID](#)⁵⁰, Y.S. Gao [ID](#)^{29,n},

F.M. Garay Walls [ID](#)^{134a}, C. García [ID](#)¹⁵⁹, J.E. García Navarro [ID](#)¹⁵⁹, J.A. García Pascual [ID](#)^{13a}, M. Garcia-Sciveres [ID](#)^{16a}, R.W. Gardner [ID](#)³⁷, D. Garg [ID](#)⁷⁷, R.B. Garg [ID](#)^{140,ak}, S. Gargiulo [ID](#)⁵², C.A. Garner [ID](#)¹⁵², V. Garonne [ID](#)¹²², S.J. Gasiorowski [ID](#)¹³⁵, P. Gaspar [ID](#)^{79b}, G. Gaudio [ID](#)^{70a}, P. Gauzzi [ID](#)^{72a,72b}, I.L. Gavrilenko [ID](#)³⁵, A. Gavrilyuk [ID](#)³⁵, C. Gay [ID](#)¹⁶⁰, G. Gaycken [ID](#)⁴⁶, E.N. Gazis [ID](#)⁹, A.A. Geanta [ID](#)^{25b}, C.M. Gee [ID](#)¹³³, C.N.P. Gee [ID](#)¹³¹, J. Geisen [ID](#)⁹⁵, M. Geisen [ID](#)⁹⁷, C. Gemme [ID](#)^{55b}, M.H. Genest [ID](#)⁵⁸, S. Gentile [ID](#)^{72a,72b}, S. George [ID](#)⁹², W.F. George [ID](#)¹⁹, T. Geralis [ID](#)⁴⁴, L.O. Gerlach [ID](#)⁵³, P. Gessinger-Befurt [ID](#)³⁴, M. Ghasemi Bostanabad [ID](#)¹⁶¹, M. Ghneimat [ID](#)¹³⁸, A. Ghosh [ID](#)¹⁵⁶, A. Ghosh [ID](#)⁷⁷, B. Giacobbe [ID](#)^{21b}, S. Giagu [ID](#)^{72a,72b}, N. Giangiacomi [ID](#)¹⁵², P. Giannetti [ID](#)^{71a}, A. Giannini [ID](#)^{69a,69b}, S.M. Gibson [ID](#)⁹², M. Gignac [ID](#)¹³³, D.T. Gil [ID](#)^{82b}, B.J. Gilbert [ID](#)³⁹, D. Gillberg [ID](#)³², G. Gilles [ID](#)¹¹¹, N.E.K. Gillwald [ID](#)⁴⁶, D.M. Gingrich [ID](#)^{2,af}, M.P. Giordani [ID](#)^{66a,66c}, P.F. Giraud [ID](#)¹³², G. Giugliarelli [ID](#)^{66a,66c}, D. Giugni [ID](#)^{68a}, F. Giuli [ID](#)^{73a,73b}, I. Gkialas [ID](#)^{8,i}, P. Gkoutoumis [ID](#)⁹, L.K. Gladilin [ID](#)³⁵, C. Glasman [ID](#)⁹⁶, G.R. Gledhill [ID](#)¹²⁰, M. Glisic [ID](#)¹²⁰, I. Gnesi [ID](#)^{41b,e}, M. Goblirsch-Kolb [ID](#)²⁴, D. Godin [ID](#)¹⁰⁵, S. Goldfarb [ID](#)¹⁰², T. Golling [ID](#)⁵⁴, D. Golubkov [ID](#)³⁵, J.P. Gombas [ID](#)¹⁰⁴, A. Gomes [ID](#)^{127a,127b}, R. Goncalves Gama [ID](#)⁵³, R. Gonçalo [ID](#)^{127a,127c}, G. Gonella [ID](#)¹²⁰, L. Gonella [ID](#)¹⁹, A. Gongadze [ID](#)³⁶, F. Gonnella [ID](#)¹⁹, J.L. Gonski [ID](#)³⁹, R.Y. González Andana [ID](#)^{134a}, S. González de la Hoz [ID](#)¹⁵⁹, S. Gonzalez Fernandez [ID](#)¹², R. Gonzalez Lopez [ID](#)⁸⁹, C. Gonzalez Renteria [ID](#)^{16a}, R. Gonzalez Suarez [ID](#)¹⁵⁷, S. Gonzalez-Sevilla [ID](#)⁵⁴, G.R. Gonzalvo Rodriguez [ID](#)¹⁵⁹, L. Goossens [ID](#)³⁴, N.A. Gorasia [ID](#)¹⁹, P.A. Gorbounov [ID](#)³⁵, B. Gorini [ID](#)³⁴, E. Gorini [ID](#)^{67a,67b}, A. Gorišek [ID](#)⁹⁰, A.T. Goshaw [ID](#)⁴⁹, M.I. Gostkin [ID](#)³⁶, C.A. Gottardo [ID](#)¹¹⁰, M. Gouighri [ID](#)^{33b}, V. Goumarre [ID](#)⁴⁶, A.G. Goussiou [ID](#)¹³⁵, N. Govender [ID](#)^{31c}, C. Goy [ID](#)⁴, I. Grabowska-Bold [ID](#)^{82a}, K. Graham [ID](#)³², E. Gramstad [ID](#)¹²², S. Grancagnolo [ID](#)¹⁷, M. Grandi [ID](#)¹⁴³, V. Gratchev [ID](#)^{35,*}, P.M. Gravila [ID](#)^{25f}, F.G. Gravili [ID](#)^{67a,67b}, H.M. Gray [ID](#)^{16a}, C. Greife [ID](#)²², I.M. Gregor [ID](#)⁴⁶, P. Grenier [ID](#)¹⁴⁰, K. Grevtsov [ID](#)⁴⁶, C. Grieco [ID](#)¹², N.A. Grieser [ID](#)¹¹⁷, A.A. Grillo [ID](#)¹³³, K. Grimm [ID](#)^{29,m}, S. Grinstein [ID](#)^{12,t}, J.-F. Grivaz [ID](#)⁶⁴, S. Groh [ID](#)⁹⁷, E. Gross [ID](#)¹⁶⁵, J. Grosse-Knetter [ID](#)⁵³, C. Grud [ID](#)¹⁰³, A. Grummer [ID](#)¹⁰⁹, J.C. Grundy [ID](#)¹²³, L. Guan [ID](#)¹⁰³, W. Guan [ID](#)¹⁶⁶, C. Gubbels [ID](#)¹⁶⁰, J. Guenther [ID](#)³⁴, J.G.R. Guerrero Rojas [ID](#)¹⁵⁹, F. Guescini [ID](#)¹⁰⁷, R. Gugel [ID](#)⁹⁷, A. Guida [ID](#)⁴⁶, T. Guillemain [ID](#)⁴, S. Guindon [ID](#)³⁴, J. Guo [ID](#)^{60c}, L. Guo [ID](#)⁶⁴, Y. Guo [ID](#)¹⁰³, R. Gupta [ID](#)⁴⁶, S. Gurbuz [ID](#)²², G. Gustavino [ID](#)¹¹⁷, M. Guth [ID](#)⁵⁴, P. Gutierrez [ID](#)¹¹⁷, L.F. Gutierrez Zagazeta [ID](#)¹²⁵, C. Gutsche [ID](#)⁹³, C. Guyot [ID](#)¹³², C. Gwenlan [ID](#)¹²³, C.B. Gwilliam [ID](#)⁸⁹, E.S. Haaland [ID](#)¹²², A. Haas [ID](#)¹¹⁴, M. Habedank [ID](#)⁴⁶, C. Haber [ID](#)^{16a}, H.K. Hadavand [ID](#)⁷, A. Hadeef [ID](#)⁹⁷, S. Hadzic [ID](#)¹⁰⁷, M. Haleem [ID](#)¹⁶², J. Haley [ID](#)¹¹⁸, J.J. Hall [ID](#)¹³⁶, G. Halladjian [ID](#)¹⁰⁴, G.D. Hallowell [ID](#)⁹⁹, L. Halser [ID](#)¹⁸, K. Hamano [ID](#)¹⁶¹, H. Hamdaoui [ID](#)^{33e}, M. Hamer [ID](#)²², G.N. Hamity [ID](#)⁵⁰, K. Han [ID](#)^{60a}, L. Han [ID](#)^{13c}, L. Han [ID](#)^{60a}, S. Han [ID](#)^{16a}, Y.F. Han [ID](#)¹⁵², K. Hanagaki [ID](#)⁸⁰, M. Hance [ID](#)¹³³, M.D. Hank [ID](#)³⁷, R. Hankache [ID](#)⁹⁸, E. Hansen [ID](#)⁹⁵, J.B. Hansen [ID](#)⁴⁰, J.D. Hansen [ID](#)⁴⁰, M.C. Hansen [ID](#)²², P.H. Hansen [ID](#)⁴⁰, K. Hara [ID](#)¹⁵⁴, T. Harenberg [ID](#)¹⁶⁷, S. Harkusha [ID](#)³⁵, Y.T. Harris [ID](#)¹²³, P.F. Harrison [ID](#)¹⁶³, N.M. Hartman [ID](#)¹⁴⁰, N.M. Hartmann [ID](#)¹⁰⁶, Y. Hasegawa [ID](#)¹³⁷, A. Hasib [ID](#)⁵⁰, S. Hassani [ID](#)¹³², S. Haug [ID](#)¹⁸, R. Hauser [ID](#)¹⁰⁴, M. Havranek [ID](#)¹²⁹, C.M. Hawkes [ID](#)¹⁹, R.J. Hawkings [ID](#)³⁴, S. Hayashida [ID](#)¹⁰⁸, D. Hayden [ID](#)¹⁰⁴, C. Hayes [ID](#)¹⁰³, R.L. Hayes [ID](#)¹⁶⁰, C.P. Hays [ID](#)¹²³, J.M. Hays [ID](#)⁹¹, H.S. Hayward [ID](#)⁸⁹, S.J. Haywood [ID](#)¹³¹, F. He [ID](#)^{60a}, Y. He [ID](#)¹⁵¹, Y. He [ID](#)¹²⁴, M.P. Heath [ID](#)⁵⁰, V. Hedberg [ID](#)⁹⁵, A.L. Heggelund [ID](#)¹²², N.D. Hehir [ID](#)⁹¹, C. Heidegger [ID](#)⁵², K.K. Heidegger [ID](#)⁵², W.D. Heidorn [ID](#)⁷⁸, J. Heilman [ID](#)³², S. Heim [ID](#)⁴⁶, T. Heim [ID](#)^{16a},

B. Heinemann [ID](#)^{46,ac}, J.G. Heinlein [ID](#)¹²⁵, J.J. Heinrich [ID](#)¹²⁰, L. Heinrich [ID](#)³⁴, J. Hejbal [ID](#)¹²⁸,
 L. Helary [ID](#)⁴⁶, A. Held [ID](#)¹¹⁴, S. Hellesund [ID](#)¹²², C.M. Helling [ID](#)¹³³, S. Hellman [ID](#)^{45a,45b},
 C. Helsens [ID](#)³⁴, R.C.W. Henderson⁸⁸, L. Henkelmann [ID](#)³⁰, A.M. Henriques Correia³⁴,
 H. Herde [ID](#)¹⁴⁰, Y. Hernández Jiménez [ID](#)¹⁴², H. Herr⁹⁷, M.G. Herrmann [ID](#)¹⁰⁶, T. Herrmann [ID](#)⁴⁸,
 G. Herten [ID](#)⁵², R. Hertenberger [ID](#)¹⁰⁶, L. Hervas [ID](#)³⁴, N.P. Hessey [ID](#)^{153a}, H. Hibi [ID](#)⁸¹,
 S. Higashino [ID](#)⁸⁰, E. Higón-Rodríguez [ID](#)¹⁵⁹, K.H. Hiller⁴⁶, S.J. Hillier [ID](#)¹⁹, M. Hils [ID](#)⁴⁸,
 I. Hinchliffe [ID](#)^{16a}, F. Hinterkeuser [ID](#)²², M. Hirose [ID](#)¹²¹, S. Hirose [ID](#)¹⁵⁴, D. Hirschbuehl [ID](#)¹⁶⁷,
 B. Hiti [ID](#)⁹⁰, O. Hladik¹²⁸, J. Hobbs [ID](#)¹⁴², R. Hobincu [ID](#)^{25e}, N. Hod [ID](#)¹⁶⁵, M.C. Hodgkinson [ID](#)¹³⁶,
 B.H. Hodgkinson [ID](#)³⁰, A. Hoecker [ID](#)³⁴, J. Hofer [ID](#)⁴⁶, D. Hohn [ID](#)⁵², T. Holm [ID](#)²², T.R. Holmes [ID](#)³⁷,
 M. Holzbock [ID](#)¹⁰⁷, L.B.A.H. Hommels [ID](#)³⁰, B.P. Honan [ID](#)⁹⁸, J. Hong [ID](#)^{60c}, T.M. Hong [ID](#)¹²⁶,
 Y. Hong [ID](#)⁵³, J.C. Honig [ID](#)⁵², A. Hönle [ID](#)¹⁰⁷, B.H. Hooberman [ID](#)¹⁵⁸, W.H. Hopkins [ID](#)⁵,
 Y. Horii [ID](#)¹⁰⁸, L.A. Horyn [ID](#)³⁷, S. Hou [ID](#)¹⁴⁵, J. Howarth [ID](#)⁵⁷, J. Hoya [ID](#)⁸⁷, M. Hrabovsky [ID](#)¹¹⁹,
 A. Hrynevich [ID](#)³⁵, T. Hryn'ova [ID](#)⁴, P.J. Hsu [ID](#)⁶³, S.-C. Hsu [ID](#)¹³⁵, Q. Hu [ID](#)³⁹, S. Hu [ID](#)^{60c},
 Y.F. Hu [ID](#)^{13a,13d,ah}, D.P. Huang [ID](#)⁹³, X. Huang [ID](#)^{13c}, Y. Huang [ID](#)^{60a}, Y. Huang [ID](#)^{13a},
 Z. Hubacek [ID](#)¹²⁹, F. Hubaut [ID](#)⁹⁹, M. Huebner [ID](#)²², F. Huegging [ID](#)²², T.B. Huffman [ID](#)¹²³,
 M. Huhtinen [ID](#)³⁴, S.K. Huiberts [ID](#)¹⁵, R. Hulsken [ID](#)⁵⁸, N. Huseynov [ID](#)^{36,x}, J. Huston [ID](#)¹⁰⁴,
 J. Huth [ID](#)⁵⁹, R. Hyneman [ID](#)¹⁴⁰, S. Hyrych [ID](#)^{26a}, G. Iacobucci [ID](#)⁵⁴, G. Iakovidis [ID](#)²⁷,
 I. Ibragimov [ID](#)¹³⁸, L. Iconomidou-Fayard [ID](#)⁶⁴, P. Iengo [ID](#)³⁴, R. Iguchi [ID](#)¹⁵⁰, T. Iizawa [ID](#)⁵⁴,
 Y. Ikegami [ID](#)⁸⁰, A. Ilg [ID](#)¹⁸, N. Ilic [ID](#)¹⁵², H. Imam [ID](#)^{33a}, T. Ingebretsen Carlson [ID](#)^{45a,45b},
 G. Introzzi [ID](#)^{70a,70b}, M. Iodice [ID](#)^{74a}, V. Ippolito [ID](#)^{72a,72b}, M. Ishino [ID](#)¹⁵⁰, W. Islam [ID](#)¹⁶⁶,
 C. Issever [ID](#)^{17,46}, S. Istin [ID](#)^{11c,ai}, J.M. Iturbe Ponce [ID](#)^{62a}, R. Iuppa [ID](#)^{75a,75b}, A. Ivina [ID](#)¹⁶⁵,
 J.M. Izen [ID](#)⁴³, V. Izzo [ID](#)^{69a}, P. Jacka [ID](#)¹²⁸, P. Jackson [ID](#)¹, R.M. Jacobs [ID](#)⁴⁶, B.P. Jaeger [ID](#)¹³⁹,
 C.S. Jagfeld [ID](#)¹⁰⁶, G. Jäkel [ID](#)¹⁶⁷, K. Jakobs [ID](#)⁵², T. Jakoubek [ID](#)¹⁶⁵, J. Jamieson [ID](#)⁵⁷,
 K.W. Janas [ID](#)^{82a}, G. Jarlskog [ID](#)⁹⁵, A.E. Jaspan [ID](#)⁸⁹, N. Javadov^{36,x}, T. Javůrek [ID](#)³⁴,
 M. Javurkova [ID](#)¹⁰⁰, F. Jeanneau [ID](#)¹³², L. Jeanty [ID](#)¹²⁰, J. Jejelava [ID](#)^{146a,y}, P. Jenni [ID](#)^{52,f},
 S. Jézéquel [ID](#)⁴, J. Jia [ID](#)¹⁴², Z. Jia [ID](#)^{13c}, Y. Jiang^{60a}, S. Jiggins [ID](#)⁵⁰, J. Jimenez Pena [ID](#)¹⁰⁷,
 S. Jin [ID](#)^{13c}, A. Jinaru [ID](#)^{25b}, O. Jinnouchi [ID](#)¹⁵¹, H. Jivan [ID](#)^{31f}, P. Johansson [ID](#)¹³⁶, K.A. Johns [ID](#)⁶,
 C.A. Johnson [ID](#)⁶⁵, D.M. Jones [ID](#)³⁰, E. Jones [ID](#)¹⁶³, R.W.L. Jones [ID](#)⁸⁸, T.J. Jones [ID](#)⁸⁹,
 J. Jovicevic [ID](#)¹⁴, X. Ju [ID](#)^{16a}, J.J. Junggeburth [ID](#)³⁴, A. Juste Rozas [ID](#)^{12,t}, S. Kabana [ID](#)^{134e},
 A. Kaczmarek [ID](#)⁸³, M. Kado [ID](#)^{72a,72b}, H. Kagan [ID](#)¹¹⁶, M. Kagan [ID](#)¹⁴⁰, A. Kahn³⁹, A. Kahn [ID](#)¹²⁵,
 C. Kahra [ID](#)⁹⁷, T. Kaji [ID](#)¹⁶⁴, E. Kajomovitz [ID](#)¹⁴⁷, C.W. Kalderon [ID](#)²⁷, A. Kamenshchikov [ID](#)³⁵,
 M. Kaneda [ID](#)¹⁵⁰, N.J. Kang [ID](#)¹³³, S. Kang [ID](#)⁷⁸, Y. Kano [ID](#)¹⁰⁸, D. Kar [ID](#)^{31f}, K. Karava [ID](#)¹²³,
 M.J. Kareem [ID](#)^{153b}, I. Karkanas [ID](#)¹⁴⁹, S.N. Karpov [ID](#)³⁶, Z.M. Karpova [ID](#)³⁶, V. Kartvelishvili [ID](#)⁸⁸,
 A.N. Karyukhin [ID](#)³⁵, E. Kasimi [ID](#)¹⁴⁹, C. Kato [ID](#)^{60d}, J. Katzy [ID](#)⁴⁶, K. Kawade [ID](#)¹³⁷,
 K. Kawagoe [ID](#)⁸⁶, T. Kawaguchi [ID](#)¹⁰⁸, T. Kawamoto [ID](#)¹³², G. Kawamura⁵³, E.F. Kay [ID](#)¹⁶¹,
 F.I. Kaya [ID](#)¹⁵⁵, S. Kazakos [ID](#)¹², V.F. Kazanin [ID](#)³⁵, Y. Ke [ID](#)¹⁴², J.M. Keaveney [ID](#)^{31a},
 R. Keeler [ID](#)¹⁶¹, J.S. Keller [ID](#)³², A.S. Kelly⁹³, D. Kelsey [ID](#)¹⁴³, J.J. Kempster [ID](#)¹⁹, J. Kendrick [ID](#)¹⁹,
 K.E. Kennedy [ID](#)³⁹, O. Kepka [ID](#)¹²⁸, S. Kersten [ID](#)¹⁶⁷, B.P. Kerševan [ID](#)⁹⁰,
 S. Ketabchi Haghighat [ID](#)¹⁵², M. Khandoga [ID](#)¹²⁴, A. Khanov [ID](#)¹¹⁸, A.G. Kharlamov [ID](#)³⁵,
 T. Kharlamova [ID](#)³⁵, E.E. Khoda [ID](#)¹³⁵, T.J. Khoo [ID](#)¹⁷, G. Khoraiuli [ID](#)¹⁶², J. Khubua [ID](#)^{146b},
 S. Kido [ID](#)⁸¹, M. Kiehn [ID](#)³⁴, A. Kilgallon [ID](#)¹²⁰, E. Kim [ID](#)¹⁵¹, Y.K. Kim [ID](#)³⁷, N. Kimura [ID](#)⁹³,
 A. Kirchhoff [ID](#)⁵³, D. Kirchmeier [ID](#)⁴⁸, C. Kirfel [ID](#)²², J. Kirk [ID](#)¹³¹, A.E. Kiryunin [ID](#)¹⁰⁷,
 T. Kishimoto [ID](#)¹⁵⁰, D.P. Kisiuk¹⁵², C. Kitsaki [ID](#)⁹, O. Kivernyk [ID](#)²²,

T. Klapdor-Kleingrothaus [ID](#)⁵², M. Klassen [ID](#)^{61a}, C. Klein [ID](#)³², L. Klein [ID](#)¹⁶², M.H. Klein [ID](#)¹⁰³, M. Klein [ID](#)⁸⁹, U. Klein [ID](#)⁸⁹, P. Klimek [ID](#)³⁴, A. Klimentov [ID](#)²⁷, F. Klimpel [ID](#)¹⁰⁷, T. Klingl [ID](#)²², T. Klioutchnikova [ID](#)³⁴, F.F. Klitzner [ID](#)¹⁰⁶, P. Kluit [ID](#)¹¹¹, S. Kluth [ID](#)¹⁰⁷, E. Kneringer [ID](#)⁷⁶, T.M. Knight [ID](#)¹⁵², A. Knue [ID](#)⁵², D. Kobayashi⁸⁶, R. Kobayashi [ID](#)⁸⁴, M. Kobel [ID](#)⁴⁸, M. Kocian [ID](#)¹⁴⁰, T. Kodama¹⁵⁰, P. Kodyš [ID](#)¹³⁰, D.M. Koeck [ID](#)¹⁴³, P.T. Koenig [ID](#)²², T. Koffas [ID](#)³², N.M. Köhler [ID](#)³⁴, M. Kolb [ID](#)¹³², I. Koletsou [ID](#)⁴, T. Komarek [ID](#)¹¹⁹, K. Köneke [ID](#)⁵², A.X.Y. Kong [ID](#)¹, T. Kono [ID](#)¹¹⁵, V. Konstantinides⁹³, N. Konstantinidis [ID](#)⁹³, B. Konya [ID](#)⁹⁵, R. Kopeliansky [ID](#)⁶⁵, S. Koperny [ID](#)^{82a}, K. Korcyl [ID](#)⁸³, K. Kordas [ID](#)¹⁴⁹, G. Koren [ID](#)¹⁴⁸, A. Korn [ID](#)⁹³, S. Korn [ID](#)⁵³, I. Korolkov [ID](#)¹², E.V. Korolkova¹³⁶, N. Korotkova [ID](#)³⁵, B. Kortman [ID](#)¹¹¹, O. Kortner [ID](#)¹⁰⁷, S. Kortner [ID](#)¹⁰⁷, W.H. Kostecka [ID](#)¹¹², V.V. Kostyukhin [ID](#)^{138,35}, A. Kotsokechagia [ID](#)⁶⁴, A. Kotwal [ID](#)⁴⁹, A. Koulouris [ID](#)³⁴, A. Kourkouveli-Charalampidi [ID](#)^{70a,70b}, C. Kourkouvelis [ID](#)⁸, E. Kourlitis [ID](#)⁵, O. Kovanda [ID](#)¹⁴³, R. Kowalewski [ID](#)¹⁶¹, W. Kozanecki [ID](#)¹³², A.S. Kozhin [ID](#)³⁵, V.A. Kramarenko [ID](#)³⁵, G. Kramberger [ID](#)⁹⁰, P. Kramer [ID](#)⁹⁷, D. Krasnopevtsev [ID](#)^{60a}, M.W. Krasny [ID](#)¹²⁴, A. Krasznahorkay [ID](#)³⁴, J.A. Kremer [ID](#)⁹⁷, J. Kretzschmar [ID](#)⁸⁹, K. Kreul [ID](#)¹⁷, P. Krieger [ID](#)¹⁵², F. Krieter [ID](#)¹⁰⁶, S. Krishnamurthy [ID](#)¹⁰⁰, A. Krishnan [ID](#)^{61b}, M. Krivos [ID](#)¹³⁰, K. Krizka [ID](#)^{16a}, K. Kroeninger [ID](#)⁴⁷, H. Kroha [ID](#)¹⁰⁷, J. Kroll [ID](#)¹²⁸, J. Kroll [ID](#)¹²⁵, K.S. Krowpman [ID](#)¹⁰⁴, U. Kruchonak [ID](#)³⁶, H. Krüger [ID](#)²², N. Krumnack⁷⁸, M.C. Kruse [ID](#)⁴⁹, J.A. Krzysiak [ID](#)⁸³, A. Kubota [ID](#)¹⁵¹, O. Kuchinskaia [ID](#)³⁵, S. Kuday [ID](#)^{3a}, D. Kuechler [ID](#)⁴⁶, J.T. Kuechler [ID](#)⁴⁶, S. Kuehn [ID](#)³⁴, T. Kuhl [ID](#)⁴⁶, V. Kukhtin [ID](#)³⁶, Y. Kulchitsky [ID](#)^{35,a}, S. Kuleshov [ID](#)^{134d}, M. Kumar [ID](#)^{31f}, N. Kumari [ID](#)⁹⁹, M. Kuna [ID](#)⁵⁸, A. Kupco [ID](#)¹²⁸, T. Kupfer⁴⁷, O. Kuprash [ID](#)⁵², H. Kurashige [ID](#)⁸¹, L.L. Kurchaninov [ID](#)^{153a}, Y.A. Kurochkin [ID](#)³⁵, A. Kurova [ID](#)³⁵, M.G. Kurth^{13a,13d}, E.S. Kuwertz [ID](#)³⁴, M. Kuze [ID](#)¹⁵¹, A.K. Kvam [ID](#)¹³⁵, J. Kvita [ID](#)¹¹⁹, T. Kwan [ID](#)¹⁰¹, K.W. Kwok [ID](#)^{62a}, C. Lacasta [ID](#)¹⁵⁹, F. Lacava [ID](#)^{72a,72b}, H. Lacker [ID](#)¹⁷, D. Lacour [ID](#)¹²⁴, N.N. Lad [ID](#)⁹³, E. Ladygin [ID](#)³⁶, R. Lafaye [ID](#)⁴, B. Laforge [ID](#)¹²⁴, T. Lagouri [ID](#)^{134e}, S. Lai [ID](#)⁵³, I.K. Lakomic [ID](#)^{82a}, N. Lalloue [ID](#)⁵⁸, J.E. Lambert [ID](#)¹¹⁷, S. Lammers [ID](#)⁶⁵, W. Lampl [ID](#)⁶, C. Lampoudis [ID](#)¹⁴⁹, E. Lançon [ID](#)²⁷, U. Landgraf [ID](#)⁵², M.P.J. Landon [ID](#)⁹¹, V.S. Lang [ID](#)⁵², J.C. Lange [ID](#)⁵³, R.J. Langenberg [ID](#)¹⁰⁰, A.J. Lankford [ID](#)¹⁵⁶, F. Lanni [ID](#)²⁷, K. Lantzsch [ID](#)²², A. Lanza [ID](#)^{70a}, A. Lapertosa [ID](#)^{55b,55a}, J.F. Laporte [ID](#)¹³², T. Lari [ID](#)^{68a}, F. Lasagni Manghi [ID](#)^{21b}, M. Lassnig [ID](#)³⁴, V. Latonova [ID](#)¹²⁸, T.S. Lau [ID](#)^{62a}, A. Laudrain [ID](#)⁹⁷, A. Laurier [ID](#)³², M. Lavorgna [ID](#)^{69a,69b}, S.D. Lawlor [ID](#)⁹², Z. Lawrence [ID](#)⁹⁸, M. Lazzaroni [ID](#)^{68a,68b}, B. Le⁹⁸, B. Leban [ID](#)⁹⁰, A. Lebedev [ID](#)⁷⁸, M. LeBlanc [ID](#)³⁴, T. LeCompte [ID](#)⁵, F. Ledroit-Guillon [ID](#)⁵⁸, A.C.A. Lee⁹³, G.R. Lee [ID](#)¹⁵, L. Lee [ID](#)⁵⁹, S.C. Lee [ID](#)¹⁴⁵, S. Lee [ID](#)⁷⁸, L.L. Leeuw [ID](#)^{31c}, B. Lefebvre [ID](#)^{153a}, H.P. Lefebvre [ID](#)⁹², M. Lefebvre [ID](#)¹⁶¹, C. Leggett [ID](#)^{16a}, K. Lehmann [ID](#)¹³⁹, N. Lehmann [ID](#)¹⁸, G. Lehmann Miotto [ID](#)³⁴, W.A. Leight [ID](#)⁴⁶, A. Leisos [ID](#)^{149,s}, M.A.L. Leite [ID](#)^{79d}, C.E. Leitgeb [ID](#)⁴⁶, R. Leitner [ID](#)¹³⁰, K.J.C. Leney [ID](#)⁴², T. Lenz [ID](#)²², S. Leone [ID](#)^{71a}, C. Leonidopoulos [ID](#)⁵⁰, A. Leopold [ID](#)¹⁴¹, C. Leroy [ID](#)¹⁰⁵, R. Les [ID](#)¹⁰⁴, C.G. Lester [ID](#)³⁰, M. Levchenko [ID](#)³⁵, J. Levêque [ID](#)⁴, D. Levin [ID](#)¹⁰³, L.J. Levinson [ID](#)¹⁶⁵, D.J. Lewis [ID](#)¹⁹, B. Li [ID](#)^{13b}, B. Li [ID](#)^{60b}, C. Li^{60a}, C-Q. Li [ID](#)^{60c,60d}, H. Li [ID](#)^{60a}, H. Li [ID](#)^{60b}, H. Li [ID](#)^{60b}, J. Li [ID](#)^{60c}, K. Li [ID](#)¹³⁵, L. Li [ID](#)^{60c}, M. Li [ID](#)^{13a,13d}, Q.Y. Li [ID](#)^{60a}, S. Li [ID](#)^{60d,60c,d}, T. Li [ID](#)^{60b}, X. Li [ID](#)⁴⁶, Y. Li [ID](#)⁴⁶, Z. Li [ID](#)^{60b}, Z. Li [ID](#)¹²³, Z. Li [ID](#)¹⁰¹, Z. Li [ID](#)⁸⁹, Z. Liang [ID](#)^{13a}, M. Liberatore [ID](#)⁴⁶, B. Liberti [ID](#)^{73a}, K. Lie [ID](#)^{62c}, J. Lieber Marin [ID](#)^{79b}, K. Lin [ID](#)¹⁰⁴, R.A. Linck [ID](#)⁶⁵, R.E. Lindley [ID](#)⁶, J.H. Lindon [ID](#)², A. Linss [ID](#)⁴⁶, E. Lipeles [ID](#)¹²⁵, A. Lipniacka [ID](#)¹⁵, T.M. Liss [ID](#)^{158,ad}, A. Lister [ID](#)¹⁶⁰, J.D. Little [ID](#)⁷, B. Liu [ID](#)^{13a}, B.X. Liu [ID](#)¹³⁹, J.B. Liu [ID](#)^{60a}, J.K.K. Liu [ID](#)³⁷, K. Liu [ID](#)^{60d,60c},

M. Liu [ID](#)^{60a}, M.Y. Liu [ID](#)^{60a}, P. Liu [ID](#)^{13a}, X. Liu [ID](#)^{60a}, Y. Liu [ID](#)⁴⁶, Y. Liu [ID](#)^{13c,13d}, Y.L. Liu [ID](#)¹⁰³,
 Y.W. Liu [ID](#)^{60a}, M. Livan [ID](#)^{70a,70b}, J. Llorente Merino [ID](#)¹³⁹, S.L. Lloyd [ID](#)⁹¹, E.M. Lobodzinska [ID](#)⁴⁶,
 P. Loch [ID](#)⁶, S. Loffredo [ID](#)^{73a,73b}, T. Lohse [ID](#)¹⁷, K. Lohwasser [ID](#)¹³⁶, M. Lokajicek [ID](#)¹²⁸,
 J.D. Long [ID](#)¹⁵⁸, I. Longarini [ID](#)^{72a,72b}, L. Longo [ID](#)³⁴, R. Longo [ID](#)¹⁵⁸, I. Lopez Paz [ID](#)¹²,
 A. Lopez Solis [ID](#)⁴⁶, J. Lorenz [ID](#)¹⁰⁶, N. Lorenzo Martinez [ID](#)⁴, A.M. Lory [ID](#)¹⁰⁶, A. Lösle [ID](#)⁵²,
 X. Lou [ID](#)^{45a,45b}, X. Lou [ID](#)^{13a,13d}, A. Lounis [ID](#)⁶⁴, J. Love [ID](#)⁵, P.A. Love [ID](#)⁸⁸,
 J.J. Lozano Bahilo [ID](#)¹⁵⁹, G. Lu [ID](#)^{13a,13d}, M. Lu [ID](#)^{60a}, S. Lu [ID](#)¹²⁵, Y.J. Lu [ID](#)⁶³, H.J. Lubatti [ID](#)¹³⁵,
 C. Luci [ID](#)^{72a,72b}, F.L. Lucio Alves [ID](#)^{13c}, A. Lucotte [ID](#)⁵⁸, F. Luehring [ID](#)⁶⁵, I. Luise [ID](#)¹⁴²,
 L. Luminari [ID](#)^{72a}, O. Lundberg [ID](#)¹⁴¹, B. Lund-Jensen [ID](#)¹⁴¹, N.A. Luongo [ID](#)¹²⁰, M.S. Lutz [ID](#)¹⁴⁸,
 D. Lynn [ID](#)²⁷, H. Lyons⁸⁹, R. Lysak [ID](#)¹²⁸, E. Lytken [ID](#)⁹⁵, F. Lyu [ID](#)^{13a}, V. Lyubushkin [ID](#)³⁶,
 T. Lyubushkina [ID](#)³⁶, H. Ma [ID](#)²⁷, L.L. Ma [ID](#)^{60b}, Y. Ma [ID](#)⁹³, D.M. Mac Donell [ID](#)¹⁶¹,
 G. Maccarrone [ID](#)⁵¹, C.M. Macdonald [ID](#)¹³⁶, J.C. MacDonald [ID](#)¹³⁶, R. Madar [ID](#)³⁸, W.F. Mader [ID](#)⁴⁸,
 M. Madugoda Ralalage Don [ID](#)¹¹⁸, N. Madysa [ID](#)⁴⁸, J. Maeda [ID](#)⁸¹, T. Maeno [ID](#)²⁷, M. Maerker [ID](#)⁴⁸,
 V. Magerl [ID](#)⁵², J. Magro [ID](#)^{66a,66c}, D.J. Mahon [ID](#)³⁹, C. Maidantchik [ID](#)^{79b}, A. Maio [ID](#)^{127a,127b,127d},
 K. Maj [ID](#)^{82a}, O. Majersky [ID](#)^{26a}, S. Majewski [ID](#)¹²⁰, N. Makovec [ID](#)⁶⁴, V. Maksimovic [ID](#)¹⁴,
 B. Malaescu [ID](#)¹²⁴, Pa. Malecki [ID](#)⁸³, V.P. Maleev [ID](#)³⁵, F. Malek [ID](#)⁵⁸, D. Malito [ID](#)^{41b,41a},
 U. Mallik [ID](#)⁷⁷, C. Malone [ID](#)³⁰, S. Maltezos⁹, S. Malyukov³⁶, J. Mamuzic [ID](#)¹⁵⁹, G. Mancini [ID](#)⁵¹,
 J.P. Mandalia [ID](#)⁹¹, I. Mandić [ID](#)⁹⁰, L. Manhaes de Andrade Filho [ID](#)^{79a}, I.M. Maniatis [ID](#)¹⁴⁹,
 M. Manisha [ID](#)¹³², J. Manjarres Ramos [ID](#)⁴⁸, K.H. Mankinen [ID](#)⁹⁵, A. Mann [ID](#)¹⁰⁶, A. Manousos [ID](#)⁷⁶,
 B. Mansoulie [ID](#)¹³², I. Mantos [ID](#)¹⁴⁹, S. Manzoni [ID](#)¹¹¹, A. Marantis [ID](#)^{149,s}, G. Marchiori [ID](#)¹²⁴,
 M. Marcisovsky [ID](#)¹²⁸, L. Marcoccia [ID](#)^{73a,73b}, C. Marcon [ID](#)⁹⁵, M. Marjanovic [ID](#)¹¹⁷,
 Z. Marshall [ID](#)^{16a}, S. Marti-Garcia [ID](#)¹⁵⁹, T.A. Martin [ID](#)¹⁶³, V.J. Martin [ID](#)⁵⁰,
 B. Martin dit Latour [ID](#)¹⁵, L. Martinelli [ID](#)^{72a,72b}, M. Martinez [ID](#)^{12,t}, P. Martinez Agullo [ID](#)¹⁵⁹,
 V.I. Martinez Outschoorn [ID](#)¹⁰⁰, S. Martin-Haugh [ID](#)¹³¹, V.S. Martoiu [ID](#)^{25b}, A.C. Martyniuk [ID](#)⁹³,
 A. Marzin [ID](#)³⁴, S.R. Maschek [ID](#)¹⁰⁷, L. Masetti [ID](#)⁹⁷, T. Mashimo [ID](#)¹⁵⁰, J. Masik [ID](#)⁹⁸,
 A.L. Maslennikov [ID](#)³⁵, L. Massa [ID](#)^{21b}, P. Massarotti [ID](#)^{69a,69b}, P. Mastrandrea [ID](#)^{71a,71b},
 A. Mastroberardino [ID](#)^{41b,41a}, T. Masubuchi [ID](#)¹⁵⁰, D. Matakias²⁷, T. Mathisen [ID](#)¹⁵⁷, A. Matic [ID](#)¹⁰⁶,
 N. Matsuzawa¹⁵⁰, J. Maurer [ID](#)^{25b}, B. Maček [ID](#)⁹⁰, D.A. Maximov [ID](#)³⁵, R. Mazini [ID](#)¹⁴⁵,
 I. Maznas [ID](#)¹⁴⁹, S.M. Mazza [ID](#)¹³³, C. Mc Ginn [ID](#)²⁷, J.P. Mc Gowan [ID](#)¹⁰¹, S.P. Mc Kee [ID](#)¹⁰³,
 T.G. McCarthy [ID](#)¹⁰⁷, W.P. McCormack [ID](#)^{16a}, E.F. McDonald [ID](#)¹⁰², A.E. McDougall [ID](#)¹¹¹,
 J.A. Mcfayden [ID](#)¹⁴³, G. Mchedlidze [ID](#)^{146b}, M.A. McKay⁴², K.D. McLean [ID](#)¹⁶¹,
 S.J. McMahan [ID](#)¹³¹, P.C. McNamara [ID](#)¹⁰², R.A. McPherson [ID](#)^{161,w}, J.E. Mdhuli [ID](#)^{31f},
 Z.A. Meadows [ID](#)¹⁰⁰, S. Meehan [ID](#)³⁴, T. Megy [ID](#)³⁸, S. Mehlhase [ID](#)¹⁰⁶, A. Mehta [ID](#)⁸⁹,
 B. Meirose [ID](#)⁴³, D. Melini [ID](#)¹⁴⁷, B.R. Mellado Garcia [ID](#)^{31f}, A.H. Melo [ID](#)⁵³, F. Meloni [ID](#)⁴⁶,
 A. Melzer [ID](#)²², E.D. Mendes Gouveia [ID](#)^{127a}, A.M. Mendes Jacques Da Costa [ID](#)¹⁹, H.Y. Meng [ID](#)¹⁵²,
 L. Meng [ID](#)³⁴, S. Menke [ID](#)¹⁰⁷, M. Mentink [ID](#)³⁴, E. Meoni [ID](#)^{41b,41a}, C. Merlassino [ID](#)¹²³,
 P. Mermod [ID](#)^{54,*}, L. Merola [ID](#)^{69a,69b}, C. Meroni [ID](#)^{68a}, G. Merz¹⁰³, O. Meshkov [ID](#)³⁵,
 J.K.R. Meshreki [ID](#)¹³⁸, J. Metcalfe [ID](#)⁵, A.S. Mete [ID](#)⁵, C. Meyer [ID](#)⁶⁵, J-P. Meyer [ID](#)¹³²,
 M. Michetti [ID](#)¹⁷, R.P. Middleton [ID](#)¹³¹, L. Mijović [ID](#)⁵⁰, G. Mikenberg [ID](#)¹⁶⁵, M. Mikestikova [ID](#)¹²⁸,
 M. Mikuž [ID](#)⁹⁰, H. Mildner [ID](#)¹³⁶, A. Milic [ID](#)¹⁵², C.D. Milke [ID](#)⁴², D.W. Miller [ID](#)³⁷, L.S. Miller [ID](#)³²,
 A. Milov [ID](#)¹⁶⁵, D.A. Milstead^{45a,45b}, T. Min^{13c}, A.A. Minaenko [ID](#)³⁵, I.A. Minashvili [ID](#)^{146b},
 L. Mince [ID](#)⁵⁷, A.I. Mincer [ID](#)¹¹⁴, B. Mindur [ID](#)^{82a}, M. Mineev [ID](#)³⁶, Y. Minegishi¹⁵⁰, Y. Mino [ID](#)⁸⁴,
 L.M. Mir [ID](#)¹², M. Miralles Lopez [ID](#)¹⁵⁹, M. Mironova [ID](#)¹²³, T. Mitani [ID](#)¹⁶⁴, V.A. Mitsou [ID](#)¹⁵⁹,

M. Mittal^{60c}, O. Miu¹⁵², P.S. Miyagawa⁹¹, Y. Miyazaki⁸⁶, A. Mizukami⁸⁰,
 J.U. Mjörnmark⁹⁵, T. Mkrtychyan^{61a}, M. Mlynarikova¹¹², T. Moa^{45a,45b}, S. Mobius⁵³,
 K. Mochizuki¹⁰⁵, P. Moder⁴⁶, P. Mogg¹⁰⁶, A.F. Mohammed^{13a,13d}, S. Mohapatra³⁹,
 G. Mokgatitswane^{31f}, B. Mondal¹³⁸, S. Mondal¹²⁹, K. Mönig⁴⁶, E. Monnier⁹⁹,
 L. Monsonis Romero¹⁵⁹, A. Montalbano¹³⁹, J. Montejo Berlingen³⁴, M. Montella¹¹⁶,
 F. Monticelli⁸⁷, N. Morange⁶⁴, A.L. Moreira De Carvalho^{127a}, M. Moreno Llácer¹⁵⁹,
 C. Moreno Martinez¹², P. Morettini^{55b}, S. Morgenstern¹⁶³, D. Mori¹³⁹, M. Morii⁵⁹,
 M. Morinaga¹⁵⁰, V. Morisbak¹²², A.K. Morley³⁴, A.P. Morris⁹³, L. Morvaj³⁴,
 P. Moschovakos³⁴, B. Moser¹¹¹, M. Mosidze^{146b}, T. Moskalets⁵², P. Moskvitina¹¹⁰,
 J. Moss^{29,o}, E.J.W. Moyses¹⁰⁰, S. Muanza⁹⁹, J. Mueller¹²⁶, D. Muenstermann⁸⁸,
 R. Müller¹⁸, G.A. Mullier⁹⁵, J.J. Mullin¹²⁵, D.P. Mungo^{68a,68b}, J.L. Munoz Martinez¹²,
 F.J. Munoz Sanchez⁹⁸, M. Murin⁹⁸, P. Murin^{26b}, W.J. Murray^{163,131},
 A. Murrone^{68a,68b}, J.M. Muse¹¹⁷, M. Muškinja^{16a}, C. Mwewa²⁷, A.G. Myagkov^{35,a},
 A.J. Myers⁷, A.A. Myers¹²⁶, G. Myers⁶⁵, M. Myska¹²⁹, B.P. Nachman^{16a},
 O. Nackenhorst⁴⁷, A. Nag⁴⁸, K. Nagai¹²³, K. Nagano⁸⁰, J.L. Nagle²⁷, E. Nagy⁹⁹,
 A.M. Nairz³⁴, Y. Nakahama¹⁰⁸, K. Nakamura⁸⁰, H. Nanjo¹²¹, F. Napolitano^{61a},
 R. Narayan⁴², E.A. Narayanan¹⁰⁹, I. Naryshkin³⁵, M. Naseri³², C. Nass²²,
 T. Naumann⁴⁶, G. Navarro^{20a}, J. Navarro-Gonzalez¹⁵⁹, R. Nayak¹⁴⁸, P.Y. Nechaeva³⁵,
 F. Nechansky⁴⁶, T.J. Neep¹⁹, A. Negri^{70a,70b}, M. Negrini^{21b}, C. Nellist¹¹⁰,
 C. Nelson¹⁰¹, K. Nelson¹⁰³, S. Nemecek¹²⁸, M. Nessi^{34,g}, M.S. Neubauer¹⁵⁸,
 F. Neuhaus⁹⁷, J. Neundorff⁴⁶, R. Newhouse¹⁶⁰, P.R. Newman¹⁹, C.W. Ng¹²⁶, Y.S. Ng¹⁷,
 Y.W.Y. Ng¹⁵⁶, B. Ngair^{33e}, H.D.N. Nguyen¹⁰⁵, R.B. Nickerson¹²³, R. Nicolaidou¹³²,
 D.S. Nielsen⁴⁰, J. Nielsen¹³³, M. Niemeyer⁵³, N. Nikiforou¹⁰, V. Nikolaenko^{35,a},
 I. Nikolic-Audit¹²⁴, K. Nikolopoulos¹⁹, P. Nilsson²⁷, H.R. Nindhito⁵⁴, A. Nisati^{72a},
 N. Nishu², R. Nisius¹⁰⁷, T. Nitta¹⁶⁴, T. Nobe¹⁵⁰, D.L. Noel³⁰, Y. Noguchi⁸⁴,
 I. Nomidis¹²⁴, M.A. Nomura²⁷, M.B. Norfolk¹³⁶, R.R.B. Norisam⁹³, J. Novak⁹⁰,
 T. Novak⁴⁶, O. Novgorodova⁴⁸, L. Novotny¹²⁹, R. Novotny¹⁰⁹, L. Nozka¹¹⁹,
 K. Ntekas¹⁵⁶, E. Nurse⁹³, F.G. Oakham^{32,af}, J. Ocariz¹²⁴, A. Ochi⁸¹, I. Ochoa^{127a},
 J.P. Ochoa-Ricoux^{134a}, S. Oda⁸⁶, S. Odaka⁸⁰, S. Oerdek¹⁵⁷, A. Ogrodnik^{82a}, A. Oh⁹⁸,
 C.C. Ohm¹⁴¹, H. Oide¹⁵¹, R. Oishi¹⁵⁰, M.L. Ojeda⁴⁶, Y. Okazaki⁸⁴, M.W. O’Keefe⁸⁹,
 Y. Okumura¹⁵⁰, A. Olariu^{25b}, L.F. Oleiro Seabra^{127a}, S.A. Olivares Pino^{134e},
 D. Oliveira Damazio²⁷, D. Oliveira Goncalves^{79a}, J.L. Oliver¹⁵⁶, M.J.R. Olsson¹⁵⁶,
 A. Olszewski⁸³, J. Olszowska^{83,*}, Ö.O. Öncel²², D.C. O’Neil¹³⁹, A.P. O’Neill¹²³,
 A. Onofre^{127a,127e}, P.U.E. Onyisi¹⁰, R.G. Oreamuno Madriz¹¹², M.J. Oreglia³⁷,
 G.E. Orellana⁸⁷, D. Orestano^{74a,74b}, N. Orlando¹², R.S. Orr¹⁵², V. O’Shea⁵⁷,
 R. Ospanov^{60a}, G. Otero y Garzon²⁸, H. Otono⁸⁶, P.S. Ott^{61a}, G.J. Ottino^{16a},
 M. Ouchrif^{33d}, J. Ouellette²⁷, F. Ould-Saada¹²², A. Ouraou^{132,*}, Q. Ouyang^{13a},
 M. Owen⁵⁷, R.E. Owen¹³¹, K.Y. Oyulmaz^{11c}, V.E. Ozcan^{11c}, N. Ozturk⁷,
 S. Ozturk^{11c}, J. Pacalt¹¹⁹, H.A. Pacey³⁰, K. Pachal⁴⁹, A. Pacheco Pages¹²,
 C. Padilla Aranda¹², S. Pagan Griso^{16a}, G. Palacino⁶⁵, S. Palazzo⁵⁰, S. Palestini³⁴,
 M. Palka^{82b}, P. Palmi^{82a}, D.K. Panchal¹⁰, C.E. Pandini⁵⁴, J.G. Panduro Vazquez⁹²,
 P. Pani⁴⁶, G. Panizzo^{66a,66c}, L. Paolozzi⁵⁴, C. Papadatos¹⁰⁵, S. Parajuli⁴²,
 A. Paramonov⁵, C. Paraskevopoulos⁹, D. Paredes Hernandez^{62b}, S.R. Paredes Saenz¹²³,

B. Parida [ID](#)¹⁶⁵, T.H. Park [ID](#)¹⁵², A.J. Parker [ID](#)²⁹, M.A. Parker [ID](#)³⁰, F. Parodi [ID](#)^{55b,55a},
 E.W. Parrish [ID](#)¹¹², J.A. Parsons [ID](#)³⁹, U. Parzefall [ID](#)⁵², L. Pascual Dominguez [ID](#)¹⁴⁸,
 V.R. Pascuzzi [ID](#)^{16a}, F. Pasquali [ID](#)¹¹¹, E. Pasqualucci [ID](#)^{72a}, S. Passaggio [ID](#)^{55b}, F. Pastore [ID](#)⁹²,
 P. Pasuwan [ID](#)^{45a,45b}, J.R. Pater [ID](#)⁹⁸, A. Pathak [ID](#)¹⁶⁶, J. Patton⁸⁹, T. Pauly [ID](#)³⁴, J. Pearkes [ID](#)¹⁴⁰,
 M. Pedersen [ID](#)¹²², L. Pedraza Diaz [ID](#)¹¹⁰, R. Pedro [ID](#)^{127a}, T. Peiffer [ID](#)⁵³, S.V. Peleganchuk [ID](#)³⁵,
 O. Penc [ID](#)¹²⁸, C. Peng [ID](#)^{62b}, H. Peng [ID](#)^{60a}, M. Penzin [ID](#)³⁵, B.S. Peralva [ID](#)^{79a},
 A.P. Pereira Peixoto [ID](#)^{127a}, L. Pereira Sanchez [ID](#)^{45a,45b}, D.V. Perepelitsa [ID](#)²⁷,
 E. Perez Codina [ID](#)^{153a}, M. Perganti [ID](#)⁹, L. Perini [ID](#)^{68a,68b,*}, H. Pernegger [ID](#)³⁴, S. Perrella [ID](#)³⁴,
 A. Perrevoort [ID](#)¹¹¹, K. Peters [ID](#)⁴⁶, R.F.Y. Peters [ID](#)⁹⁸, B.A. Petersen [ID](#)³⁴, T.C. Petersen [ID](#)⁴⁰,
 E. Petit [ID](#)⁹⁹, V. Petousis [ID](#)¹²⁹, C. Petridou [ID](#)¹⁴⁹, P. Petroff⁶⁴, F. Petrucci [ID](#)^{74a,74b},
 A. Petrukhin [ID](#)¹³⁸, M. Pettee [ID](#)¹⁶⁸, N.E. Pettersson [ID](#)³⁴, K. Petukhova [ID](#)¹³⁰, A. Peyaud [ID](#)¹³²,
 R. Pezoa [ID](#)^{134f}, L. Pezzotti [ID](#)³⁴, G. Pezzullo [ID](#)¹⁶⁸, T. Pham [ID](#)¹⁰², P.W. Phillips [ID](#)¹³¹,
 M.W. Phipps [ID](#)¹⁵⁸, G. Piacquadio [ID](#)¹⁴², E. Pianori [ID](#)^{16a}, F. Piazza [ID](#)^{68a,68b}, A. Picazio [ID](#)¹⁰⁰,
 R. Piegai [ID](#)²⁸, D. Pietreanu [ID](#)^{25b}, J.E. Pilcher [ID](#)³⁷, A.D. Pilkington [ID](#)⁹⁸, M. Pinamonti [ID](#)^{66a,66c},
 J.L. Pinfold [ID](#)², C. Pitman Donaldson⁹³, D.A. Pizzi [ID](#)³², L. Pizzimento [ID](#)^{73a,73b}, A. Pizzini [ID](#)¹¹¹,
 M.-A. Pleier [ID](#)²⁷, V. Plesanovs⁵², V. Pleskot [ID](#)¹³⁰, E. Plotnikova³⁶, P. Podberezko [ID](#)³⁵,
 R. Poettgen [ID](#)⁹⁵, R. Poggi [ID](#)⁵⁴, L. Poggioli [ID](#)¹²⁴, I. Pogrebnnyak [ID](#)¹⁰⁴, D. Pohl [ID](#)²², I. Pokharel [ID](#)⁵³,
 G. Polesello [ID](#)^{70a}, A. Poley [ID](#)^{139,153a}, A. Policicchio [ID](#)^{72a,72b}, R. Polifka [ID](#)¹³⁰, A. Polini [ID](#)^{21b},
 C.S. Pollard [ID](#)¹²³, Z.B. Pollock [ID](#)¹¹⁶, V. Polychronakos [ID](#)²⁷, D. Ponomarenko [ID](#)³⁵,
 L. Pontecorvo [ID](#)³⁴, S. Popa [ID](#)^{25a}, G.A. Popeneciu [ID](#)^{25d}, L. Portales [ID](#)⁴,
 D.M. Portillo Quintero [ID](#)^{153a}, S. Pospisil [ID](#)¹²⁹, P. Postolache [ID](#)^{25c}, K. Potamianos [ID](#)¹²³,
 I.N. Potrap [ID](#)³⁶, C.J. Potter [ID](#)³⁰, H. Potti [ID](#)¹, T. Poulsen [ID](#)⁴⁶, J. Poveda [ID](#)¹⁵⁹, T.D. Powell [ID](#)¹³⁶,
 G. Pownall [ID](#)⁴⁶, M.E. Pozo Astigarraga [ID](#)³⁴, A. Prades Ibanez [ID](#)¹⁵⁹, P. Pralavorio [ID](#)⁹⁹,
 M.M. Prapa [ID](#)⁴⁴, S. Prell [ID](#)⁷⁸, D. Price [ID](#)⁹⁸, M. Primavera [ID](#)^{67a}, M.A. Principe Martin [ID](#)⁹⁶,
 M.L. Proffitt [ID](#)¹³⁵, N. Proklova [ID](#)³⁵, K. Prokofiev [ID](#)^{62c}, S. Protopopescu [ID](#)²⁷, J. Proudfoot [ID](#)⁵,
 M. Przybycien [ID](#)^{82a}, D. Pudzha [ID](#)³⁵, P. Puzo⁶⁴, D. Pyatiizbyantseva [ID](#)³⁵, J. Qian [ID](#)¹⁰³,
 Y. Qin [ID](#)⁹⁸, T. Qiu [ID](#)⁹¹, A. Quadt [ID](#)⁵³, M. Queitsch-Maitland [ID](#)³⁴, G. Rabanal Bolanos [ID](#)⁵⁹,
 F. Ragusa [ID](#)^{68a,68b}, J.A. Raine [ID](#)⁵⁴, S. Rajagopalan [ID](#)²⁷, K. Ran [ID](#)^{13a,13d}, D.F. Rassloff [ID](#)^{61a},
 D.M. Rauch [ID](#)⁴⁶, S. Rave [ID](#)⁹⁷, B. Ravina [ID](#)⁵⁷, I. Ravinovich [ID](#)¹⁶⁵, M. Raymond [ID](#)³⁴,
 A.L. Read [ID](#)¹²², N.P. Readioff [ID](#)¹³⁶, D.M. Rebuzzi [ID](#)^{70a,70b}, G. Redlinger [ID](#)²⁷, K. Reeves [ID](#)⁴³,
 D. Reikher [ID](#)¹⁴⁸, A. Reiss⁹⁷, A. Rej [ID](#)¹³⁸, C. Rembser [ID](#)³⁴, A. Renardi [ID](#)⁴⁶, M. Renda [ID](#)^{25b},
 M.B. Rendel¹⁰⁷, A.G. Rennie [ID](#)⁵⁷, S. Resconi [ID](#)^{68a}, M. Ressegotti [ID](#)^{55b,55a}, E.D. Resseguie [ID](#)^{16a},
 S. Rettie [ID](#)⁹³, B. Reynolds¹¹⁶, E. Reynolds [ID](#)¹⁹, M. Rezaei Estabragh [ID](#)¹⁶⁷, O.L. Rezanova [ID](#)³⁵,
 P. Reznicek [ID](#)¹³⁰, E. Ricci [ID](#)^{75a,75b}, R. Richter [ID](#)¹⁰⁷, S. Richter [ID](#)⁴⁶, E. Richter-Was [ID](#)^{82b},
 M. Ridel [ID](#)¹²⁴, P. Rieck [ID](#)¹⁰⁷, P. Riedler [ID](#)³⁴, O. Rifki [ID](#)⁴⁶, M. Rijssenbeek [ID](#)¹⁴²,
 A. Rimoldi [ID](#)^{70a,70b}, M. Rimoldi [ID](#)⁴⁶, L. Rinaldi [ID](#)^{21b,21a}, T.T. Rinn [ID](#)¹⁵⁸, M.P. Rinnagel [ID](#)¹⁰⁶,
 G. Ripellino [ID](#)¹⁴¹, I. Riu [ID](#)¹², P. Rivadeneira [ID](#)⁴⁶, J.C. Rivera Vergara [ID](#)¹⁶¹, F. Rizatdinova [ID](#)¹¹⁸,
 E. Rizvi [ID](#)⁹¹, C. Rizzi [ID](#)⁵⁴, B.A. Roberts [ID](#)¹⁶³, B.R. Roberts [ID](#)^{16a}, S.H. Robertson [ID](#)^{101,w},
 M. Robin [ID](#)⁴⁶, D. Robinson [ID](#)³⁰, C.M. Robles Gajardo^{134f}, M. Robles Manzano [ID](#)⁹⁷,
 A. Robson [ID](#)⁵⁷, A. Rocchi [ID](#)^{73a,73b}, C. Roda [ID](#)^{71a,71b}, S. Rodriguez Bosca [ID](#)^{61a},
 A. Rodriguez Rodriguez [ID](#)⁵², A.M. Rodríguez Vera [ID](#)^{153b}, S. Roe³⁴, A.R. Roepe-Gier [ID](#)¹¹⁷,
 J. Roggel [ID](#)¹⁶⁷, O. Røhne [ID](#)¹²², R.A. Rojas [ID](#)¹⁶¹, B. Roland [ID](#)⁵², C.P.A. Roland [ID](#)⁶⁵, J. Roloff [ID](#)²⁷,
 A. Romaniouk [ID](#)³⁵, M. Romano [ID](#)^{21b}, A.C. Romero Hernandez [ID](#)¹⁵⁸, N. Rompotis [ID](#)⁸⁹,

M. Ronzani [ID](#)¹¹⁴, L. Roos [ID](#)¹²⁴, S. Rosati [ID](#)^{72a}, B.J. Rosser [ID](#)¹²⁵, E. Rossi [ID](#)¹⁵², E. Rossi [ID](#)⁴,
 E. Rossi [ID](#)^{69a,69b}, L.P. Rossi [ID](#)^{55b}, L. Rossini [ID](#)⁴⁶, R. Rosten [ID](#)¹¹⁶, M. Rotaru [ID](#)^{25b}, B. Rottler [ID](#)⁵²,
 D. Rousseau [ID](#)⁶⁴, D. Rousso [ID](#)³⁰, G. Rovelli [ID](#)^{70a,70b}, A. Roy [ID](#)¹⁰, A. Rozanov [ID](#)⁹⁹, Y. Rozen [ID](#)¹⁴⁷,
 X. Ruan [ID](#)^{31f}, A.J. Ruby [ID](#)⁸⁹, T.A. Ruggeri [ID](#)¹, F. Rühr [ID](#)⁵², A. Ruiz-Martinez [ID](#)¹⁵⁹,
 A. Rummler [ID](#)³⁴, Z. Rurikova [ID](#)⁵², N.A. Rusakovich [ID](#)³⁶, H.L. Russell [ID](#)³⁴, L. Rustige [ID](#)³⁸,
 J.P. Rutherford [ID](#)⁶, E.M. Rüttinger [ID](#)¹³⁶, M. Rybar [ID](#)¹³⁰, E.B. Rye [ID](#)¹²², A. Ryzhov [ID](#)³⁵,
 J.A. Sabater Iglesias [ID](#)⁴⁶, P. Sabatini [ID](#)¹⁵⁹, L. Sabetta [ID](#)^{72a,72b}, H.F.W. Sadrozinski [ID](#)¹³³,
 F. Safai Tehrani [ID](#)^{72a}, B. Safarzadeh Samani [ID](#)¹⁴³, M. Safdari [ID](#)¹⁴⁰, S. Saha [ID](#)¹⁰¹,
 M. Sahinsoy [ID](#)¹⁰⁷, A. Sahu [ID](#)¹⁶⁷, M. Saimpert [ID](#)¹³², M. Saito [ID](#)¹⁵⁰, T. Saito [ID](#)¹⁵⁰, D. Salamani [ID](#)³⁴,
 G. Salamanna [ID](#)^{74a,74b}, A. Salnikov [ID](#)¹⁴⁰, J. Salt [ID](#)¹⁵⁹, A. Salvador Salas [ID](#)¹²,
 D. Salvatore [ID](#)^{41b,41a}, F. Salvatore [ID](#)¹⁴³, A. Salzburger [ID](#)³⁴, D. Sammel [ID](#)⁵², D. Sampsonidis [ID](#)¹⁴⁹,
 D. Sampsonidou [ID](#)^{60d,60c}, J. Sánchez [ID](#)¹⁵⁹, A. Sanchez Pineda [ID](#)⁴, V. Sanchez Sebastian [ID](#)¹⁵⁹,
 H. Sandaker [ID](#)¹²², C.O. Sander [ID](#)⁴⁶, I.G. Sanderswood [ID](#)⁸⁸, J.A. Sandesara [ID](#)¹⁰⁰,
 M. Sandhoff [ID](#)¹⁶⁷, C. Sandoval [ID](#)^{20b}, D.P.C. Sankey [ID](#)¹³¹, M. Sannino [ID](#)^{55b,55a}, A. Sansoni [ID](#)⁵¹,
 C. Santoni [ID](#)³⁸, H. Santos [ID](#)^{127a,127b}, S.N. Santpur [ID](#)^{16a}, A. Santra [ID](#)¹⁶⁵, K.A. Saoucha [ID](#)¹³⁶,
 J.G. Saraiva [ID](#)^{127a,127d}, J. Sardain [ID](#)⁹⁹, O. Sasaki [ID](#)⁸⁰, K. Sato [ID](#)¹⁵⁴, C. Sauer [ID](#)^{61b},
 F. Sauerburger [ID](#)⁵², E. Sauvan [ID](#)⁴, P. Savard [ID](#)^{152,af}, R. Sawada [ID](#)¹⁵⁰, C. Sawyer [ID](#)¹³¹,
 L. Sawyer [ID](#)⁹⁴, I. Sayago Galvan [ID](#)¹⁵⁹, C. Sbarra [ID](#)^{21b}, A. Sbrizzi [ID](#)^{21b,21a}, T. Scanlon [ID](#)⁹³,
 J. Schaarschmidt [ID](#)¹³⁵, P. Schacht [ID](#)¹⁰⁷, D. Schaefer [ID](#)³⁷, U. Schäfer [ID](#)⁹⁷, A.C. Schaffer [ID](#)⁶⁴,
 D. Schaile [ID](#)¹⁰⁶, R.D. Schamberger [ID](#)¹⁴², E. Schanet [ID](#)¹⁰⁶, C. Scharf [ID](#)¹⁷, N. Scharmberg [ID](#)⁹⁸,
 V.A. Schegelsky [ID](#)³⁵, D. Scheirich [ID](#)¹³⁰, F. Schenck [ID](#)¹⁷, M. Schernau [ID](#)¹⁵⁶, C. Schiavi [ID](#)^{55b,55a},
 L.K. Schildgen [ID](#)²², Z.M. Schillaci [ID](#)²⁴, E.J. Schioppa [ID](#)^{67a,67b}, M. Schioppa [ID](#)^{41b,41a},
 B. Schlag [ID](#)⁹⁷, K.E. Schleicher [ID](#)⁵², S. Schlenker [ID](#)³⁴, K. Schmieden [ID](#)⁹⁷, C. Schmitt [ID](#)⁹⁷,
 S. Schmitt [ID](#)⁴⁶, L. Schoeffel [ID](#)¹³², A. Schoening [ID](#)^{61b}, P.G. Scholer [ID](#)⁵², E. Schopf [ID](#)¹²³,
 M. Schott [ID](#)⁹⁷, J. Schovancova [ID](#)³⁴, S. Schramm [ID](#)⁵⁴, F. Schroeder [ID](#)¹⁶⁷, H-C. Schultz-Coulon [ID](#)^{61a},
 M. Schumacher [ID](#)⁵², B.A. Schumm [ID](#)¹³³, Ph. Schune [ID](#)¹³², A. Schwartzman [ID](#)¹⁴⁰,
 T.A. Schwarz [ID](#)¹⁰³, Ph. Schwemling [ID](#)¹³², R. Schwienhorst [ID](#)¹⁰⁴, A. Sciandra [ID](#)¹³³, G. Sciolla [ID](#)²⁴,
 F. Scuri [ID](#)^{71a}, F. Scutti [ID](#)¹⁰², C.D. Sebastiani [ID](#)⁸⁹, K. Sedlaczek [ID](#)⁴⁷, P. Seema [ID](#)¹⁷, S.C. Seidel [ID](#)¹⁰⁹,
 A. Seiden [ID](#)¹³³, B.D. Seidlitz [ID](#)²⁷, T. Seiss [ID](#)³⁷, C. Seitz [ID](#)⁴⁶, J.M. Seixas [ID](#)^{79b},
 G. Sekhniaidze [ID](#)^{69a}, S.J. Sekula [ID](#)⁴², L. Selem [ID](#)⁴, N. Semprini-Cesari [ID](#)^{21b,21a}, S. Sen [ID](#)⁴⁹,
 C. Serfon [ID](#)²⁷, L. Serin [ID](#)⁶⁴, L. Serkin [ID](#)^{66a,66b}, M. Sessa [ID](#)^{74a,74b}, H. Severini [ID](#)¹¹⁷, S. Sevova [ID](#)¹⁴⁰,
 F. Sforza [ID](#)^{55b,55a}, A. Sfyrla [ID](#)⁵⁴, E. Shabalina [ID](#)⁵³, R. Shaheen [ID](#)¹⁴¹, J.D. Shahinian [ID](#)¹²⁵,
 N.W. Shaikh [ID](#)^{45a,45b}, D. Shaked Renous [ID](#)¹⁶⁵, L.Y. Shan [ID](#)^{13a}, M. Shapiro [ID](#)^{16a}, A. Sharma [ID](#)³⁴,
 A.S. Sharma [ID](#)¹, S. Sharma [ID](#)⁴⁶, P.B. Shatalov [ID](#)³⁵, K. Shaw [ID](#)¹⁴³, S.M. Shaw [ID](#)⁹⁸,
 P. Sherwood [ID](#)⁹³, L. Shi [ID](#)⁹³, C.O. Shimmin [ID](#)¹⁶⁸, Y. Shimogama [ID](#)¹⁶⁴, J.D. Shinner [ID](#)⁹²,
 I.P.J. Shipsey [ID](#)¹²³, S. Shirabe [ID](#)⁵⁴, M. Shiyakova [ID](#)^{36,al}, J. Shlomi [ID](#)¹⁶⁵, M.J. Shochet [ID](#)³⁷,
 J. Shojaii [ID](#)¹⁰², D.R. Shope [ID](#)¹⁴¹, S. Shrestha [ID](#)¹¹⁶, E.M. Shrif [ID](#)^{31f}, M.J. Shroff [ID](#)¹⁶¹,
 E. Shulga [ID](#)¹⁶⁵, P. Sicho [ID](#)¹²⁸, A.M. Sickles [ID](#)¹⁵⁸, E. Sideras Haddad [ID](#)^{31f}, O. Sidiropoulou [ID](#)³⁴,
 A. Sidoti [ID](#)^{21b}, F. Siegert [ID](#)⁴⁸, Dj. Sijacki [ID](#)¹⁴, J.M. Silva [ID](#)¹⁹, M.V. Silva Oliveira [ID](#)³⁴,
 S.B. Silverstein [ID](#)^{45a}, S. Simion [ID](#)⁶⁴, R. Simoniello [ID](#)³⁴, N.D. Simpson [ID](#)⁹⁵, S. Simsek [ID](#)^{11b},
 P. Sinervo [ID](#)¹⁵², V. Sinetckii [ID](#)³⁵, S. Singh [ID](#)¹³⁹, S. Singh [ID](#)¹⁵², S. Sinha [ID](#)⁴⁶, S. Sinha [ID](#)^{31f},
 M. Sioli [ID](#)^{21b,21a}, I. Siral [ID](#)¹²⁰, S.Yu. Sivoklov [ID](#)^{35,*}, J. Sjölin [ID](#)^{45a,45b}, A. Skaf [ID](#)⁵³,
 E. Skorda [ID](#)⁹⁵, P. Skubic [ID](#)¹¹⁷, M. Slawinska [ID](#)⁸³, K. Sliwa [ID](#)¹⁵⁵, V. Smakhtin [ID](#)¹⁶⁵, B.H. Smart [ID](#)¹³¹,

J. Smiesko [ID](#)¹³⁰, S.Yu. Smirnov [ID](#)³⁵, Y. Smirnov [ID](#)³⁵, L.N. Smirnova [ID](#)^{35,a}, O. Smirnova [ID](#)⁹⁵,
 E.A. Smith [ID](#)³⁷, H.A. Smith [ID](#)¹²³, M. Smizanska [ID](#)⁸⁸, K. Smolek [ID](#)¹²⁹, A. Smykiewicz [ID](#)⁸³,
 A.A. Snesarev [ID](#)³⁵, H.L. Snoek [ID](#)¹¹¹, S. Snyder [ID](#)²⁷, R. Sobie [ID](#)^{161,w}, A. Soffer [ID](#)¹⁴⁸, F. Sohns [ID](#)⁵³,
 C.A. Solans Sanchez [ID](#)³⁴, E.Yu. Soldatov [ID](#)³⁵, U. Soldevila [ID](#)¹⁵⁹, A.A. Solodkov [ID](#)³⁵,
 S. Solomon [ID](#)⁵², A. Soloshenko [ID](#)³⁶, O.V. Solovyanov [ID](#)³⁵, V. Solovyev [ID](#)³⁵, P. Sommer [ID](#)¹³⁶,
 H. Son [ID](#)¹⁵⁵, A. Sonay [ID](#)¹², W.Y. Song [ID](#)^{153b}, A. Sopczak [ID](#)¹²⁹, A.L. Soppio [ID](#)⁹³, F. Sopkova [ID](#)^{26b},
 S. Sottocornola [ID](#)^{70a,70b}, R. Soualah [ID](#)^{66a,66c}, Z. Soumami [ID](#)^{33e}, D. South [ID](#)⁴⁶,
 S. Spagnolo [ID](#)^{67a,67b}, M. Spalla [ID](#)¹⁰⁷, M. Spangenberg [ID](#)¹⁶³, F. Spanò [ID](#)⁹², D. Sperlich [ID](#)⁵²,
 T.M. Spieker [ID](#)^{61a}, G. Spigo [ID](#)³⁴, M. Spina [ID](#)¹⁴³, D.P. Spiteri [ID](#)⁵⁷, M. Spousta [ID](#)¹³⁰,
 A. Stabile [ID](#)^{68a,68b}, R. Stamen [ID](#)^{61a}, M. Stamenkovic [ID](#)¹¹¹, A. Stampekis [ID](#)¹⁹, M. Standke [ID](#)²²,
 E. Stanecka [ID](#)⁸³, B. Stanislaus [ID](#)³⁴, M.M. Stanitzki [ID](#)⁴⁶, M. Stankaityte [ID](#)¹²³, B. Stapf [ID](#)⁴⁶,
 E.A. Starchenko [ID](#)³⁵, G.H. Stark [ID](#)¹³³, J. Stark [ID](#)^{99,aj}, D.M. Starko [ID](#)^{153b}, P. Staroba [ID](#)¹²⁸,
 P. Starovoitov [ID](#)^{61a}, S. Stärz [ID](#)¹⁰¹, R. Staszewski [ID](#)⁸³, G. Stavropoulos [ID](#)⁴⁴, P. Steinberg [ID](#)²⁷,
 A.L. Steinhebel [ID](#)¹²⁰, B. Stelzer [ID](#)^{139,153a}, H.J. Stelzer [ID](#)¹²⁶, O. Stelzer-Chilton [ID](#)^{153a},
 H. Stenzel [ID](#)⁵⁶, T.J. Stevenson [ID](#)¹⁴³, G.A. Stewart [ID](#)³⁴, M.C. Stockton [ID](#)³⁴, G. Stoicea [ID](#)^{25b},
 M. Stolarski [ID](#)^{127a}, S. Stonjek [ID](#)¹⁰⁷, A. Straessner [ID](#)⁴⁸, J. Strandberg [ID](#)¹⁴¹, S. Strandberg [ID](#)^{45a,45b},
 M. Strauss [ID](#)¹¹⁷, T. Strebler [ID](#)⁹⁹, P. Strizenc [ID](#)^{26b}, R. Ströhmer [ID](#)¹⁶², D.M. Strom [ID](#)¹²⁰,
 L.R. Strom [ID](#)⁴⁶, R. Stroynowski [ID](#)⁴², A. Strubig [ID](#)^{45a,45b}, S.A. Stucci [ID](#)²⁷, B. Stugu [ID](#)¹⁵,
 J. Stupak [ID](#)¹¹⁷, N.A. Styles [ID](#)⁴⁶, D. Su [ID](#)¹⁴⁰, S. Su [ID](#)^{60a}, W. Su [ID](#)^{60d,135,60c}, X. Su [ID](#)^{60a},
 K. Sugizaki [ID](#)¹⁵⁰, V.V. Sulim [ID](#)³⁵, M.J. Sullivan [ID](#)⁸⁹, D.M.S. Sultan [ID](#)⁵⁴, L. Sultanaliyeva [ID](#)³⁵,
 S. Sultansoy [ID](#)^{3c}, T. Sumida [ID](#)⁸⁴, S. Sun [ID](#)¹⁰³, S. Sun [ID](#)¹⁶⁶, X. Sun [ID](#)⁹⁸,
 O. Sunneborn Gudnadottir [ID](#)¹⁵⁷, C.J.E. Suster [ID](#)¹⁴⁴, M.R. Sutton [ID](#)¹⁴³, M. Svatos [ID](#)¹²⁸,
 M. Swiatlowski [ID](#)^{153a}, T. Swirski [ID](#)¹⁶², I. Sykora [ID](#)^{26a}, M. Sykora [ID](#)¹³⁰, T. Sykora [ID](#)¹³⁰, D. Ta [ID](#)⁹⁷,
 K. Tackmann [ID](#)^{46,u}, A. Taffard [ID](#)¹⁵⁶, R. Tafirout [ID](#)^{153a}, R.H.M. Taibah [ID](#)¹²⁴, R. Takashima [ID](#)⁸⁵,
 K. Takeda [ID](#)⁸¹, T. Takeshita [ID](#)¹³⁷, E.P. Takeva [ID](#)⁵⁰, Y. Takubo [ID](#)⁸⁰, M. Talby [ID](#)⁹⁹,
 A.A. Talyshev [ID](#)³⁵, K.C. Tam [ID](#)^{62b}, N.M. Tamer [ID](#)¹⁴⁸, A. Tanaka [ID](#)¹⁵⁰, J. Tanaka [ID](#)¹⁵⁰,
 R. Tanaka [ID](#)⁶⁴, J. Tang [ID](#)^{60c}, Z. Tao [ID](#)¹⁶⁰, S. Tapia Araya [ID](#)⁷⁸, S. Tapprogge [ID](#)⁹⁷,
 A. Tarek Abouelfadl Mohamed [ID](#)¹⁰⁴, S. Tarem [ID](#)¹⁴⁷, K. Tariq [ID](#)^{60b}, G. Tarna [ID](#)^{25b},
 G.F. Tartarelli [ID](#)^{68a}, P. Tas [ID](#)¹³⁰, M. Tasevsky [ID](#)¹²⁸, E. Tassi [ID](#)^{41b,41a}, G. Tateno [ID](#)¹⁵⁰,
 Y. Tayalati [ID](#)^{33e}, G.N. Taylor [ID](#)¹⁰², W. Taylor [ID](#)^{153b}, H. Teagle [ID](#)⁸⁹, A.S. Tee [ID](#)¹⁶⁶,
 R. Teixeira De Lima [ID](#)¹⁴⁰, P. Teixeira-Dias [ID](#)⁹², H. Ten Kate [ID](#)³⁴, J.J. Teoh [ID](#)¹¹¹, K. Terashi [ID](#)¹⁵⁰,
 J. Terron [ID](#)⁹⁶, S. Terzo [ID](#)¹², M. Testa [ID](#)⁵¹, R.J. Teuscher [ID](#)^{152,w}, N. Themistokleous [ID](#)⁵⁰,
 T. Theveneaux-Pelzer [ID](#)¹⁷, O. Thielmann [ID](#)¹⁶⁷, D.W. Thomas [ID](#)⁹², J.P. Thomas [ID](#)¹⁹,
 E.A. Thompson [ID](#)⁴⁶, P.D. Thompson [ID](#)¹⁹, E. Thomson [ID](#)¹²⁵, E.J. Thorpe [ID](#)⁹¹, Y. Tian [ID](#)⁵³,
 V. Tikhomirov [ID](#)^{35,a}, Yu.A. Tikhonov [ID](#)³⁵, S. Timoshenko [ID](#)³⁵, P. Tipton [ID](#)¹⁶⁸, S. Tisserant [ID](#)⁹⁹,
 S.H. Tlou [ID](#)^{31f}, A. Tmourji [ID](#)³⁸, K. Todome [ID](#)^{21b,21a}, S. Todorova-Nova [ID](#)¹³⁰, S. Todt [ID](#)⁴⁸,
 M. Togawa [ID](#)⁸⁰, J. Tojo [ID](#)⁸⁶, S. Tokár [ID](#)^{26a}, K. Tokushuku [ID](#)⁸⁰, E. Tolley [ID](#)¹¹⁶, R. Tombs [ID](#)³⁰,
 M. Tomoto [ID](#)^{80,108}, L. Tompkins [ID](#)^{140,ak}, P. Tornambe [ID](#)¹⁰⁰, E. Torrence [ID](#)¹²⁰, H. Torres [ID](#)⁴⁸,
 E. Torró Pastor [ID](#)¹⁵⁹, M. Toscani [ID](#)²⁸, C. Toscirì [ID](#)³⁷, J. Toth [ID](#)^{99,v}, D.R. Tovey [ID](#)¹³⁶, A. Traeet [ID](#)¹⁵,
 C.J. Treado [ID](#)¹¹⁴, T. Trefzger [ID](#)¹⁶², A. Tricoli [ID](#)²⁷, I.M. Trigger [ID](#)^{153a}, S. Trincaz-Duvoid [ID](#)¹²⁴,
 D.A. Trischuk [ID](#)¹⁶⁰, B. Trocmé [ID](#)⁵⁸, A. Trofymov [ID](#)⁶⁴, C. Troncon [ID](#)^{68a}, F. Trovato [ID](#)¹⁴³,
 L. Truong [ID](#)^{31c}, M. Trzebinski [ID](#)⁸³, A. Trzupek [ID](#)⁸³, F. Tsai [ID](#)¹⁴², A. Tsiamis [ID](#)¹⁴⁹,
 P.V. Tsiarehka [ID](#)^{35,a}, A. Tsirigotis [ID](#)^{149,s}, V. Tsiskaridze [ID](#)¹⁴², E.G. Tskhadadze [ID](#)^{146a},

M. Tsopoulou [ID](#)¹⁴⁹, Y. Tsujikawa [ID](#)⁸⁴, I.I. Tsukerman [ID](#)³⁵, V. Tsulaia [ID](#)^{16a}, S. Tsuno [ID](#)⁸⁰,
 O. Tsur [ID](#)¹⁴⁷, D. Tsybychev [ID](#)¹⁴², Y. Tu [ID](#)^{62b}, A. Tudorache [ID](#)^{25b}, V. Tudorache [ID](#)^{25b},
 A.N. Tuna [ID](#)³⁴, S. Turchikhin [ID](#)³⁶, I. Turk Cakir [ID](#)^{3a}, R.J. Turner [ID](#)¹⁹, R. Turra [ID](#)^{68a},
 P.M. Tuts [ID](#)³⁹, S. Tzamarias [ID](#)¹⁴⁹, P. Tzani [ID](#)⁹, E. Tzovara [ID](#)⁹⁷, K. Uchida [ID](#)¹⁵⁰, F. Ukegawa [ID](#)¹⁵⁴,
 P.A. Ulloa Poblete [ID](#)^{134c}, G. Unal [ID](#)³⁴, M. Unal [ID](#)¹⁰, A. Undrus [ID](#)²⁷, G. Unel [ID](#)¹⁵⁶,
 F.C. Ungaro [ID](#)¹⁰², K. Uno [ID](#)¹⁵⁰, J. Urban [ID](#)^{26b}, P. Urquijo [ID](#)¹⁰², G. Usai [ID](#)⁷, R. Ushioda [ID](#)¹⁵¹,
 M. Usman [ID](#)¹⁰⁵, Z. Uysal [ID](#)^{11d}, V. Vacek [ID](#)¹²⁹, B. Vachon [ID](#)¹⁰¹, K.O.H. Vadla [ID](#)¹²²,
 T. Vafeiadis [ID](#)³⁴, C. Valderanis [ID](#)¹⁰⁶, E. Valdes Santurio [ID](#)^{45a,45b}, M. Valente [ID](#)^{153a},
 S. Valentineti [ID](#)^{21b,21a}, A. Valero [ID](#)¹⁵⁹, R.A. Vallance [ID](#)¹⁹, A. Vallier [ID](#)^{99,aj},
 J.A. Valls Ferrer [ID](#)¹⁵⁹, T.R. Van Daalen [ID](#)¹³⁵, P. Van Gemmeren [ID](#)⁵, S. Van Stroud [ID](#)⁹³,
 I. Van Vulpen [ID](#)¹¹¹, M. Vanadia [ID](#)^{73a,73b}, W. Vandelli [ID](#)³⁴, M. Vandenbroucke [ID](#)¹³²,
 E.R. Vandewall [ID](#)¹¹⁸, D. Vannicola [ID](#)¹⁴⁸, L. Vannoli [ID](#)^{55b,55a}, R. Vari [ID](#)^{72a}, E.W. Varnes [ID](#)⁶,
 C. Varni [ID](#)^{16a}, T. Varol [ID](#)¹⁴⁵, D. Varouchas [ID](#)⁶⁴, K.E. Varvell [ID](#)¹⁴⁴, M.E. Vasile [ID](#)^{25b}, L. Vaslin [ID](#)³⁸,
 G.A. Vasquez [ID](#)¹⁶¹, F. Vazeille [ID](#)³⁸, D. Vazquez Furelos [ID](#)¹², T. Vazquez Schroeder [ID](#)³⁴,
 J. Veatch [ID](#)⁵³, V. Vecchio [ID](#)⁹⁸, M.J. Veen [ID](#)¹¹¹, I. Veliscek [ID](#)¹²³, L.M. Veloce [ID](#)¹⁵²,
 F. Veloso [ID](#)^{127a,127c}, S. Veneziano [ID](#)^{72a}, A. Ventura [ID](#)^{67a,67b}, A. Verbytskyi [ID](#)¹⁰⁷,
 M. Verducci [ID](#)^{71a,71b}, C. Vergis [ID](#)²², M. Verissimo De Araujo [ID](#)^{79b}, W. Verkerke [ID](#)¹¹¹,
 A.T. Vermeulen [ID](#)¹¹¹, J.C. Vermeulen [ID](#)¹¹¹, C. Vernieri [ID](#)¹⁴⁰, P.J. Verschuuren [ID](#)⁹²,
 M. Vessella [ID](#)¹⁰⁰, M.L. Vesterbacka [ID](#)¹¹⁴, M.C. Vetterli [ID](#)^{139,af}, A. Vgenopoulos [ID](#)¹⁴⁹,
 N. Viaux Maira [ID](#)^{134f}, T. Vickey [ID](#)¹³⁶, O.E. Vickey Boeriu [ID](#)¹³⁶, G.H.A. Viehhauser [ID](#)¹²³,
 L. Vigani [ID](#)^{61b}, M. Villa [ID](#)^{21b,21a}, M. Villaplana Perez [ID](#)¹⁵⁹, E.M. Villhauer [ID](#)⁵⁰, E. Vilucchi [ID](#)⁵¹,
 M.G. Vincter [ID](#)³², G.S. Virdee [ID](#)¹⁹, A. Vishwakarma [ID](#)⁵⁰, C. Vittori [ID](#)^{21b,21a}, I. Vivarelli [ID](#)¹⁴³,
 V. Vladimirov [ID](#)¹⁶³, E. Voevodina [ID](#)¹⁰⁷, M. Vogel [ID](#)¹⁶⁷, P. Vokac [ID](#)¹²⁹, J. Von Ahnen [ID](#)⁴⁶,
 E. Von Toerne [ID](#)²², V. Vorobel [ID](#)¹³⁰, K. Vorobev [ID](#)³⁵, M. Vos [ID](#)¹⁵⁹, J.H. Vosseveld [ID](#)⁸⁹,
 M. Vozak [ID](#)⁹⁸, L. Vozdecky [ID](#)⁹¹, N. Vranjes [ID](#)¹⁴, M. Vranjes Milosavljevic [ID](#)¹⁴, V. Vrba [ID](#)^{129,*},
 M. Vreeswijk [ID](#)¹¹¹, R. Vuillermet [ID](#)³⁴, O. Vujanovic [ID](#)⁹⁷, I. Vukotic [ID](#)³⁷, S. Wada [ID](#)¹⁵⁴,
 C. Wagner [ID](#)¹⁰⁰, W. Wagner [ID](#)¹⁶⁷, S. Wahdan [ID](#)¹⁶⁷, H. Wahlberg [ID](#)⁸⁷, R. Wakasa [ID](#)¹⁵⁴,
 M. Wakida [ID](#)¹⁰⁸, V.M. Walbrecht [ID](#)¹⁰⁷, J. Walder [ID](#)¹³¹, R. Walker [ID](#)¹⁰⁶, S.D. Walker [ID](#)⁹²,
 W. Walkowiak [ID](#)¹³⁸, A.M. Wang [ID](#)⁵⁹, A.Z. Wang [ID](#)¹⁶⁶, C. Wang [ID](#)^{60a}, C. Wang [ID](#)^{60c},
 H. Wang [ID](#)^{16a}, J. Wang [ID](#)^{62a}, P. Wang [ID](#)⁴², R.-J. Wang [ID](#)⁹⁷, R. Wang [ID](#)⁵⁹, R. Wang [ID](#)¹¹²,
 S.M. Wang [ID](#)¹⁴⁵, S. Wang [ID](#)^{60b}, T. Wang [ID](#)^{60a}, W.T. Wang [ID](#)⁷⁷, W.X. Wang [ID](#)^{60a}, X. Wang [ID](#)^{13c},
 X. Wang [ID](#)¹⁵⁸, X. Wang [ID](#)^{60c}, Y. Wang [ID](#)^{60a}, Z. Wang [ID](#)¹⁰³, C. Wanotayaroj [ID](#)³⁴,
 A. Warburton [ID](#)¹⁰¹, C.P. Ward [ID](#)³⁰, R.J. Ward [ID](#)¹⁹, N. Warrack [ID](#)⁵⁷, A.T. Watson [ID](#)¹⁹,
 M.F. Watson [ID](#)¹⁹, G. Watts [ID](#)¹³⁵, B.M. Waugh [ID](#)⁹³, A.F. Webb [ID](#)¹⁰, C. Weber [ID](#)²⁷,
 M.S. Weber [ID](#)¹⁸, S.A. Weber [ID](#)³², S.M. Weber [ID](#)^{61a}, C. Wei [ID](#)^{60a}, Y. Wei [ID](#)¹²³, A.R. Weidberg [ID](#)¹²³,
 J. Weingarten [ID](#)⁴⁷, M. Weirich [ID](#)⁹⁷, C. Weiser [ID](#)⁵², T. Wenaus [ID](#)²⁷, B. Wendland [ID](#)⁴⁷,
 T. Wengler [ID](#)³⁴, S. Wenig [ID](#)³⁴, N. Wermes [ID](#)²², M. Wessels [ID](#)^{61a}, K. Whalen [ID](#)¹²⁰,
 A.M. Wharton [ID](#)⁸⁸, A.S. White [ID](#)⁵⁹, A. White [ID](#)⁷, M.J. White [ID](#)¹, D. Whiteson [ID](#)¹⁵⁶,
 L. Wickremasinghe [ID](#)¹²¹, W. Wiedenmann [ID](#)¹⁶⁶, C. Wiel [ID](#)⁴⁸, M. Wielers [ID](#)¹³¹, N. Wieseotte [ID](#)⁹⁷,
 C. Wiglesworth [ID](#)⁴⁰, L.A.M. Wiik-Fuchs [ID](#)⁵², D.J. Wilbern [ID](#)¹¹⁷, H.G. Wilkens [ID](#)³⁴,
 L.J. Wilkins [ID](#)⁹², D.M. Williams [ID](#)³⁹, H.H. Williams [ID](#)¹²⁵, S. Williams [ID](#)³⁰, S. Willocq [ID](#)¹⁰⁰,
 P.J. Windischhofer [ID](#)¹²³, I. Wingerter-Seez [ID](#)⁴, F. Winklmeier [ID](#)¹²⁰, B.T. Winter [ID](#)⁵²,
 M. Wittgen [ID](#)¹⁴⁰, M. Wobisch [ID](#)⁹⁴, A. Wolf [ID](#)⁹⁷, R. Wölker [ID](#)¹²³, J. Wollrath [ID](#)¹⁵⁶, M.W. Wolter [ID](#)⁸³,

H. Wolters^{127a,127c}, V.W.S. Wong¹⁶⁰, A.F. Wongel⁴⁶, S.D. Worm⁴⁶, B.K. Wosiek⁸³, K.W. Woźniak⁸³, K. Wraight⁵⁷, J. Wu^{13a,13d}, S.L. Wu¹⁶⁶, X. Wu⁵⁴, Y. Wu^{60a}, Z. Wu^{132,60a}, J. Wuerzinger¹²³, T.R. Wyatt⁹⁸, B.M. Wynne⁵⁰, S. Xella⁴⁰, L. Xia^{13c}, M. Xia^{13b}, J. Xiang^{62c}, X. Xiao¹⁰³, M. Xie^{60a}, X. Xie^{60a}, I. Xioidis¹⁴³, D. Xu^{13a}, H. Xu^{60a}, H. Xu^{60a}, L. Xu^{60a}, R. Xu¹²⁵, T. Xu^{60a}, W. Xu¹⁰³, Y. Xu^{13b}, Z. Xu^{60b}, Z. Xu¹⁴⁰, B. Yabsley¹⁴⁴, S. Yacoob^{31a}, N. Yamaguchi⁸⁶, Y. Yamaguchi¹⁵¹, M. Yamatani¹⁵⁰, H. Yamauchi¹⁵⁴, T. Yamazaki^{16a}, Y. Yamazaki⁸¹, J. Yan^{60c}, S. Yan¹²³, Z. Yan²³, H.J. Yang^{60c,60d}, H.T. Yang^{16a}, S. Yang^{60a}, T. Yang^{62c}, X. Yang^{60a}, X. Yang^{13a}, Y. Yang¹⁵⁰, Z. Yang^{60a,103}, W-M. Yao^{16a}, Y.C. Yap⁴⁶, H. Ye^{13c}, J. Ye⁴², S. Ye²⁷, I. Yeletsikh³⁶, M.R. Yexley⁸⁸, P. Yin³⁹, K. Yorita¹⁶⁴, K. Yoshihara⁷⁸, C.J.S. Young⁵², C. Young¹⁴⁰, M. Yuan¹⁰³, R. Yuan^{60b,j}, X. Yue^{61a}, M. Zaazoua^{33e}, B. Zabinski⁸³, G. Zacharis⁹, E. Zaid⁵⁰, T. Zakareishvili^{146b}, N. Zakharchuk³², S. Zambito³⁴, D. Zanzi⁵², S.V. Zeibner⁴⁷, C. Zeitnitz¹⁶⁷, J.C. Zeng¹⁵⁸, D.T. Zenger Jr²⁴, O. Zenin³⁵, T. Ženiš^{26a}, S. Zenz⁹¹, S. Zerradi^{33a}, D. Zerwas⁶⁴, B. Zhang^{13c}, D.F. Zhang¹³⁶, G. Zhang^{13b}, J. Zhang⁵, K. Zhang^{13a,13d}, L. Zhang^{13c}, M. Zhang¹⁵⁸, R. Zhang¹⁶⁶, S. Zhang¹⁰³, X. Zhang^{60c}, X. Zhang^{60b}, Z. Zhang⁶⁴, P. Zhao⁴⁹, T. Zhao^{60b}, Y. Zhao¹³³, Z. Zhao^{60a}, A. Zhemchugov³⁶, Z. Zheng¹⁴⁰, D. Zhong¹⁵⁸, B. Zhou¹⁰³, C. Zhou¹⁶⁶, H. Zhou⁶, N. Zhou^{60c}, Y. Zhou⁶, C.G. Zhu^{60b}, C. Zhu^{13a,13d}, H.L. Zhu^{60a}, H. Zhu^{13a}, J. Zhu¹⁰³, Y. Zhu^{60a}, X. Zhuang^{13a}, K. Zhukov³⁵, V. Zhulanov³⁵, D. Zieminska⁶⁵, N.I. Zimine³⁶, S. Zimmermann^{52,*}, J. Zinsser^{61b}, M. Ziolkowski¹³⁸, L. Živković¹⁴, A. Zoccoli^{21b,21a}, K. Zoch⁵⁴, T.G. Zorbas¹³⁶, O. Zormpa⁴⁴, W. Zou³⁹, L. Zwalinski³⁴

¹ Department of Physics, University of Adelaide, Adelaide; Australia

² Department of Physics, University of Alberta, Edmonton AB; Canada

³ ^(a) Department of Physics, Ankara University, Ankara; ^(b) Istanbul Aydin University, Application and Research Center for Advanced Studies, Istanbul; ^(c) Division of Physics, TOBB University of Economics and Technology, Ankara; Türkiye

⁴ LAPP, Univ. Savoie Mont Blanc, CNRS/IN2P3, Annecy; France

⁵ High Energy Physics Division, Argonne National Laboratory, Argonne IL; United States of America

⁶ Department of Physics, University of Arizona, Tucson AZ; United States of America

⁷ Department of Physics, University of Texas at Arlington, Arlington TX; United States of America

⁸ Physics Department, National and Kapodistrian University of Athens, Athens; Greece

⁹ Physics Department, National Technical University of Athens, Zografou; Greece

¹⁰ Department of Physics, University of Texas at Austin, Austin TX; United States of America

¹¹ ^(a) Bahcesehir University, Faculty of Engineering and Natural Sciences, Istanbul; ^(b) Istanbul Bilgi University, Faculty of Engineering and Natural Sciences, Istanbul; ^(c) Department of Physics, Bogazici University, Istanbul; ^(d) Department of Physics Engineering, Gaziantep University, Gaziantep; Türkiye

¹² Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona; Spain

¹³ ^(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; ^(b) Physics Department, Tsinghua University, Beijing; ^(c) Department of Physics, Nanjing University, Nanjing; ^(d) University of Chinese Academy of Science (UCAS), Beijing; China

¹⁴ Institute of Physics, University of Belgrade, Belgrade; Serbia

¹⁵ Department for Physics and Technology, University of Bergen, Bergen; Norway

¹⁶ ^(a) Physics Division, Lawrence Berkeley National Laboratory, Berkeley CA; ^(b) University of California, Berkeley CA; United States of America

¹⁷ Institut für Physik, Humboldt Universität zu Berlin, Berlin; Germany

- ¹⁸ *Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern; Switzerland*
- ¹⁹ *School of Physics and Astronomy, University of Birmingham, Birmingham; United Kingdom*
- ²⁰ ^(a) *Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá;* ^(b) *Departamento de Física, Universidad Nacional de Colombia, Bogotá; Colombia*
- ²¹ ^(a) *Dipartimento di Fisica e Astronomia A. Righi, Università di Bologna, Bologna;* ^(b) *INFN Sezione di Bologna; Italy*
- ²² *Physikalisches Institut, Universität Bonn, Bonn; Germany*
- ²³ *Department of Physics, Boston University, Boston MA; United States of America*
- ²⁴ *Department of Physics, Brandeis University, Waltham MA; United States of America*
- ²⁵ ^(a) *Transilvania University of Brasov, Brasov;* ^(b) *Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest;* ^(c) *Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi;* ^(d) *National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj-Napoca;* ^(e) *University Politehnica Bucharest, Bucharest;* ^(f) *West University in Timisoara, Timisoara; Romania*
- ²⁶ ^(a) *Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava;* ^(b) *Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice; Slovak Republic*
- ²⁷ *Physics Department, Brookhaven National Laboratory, Upton NY; United States of America*
- ²⁸ *Universidad de Buenos Aires, Facultad de Ciencias Exactas y Naturales, Departamento de Física, y CONICET, Instituto de Física de Buenos Aires (IFIBA), Buenos Aires; Argentina*
- ²⁹ *California State University, CA; United States of America*
- ³⁰ *Cavendish Laboratory, University of Cambridge, Cambridge; United Kingdom*
- ³¹ ^(a) *Department of Physics, University of Cape Town, Cape Town;* ^(b) *iThemba Labs, Western Cape;* ^(c) *Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg;* ^(d) *National Institute of Physics, University of the Philippines Diliman (Philippines);* ^(e) *University of South Africa, Department of Physics, Pretoria;* ^(f) *School of Physics, University of the Witwatersrand, Johannesburg; South Africa*
- ³² *Department of Physics, Carleton University, Ottawa ON; Canada*
- ³³ ^(a) *Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies — Université Hassan II, Casablanca;* ^(b) *Faculté des Sciences, Université Ibn-Tofail, Kénitra;* ^(c) *Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech;* ^(d) *LPMR, Faculté des Sciences, Université Mohamed Premier, Oujda;* ^(e) *Faculté des sciences, Université Mohammed V, Rabat;* ^(f) *Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir; Morocco*
- ³⁴ *CERN, Geneva; Switzerland*
- ³⁵ *Affiliated with an institute covered by a cooperation agreement with CERN*
- ³⁶ *Affiliated with an international laboratory covered by a cooperation agreement with CERN*
- ³⁷ *Enrico Fermi Institute, University of Chicago, Chicago IL; United States of America*
- ³⁸ *LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand; France*
- ³⁹ *Nevis Laboratory, Columbia University, Irvington NY; United States of America*
- ⁴⁰ *Niels Bohr Institute, University of Copenhagen, Copenhagen; Denmark*
- ⁴¹ ^(a) *Dipartimento di Fisica, Università della Calabria, Rende;* ^(b) *INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati; Italy*
- ⁴² *Physics Department, Southern Methodist University, Dallas TX; United States of America*
- ⁴³ *Physics Department, University of Texas at Dallas, Richardson TX; United States of America*
- ⁴⁴ *National Centre for Scientific Research "Demokritos", Agia Paraskevi; Greece*
- ⁴⁵ ^(a) *Department of Physics, Stockholm University;* ^(b) *Oskar Klein Centre, Stockholm; Sweden*
- ⁴⁶ *Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen; Germany*
- ⁴⁷ *Fakultät Physik, Technische Universität Dortmund, Dortmund; Germany*
- ⁴⁸ *Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden; Germany*
- ⁴⁹ *Department of Physics, Duke University, Durham NC; United States of America*
- ⁵⁰ *SUPA — School of Physics and Astronomy, University of Edinburgh, Edinburgh; United Kingdom*

- ⁵¹ INFN e Laboratori Nazionali di Frascati, Frascati; Italy
- ⁵² Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany
- ⁵³ II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen; Germany
- ⁵⁴ Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland
- ⁵⁵ ^(a) Dipartimento di Fisica, Università di Genova, Genova; ^(b) INFN Sezione di Genova; Italy
- ⁵⁶ II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen; Germany
- ⁵⁷ SUPA — School of Physics and Astronomy, University of Glasgow, Glasgow; United Kingdom
- ⁵⁸ LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble; France
- ⁵⁹ Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA; United States of America
- ⁶⁰ ^(a) Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei; ^(b) Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao; ^(c) School of Physics and Astronomy, Shanghai Jiao Tong University, Key Laboratory for Particle Astrophysics and Cosmology (MOE), SKLPPC, Shanghai; ^(d) Tsung-Dao Lee Institute, Shanghai; China
- ⁶¹ ^(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; Germany
- ⁶² ^(a) Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, University of Hong Kong, Hong Kong; ^(c) Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong; China
- ⁶³ Department of Physics, National Tsing Hua University, Hsinchu; Taiwan
- ⁶⁴ IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91140, Orsay; France
- ⁶⁵ Department of Physics, Indiana University, Bloomington IN; United States of America
- ⁶⁶ ^(a) INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine; Italy
- ⁶⁷ ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce; Italy
- ⁶⁸ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano; Italy
- ⁶⁹ ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Fisica, Università di Napoli, Napoli; Italy
- ⁷⁰ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia; Italy
- ⁷¹ ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa; Italy
- ⁷² ^(a) INFN Sezione di Roma; ^(b) Dipartimento di Fisica, Sapienza Università di Roma, Roma; Italy
- ⁷³ ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma; Italy
- ⁷⁴ ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma; Italy
- ⁷⁵ ^(a) INFN-TIFPA; ^(b) Università degli Studi di Trento, Trento; Italy
- ⁷⁶ Universität Innsbruck, Department of Astro and Particle Physics, Innsbruck; Austria
- ⁷⁷ University of Iowa, Iowa City IA; United States of America
- ⁷⁸ Department of Physics and Astronomy, Iowa State University, Ames IA; United States of America
- ⁷⁹ ^(a) Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora; ^(b) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; ^(c) Universidade Federal de São João del Rei (UFSJ), São João del Rei; ^(d) Instituto de Física, Universidade de São Paulo, São Paulo; Brazil
- ⁸⁰ KEK, High Energy Accelerator Research Organization, Tsukuba; Japan
- ⁸¹ Graduate School of Science, Kobe University, Kobe; Japan
- ⁸² ^(a) AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow; ^(b) Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow; Poland
- ⁸³ Institute of Nuclear Physics Polish Academy of Sciences, Krakow; Poland
- ⁸⁴ Faculty of Science, Kyoto University, Kyoto; Japan
- ⁸⁵ Kyoto University of Education, Kyoto; Japan

- ⁸⁶ *Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka; Japan*
- ⁸⁷ *Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata; Argentina*
- ⁸⁸ *Physics Department, Lancaster University, Lancaster; United Kingdom*
- ⁸⁹ *Oliver Lodge Laboratory, University of Liverpool, Liverpool; United Kingdom*
- ⁹⁰ *Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana; Slovenia*
- ⁹¹ *School of Physics and Astronomy, Queen Mary University of London, London; United Kingdom*
- ⁹² *Department of Physics, Royal Holloway University of London, Egham; United Kingdom*
- ⁹³ *Department of Physics and Astronomy, University College London, London; United Kingdom*
- ⁹⁴ *Louisiana Tech University, Ruston LA; United States of America*
- ⁹⁵ *Fysiska institutionen, Lunds universitet, Lund; Sweden*
- ⁹⁶ *Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid; Spain*
- ⁹⁷ *Institut für Physik, Universität Mainz, Mainz; Germany*
- ⁹⁸ *School of Physics and Astronomy, University of Manchester, Manchester; United Kingdom*
- ⁹⁹ *CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille; France*
- ¹⁰⁰ *Department of Physics, University of Massachusetts, Amherst MA; United States of America*
- ¹⁰¹ *Department of Physics, McGill University, Montreal QC; Canada*
- ¹⁰² *School of Physics, University of Melbourne, Victoria; Australia*
- ¹⁰³ *Department of Physics, University of Michigan, Ann Arbor MI; United States of America*
- ¹⁰⁴ *Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America*
- ¹⁰⁵ *Group of Particle Physics, University of Montreal, Montreal QC; Canada*
- ¹⁰⁶ *Fakultät für Physik, Ludwig-Maximilians-Universität München, München; Germany*
- ¹⁰⁷ *Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München; Germany*
- ¹⁰⁸ *Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya; Japan*
- ¹⁰⁹ *Department of Physics and Astronomy, University of New Mexico, Albuquerque NM; United States of America*
- ¹¹⁰ *Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen; Netherlands*
- ¹¹¹ *Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam; Netherlands*
- ¹¹² *Department of Physics, Northern Illinois University, DeKalb IL; United States of America*
- ¹¹³ *(^a) New York University Abu Dhabi, Abu Dhabi; (^b) United Arab Emirates University, Al Ain; (^c) University of Sharjah, Sharjah; United Arab Emirates*
- ¹¹⁴ *Department of Physics, New York University, New York NY; United States of America*
- ¹¹⁵ *Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo; Japan*
- ¹¹⁶ *Ohio State University, Columbus OH; United States of America*
- ¹¹⁷ *Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK; United States of America*
- ¹¹⁸ *Department of Physics, Oklahoma State University, Stillwater OK; United States of America*
- ¹¹⁹ *Palacký University, Joint Laboratory of Optics, Olomouc; Czech Republic*
- ¹²⁰ *Institute for Fundamental Science, University of Oregon, Eugene, OR; United States of America*
- ¹²¹ *Graduate School of Science, Osaka University, Osaka; Japan*
- ¹²² *Department of Physics, University of Oslo, Oslo; Norway*
- ¹²³ *Department of Physics, Oxford University, Oxford; United Kingdom*
- ¹²⁴ *LPNHE, Sorbonne Université, Université Paris Cité, CNRS/IN2P3, Paris; France*
- ¹²⁵ *Department of Physics, University of Pennsylvania, Philadelphia PA; United States of America*
- ¹²⁶ *Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA; United States of America*
- ¹²⁷ *(^a) Laboratório de Instrumentação e Física Experimental de Partículas — LIP, Lisboa; (^b) Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa; (^c) Departamento de Física,*

- Universidade de Coimbra, Coimbra;^(d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa;^(e) Departamento de Física, Universidade do Minho, Braga;^(f) Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain);^(g) Departamento de Física, Instituto Superior Técnico, Universidade de Lisboa, Lisboa; Portugal
- ¹²⁸ Institute of Physics of the Czech Academy of Sciences, Prague; Czech Republic
- ¹²⁹ Czech Technical University in Prague, Prague; Czech Republic
- ¹³⁰ Charles University, Faculty of Mathematics and Physics, Prague; Czech Republic
- ¹³¹ Particle Physics Department, Rutherford Appleton Laboratory, Didcot; United Kingdom
- ¹³² IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette; France
- ¹³³ Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA; United States of America
- ¹³⁴ ^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago;^(b) Millennium Institute for Subatomic physics at high energy frontier (SAPHIR), Santiago;^(c) Instituto de Investigación Multidisciplinario en Ciencia y Tecnología, y Departamento de Física, Universidad de La Serena;^(d) Universidad Andres Bello, Department of Physics, Santiago;^(e) Instituto de Alta Investigación, Universidad de Tarapacá, Arica;^(f) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso; Chile
- ¹³⁵ Department of Physics, University of Washington, Seattle WA; United States of America
- ¹³⁶ Department of Physics and Astronomy, University of Sheffield, Sheffield; United Kingdom
- ¹³⁷ Department of Physics, Shinshu University, Nagano; Japan
- ¹³⁸ Department Physik, Universität Siegen, Siegen; Germany
- ¹³⁹ Department of Physics, Simon Fraser University, Burnaby BC; Canada
- ¹⁴⁰ SLAC National Accelerator Laboratory, Stanford CA; United States of America
- ¹⁴¹ Department of Physics, Royal Institute of Technology, Stockholm; Sweden
- ¹⁴² Departments of Physics and Astronomy, Stony Brook University, Stony Brook NY; United States of America
- ¹⁴³ Department of Physics and Astronomy, University of Sussex, Brighton; United Kingdom
- ¹⁴⁴ School of Physics, University of Sydney, Sydney; Australia
- ¹⁴⁵ Institute of Physics, Academia Sinica, Taipei; Taiwan
- ¹⁴⁶ ^(a) E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi;^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi;^(c) University of Georgia, Tbilisi; Georgia
- ¹⁴⁷ Department of Physics, Technion, Israel Institute of Technology, Haifa; Israel
- ¹⁴⁸ Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv; Israel
- ¹⁴⁹ Department of Physics, Aristotle University of Thessaloniki, Thessaloniki; Greece
- ¹⁵⁰ International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo; Japan
- ¹⁵¹ Department of Physics, Tokyo Institute of Technology, Tokyo; Japan
- ¹⁵² Department of Physics, University of Toronto, Toronto ON; Canada
- ¹⁵³ ^(a) TRIUMF, Vancouver BC;^(b) Department of Physics and Astronomy, York University, Toronto ON; Canada
- ¹⁵⁴ Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba; Japan
- ¹⁵⁵ Department of Physics and Astronomy, Tufts University, Medford MA; United States of America
- ¹⁵⁶ Department of Physics and Astronomy, University of California Irvine, Irvine CA; United States of America
- ¹⁵⁷ Department of Physics and Astronomy, University of Uppsala, Uppsala; Sweden
- ¹⁵⁸ Department of Physics, University of Illinois, Urbana IL; United States of America
- ¹⁵⁹ Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia — CSIC, Valencia; Spain
- ¹⁶⁰ Department of Physics, University of British Columbia, Vancouver BC; Canada
- ¹⁶¹ Department of Physics and Astronomy, University of Victoria, Victoria BC; Canada
- ¹⁶² Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg; Germany
- ¹⁶³ Department of Physics, University of Warwick, Coventry; United Kingdom

- ¹⁶⁴ *Waseda University, Tokyo; Japan*
- ¹⁶⁵ *Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot; Israel*
- ¹⁶⁶ *Department of Physics, University of Wisconsin, Madison WI; United States of America*
- ¹⁶⁷ *Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal; Germany*
- ¹⁶⁸ *Department of Physics, Yale University, New Haven CT; United States of America*
- ^a *Also Affiliated with an institute covered by a cooperation agreement with CERN*
- ^b *Also at Borough of Manhattan Community College, City University of New York, New York NY; United States of America*
- ^c *Also at Bruno Kessler Foundation, Trento; Italy*
- ^d *Also at Center for High Energy Physics, Peking University; China*
- ^e *Also at Centro Studi e Ricerche Enrico Fermi; Italy*
- ^f *Also at CERN, Geneva; Switzerland*
- ^g *Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève; Switzerland*
- ^h *Also at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona; Spain*
- ⁱ *Also at Department of Financial and Management Engineering, University of the Aegean, Chios; Greece*
- ^j *Also at Department of Physics and Astronomy, Michigan State University, East Lansing MI; United States of America*
- ^k *Also at Department of Physics and Astronomy, University of Louisville, Louisville, KY; United States of America*
- ^l *Also at Department of Physics, Ben Gurion University of the Negev, Beer Sheva; Israel*
- ^m *Also at Department of Physics, California State University, East Bay; United States of America*
- ⁿ *Also at Department of Physics, California State University, Fresno; United States of America*
- ^o *Also at Department of Physics, California State University, Sacramento; United States of America*
- ^p *Also at Department of Physics, King's College London, London; United Kingdom*
- ^q *Also at Department of Physics, University of Fribourg, Fribourg; Switzerland*
- ^r *Also at Faculty of Physics, Sofia University, 'St. Kliment Ohridski', Sofia; Bulgaria*
- ^s *Also at Hellenic Open University, Patras; Greece*
- ^t *Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona; Spain*
- ^u *Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg; Germany*
- ^v *Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest; Hungary*
- ^w *Also at Institute of Particle Physics (IPP); Canada*
- ^x *Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku; Azerbaijan*
- ^y *Also at Institute of Theoretical Physics, Ilia State University, Tbilisi; Georgia*
- ^z *Also at Instituto de Física Teórica, IFT-UAM/CSIC, Madrid; Spain*
- ^{aa} *Also at Istanbul University, Dept. of Physics, Istanbul; Türkiye*
- ^{ab} *Also at Physics Department, An-Najah National University, Nablus; Palestine*
- ^{ac} *Also at Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg; Germany*
- ^{ad} *Also at The City College of New York, New York NY; United States of America*
- ^{ae} *Also at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing; China*
- ^{af} *Also at TRIUMF, Vancouver BC; Canada*
- ^{ag} *Also at Università di Napoli Parthenope, Napoli; Italy*
- ^{ah} *Also at University of Chinese Academy of Sciences (UCAS), Beijing; China*
- ^{ai} *Also at Yeditepe University, Physics Department, Istanbul; Türkiye*
- ^{aj} *Also at L2IT, Université de Toulouse, CNRS/IN2P3, UPS, Toulouse, France*
- ^{ak} *Also at Department of Physics, Stanford University, Stanford CA, United States of America*
- ^{al} *Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria*
- ^{*} *Deceased*