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Key Points:

- Powerful new laboratory technique for resolving weak and overlapping spectra
- First measurement of the combined UV and visible spectra of CIOOCI
- Determination of the ClOOCI photolysis rate

Supporting Information:

- Readme
- Table S1
- Text S1Figure S1
- Figure S2
- Figure S3

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The UV and visible spectra of chlorine peroxide: Constraining the atmospheric photolysis rate

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Abstract The photolysis of chlorine peroxide (CIOOCI) is a key chemical step in the depletion of polar stratospheric ozone. As such, precise measurements of the absorption cross sections for CIOOCI are required. In this paper we provide two critical pieces of laboratory data with which to constrain the rate of ozone depletion. First, we provide an optically pure CIOOCI spectrum in the photolytically important UV wavelength region (200–350 nm). Second, we provide the first ever measurement of the CIOOCI spectrum in the visible region (509–536 nm) conclusively demonstrating the photolysis of CIOOCI to be negligible in this spectral region. The visible measurement is important because it can be used to anchor other previously recorded UV spectra and hence reduce the uncertainty in the CIOOCI photolysis rate. The calculated photolysis rate from this study indicates that the current models of stratospheric ozone depletion are correct.

1. Introduction

1.1. Chlorine Peroxide and Ozone Depletion

The springtime destruction of polar ozone (O₃) was first observed by *Farman et al.* [1985], and this destruction continues to the present day [World Meteorological Organization (WMO), 2011]. The link between ozone depletion and the presence of halogen, and in particular chlorine, containing gases in the stratosphere was quickly elucidated after the *Farman et al.* publication and ultimately led to the phasing out of the chlorine-containing chlorofluorocarbon gases through the 1989 Montreal Protocol and subsequent treaties [WMO, 2011].

The photolysis of the ClOOCI molecule is the rate-limiting step in the formation of the chlorine oxide species (ClOx = ClO + Cl), which are responsible for the majority of the ozone depletion in the stratospheric polar springtime [Solomon, 1999; $von\ Hobe\ et\ al.$, 2013]. During winter months, the stratosphere is dark and cold thereby allowing ClOOCI to form as a reservoir species of the more reactive ClOx which catalytically destroy O_3 [$Molina\ and\ Molina$, 1987]. The return of light in spring to the polar regions initiates the destruction of the photolabile ClOOCI thereby regenerating ClOx. The key reactions in the generation and loss of ClOOCI are shown below in reactions (R1-R3).

$$CIO + CIO + M \iff CIOOCI + M$$
 R1, R-1 (at low temperature RHS is favored.)
 $CIOOCI + hv \rightarrow CIOO + CI$ R2 (slow)
 $CIOO + M \rightarrow CI + O_2 + M$ R3 (rapid)

The photolysis rate of ClOOCI (J_{ClOOCI}) is calculated by equation (E1), where σ_{ClOOCI} is the absorption cross section, ϕ_{ClOOCI} is the photolysis quantum yield, F is the actinic flux, and λ is the wavelength. The quantum yield [$Huang\ et\ al.$, 2011; $Huang\ et\ al.$, 2011; $Huang\ et\ al.$, 2011; $Huang\ et\ al.$, 2004] has been sufficiently investigated through previous laboratory measurements as stated in $Huang\ et\ al.$ [2011]. The actinic fluxes can be modeled, both spatially and temporally, to a high accuracy (e.g., MYSTIC) [$Huang\ et\ al.$]

$$J_{\text{CIOOCI}} = \int \sigma_{\text{CIOOCI}}(\lambda) \phi_{\text{CIOOCI}}(\lambda) F(\lambda) d\lambda$$
 (E1)

1.2. Previous Spectral Determinations of ClOOCI

There have been numerous measurements of the ClOOCI absorption cross section, and these are shown in Figures S1 and S2 of the supporting information [Bloss et al., 2001; Burkholder et al., 1990; Chen et al., 2009; Cox and Hayman, 1988; DeMore and Tschuikow-Roux, 1990; Huder and DeMore, 1995; Jin et al., 2010; Lien et al., 2009; Papanastasiou et al., 2009; Pope et al., 2007; von Hobe et al., 2009; Wilmouth et al., 2009]. Large

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discrepancies still exist between the different laboratory determinations of the absorption cross sections of CIOOCI in the long-wavelength tail region. In particular, the study of Pope et al. [2007] was controversial because of its very low proposed values [Pope et al., 2007], which if correct, would indicate that the photolysis of CIOOCI was too slow to explain the generation of CIOx required to generate the observed polar ozone depletions [Schiermeier, 2007; von Hobe, 2007]. The Pope et al. study highlighted the lack of scientific consensus on the ClOOCI absorption cross section and renewed interest in the laboratory measurement of CIOOCI. Subsequently, six new studies measured the UV absorption cross sections using a variety of advanced techniques [Chen et al., 2009; Jin et al., 2010; Lien et al., 2009; Papanastasiou et al., 2009; von Hobe et al., 2009; Wilmouth et al., 2009]. However, within these studies, there are still significant differences between measurements at wavelengths longer than 300 nm (see Figure S2 of the supporting information). Currently, the Jet Propulsion Laboratory (JPL) kinetics evaluation [Sander et al., 2011] recommends the ClOOCI spectrum reported by Papanastasiou et al. [2009], recorded in the wavelength range 200-420 nm. At longer wavelengths, a log linear extrapolation of the Papanastasiou et al. data is recommended [Sander et al., 2011]. The Papanastasiou et al. study used diode array spectroscopy and isosbestic points to remove the contribution of the Cl₂ absorbance from the spectrum of the Cl₂/ClOOCl gas mixture [Papanastasiou et al., 2009].

One of the greatest difficulties in measuring the CIOOCI spectrum is generating pure CIOOCI within the laboratory. The synthetic routes employed involve precursors (e.g., Cl2, O3, and Cl2O) and generate coproducts (e.g., OCIO and Cl₂O₃) which are also absorbed in the UV/visible region (see Figure S3 of the supporting information). Distillation has been utilized as an effective method for removing most of the impurities [Pope et al., 2007]. However, Cl₂ is stubbornly resistant to complete removal because the vapor pressure is similar to CIOOCI. Furthermore, the CI₂ spectrum is smooth, broad, and similar to CIOOCI, making it problematic to unambiguously subtract from the CIOOCI spectrum. Several studies have employed indirect measures of the absorption cross section, which remove the importance of spectral impurities [Chen et al., 2009; Jin et al., 2010; Lien et al., 2009; Wilmouth et al., 2009]. Unfortunately, these studies rely on powerful lasers which can only be utilized at discrete wavelengths, and thus, the complete CIOOCI spectrum cannot be elucidated.

Prior to this study, the longest wavelength measurement of ClOOCI was 420 nm [Papanastasiou et al., 2009]. Theoretical calculations suggest that there is a possibility of additional electronic transitions occurring in the long-wavelength tail of the spectrum [Peterson and Francisco, 2004; Stanton and Bartlett, 1993]. Absorptions of low magnitude at long wavelengths could be significant because of the greater actinic flux in the visible region of the spectrum compared to the UV region. The recent stratospheric aircraft campaign carried out in the RECONCILE project [von Hobe et al., 2013] indirectly indicated that the long-wavelength tail of the CIOOCI spectrum should be photolytically negligible at wavelengths longer than 420 nm [Suminska-Ebersoldt et al., 2012].

2. Materials and Methods

2.1. Experimental Strategy

This study utilizes a twin-channel spectrometer [Young et al., 2011] to make absorption measurements of CIOOCI and CI₂ mixtures in both the UV (200-350 nm) and visible (509-536 nm) spectral regions. The UV channel utilizes conventional single-pass spectroscopy, and the visible channel exploits the highly sensitive Cavity-Enhanced Absorption Spectroscopy (CEAS) technique [Engeln et al., 1998; Langridge et al., 2008; O'Keefe, 1998] to measure the much weaker spectrum of Cl₂ in the visible part of the spectrum.

The experimental strategy was to generate mixtures of ClOOCI and Cl₂ using the methodology of *Pope et al.* [2007], then optimize the CIOOCI:Cl₂ ratio through distillation and temperature-dependent desorption of the mixture from a cold trap. The highly structured and well-characterized visible Cl₂ spectrum [Maric et al., 1993; Young et al., 2011] allows for the unambiguous measurement of the temporally evolving Cl₂ concentration.

Once the concentration of Cl₂ is known, the absorption due to Cl₂ in both spectral regions can be calculated, through knowledge of the UV and visible absorption cross sections, and subtracted from the composite ClOOCI/Cl₂ spectra resulting in pure ClOOCI spectra. The UV reference spectrum is taken from 2011 JPL kinetics evaluation [Sander et al., 2011], and the visible spectrum is taken from the study of Young et al. [2011]. This study did not measure the absolute concentration of ClOOCI within the absorption cell. The peak UV absorption, measured at 244 nm, is scaled to the corresponding peak absorption cross section recommended



by the 2010 JPL kinetics evaluation [Sander et al., 2011], which is based on the study of Papanastasiou et al. [2009]. Hence, the measurements reported herein are subject to the stated $\pm 35\%$ uncertainty of the JPL recommendation.

2.2. Experiment

The dual-beam spectrometer apparatus has been described in detail before [Young et al., 2011], and only a brief description is given here. The apparatus is composed of three distinct sequential sections: a photolysis cell in which ClOOCI is generated, a cold temperature trap in which to purify ClOOCI, and an absorption cell in which to measure ClOOCI and the optical impurities. All three components are temperature controlled. Throughout all stages of the experiment, the total flow through the photolysis cell, trap, and absorption cell was maintained at ~2 L/min and was controlled by mass flow controllers (MKS). A schematic diagram of the dual-beam spectrometer can be found in the study of Young et al. [2011].

Within the photolysis cell, ClOOCI is generated via the photolysis of Cl₂ in the presence of O₃. The reaction mechanism is shown in reactions (R4-R6), where M is any third body collision partner. The resulting mixture is composed of the following optically absorbing gases: CIOOCI, Cl_2 , O_3 , and trace amounts of OCIO and Cl_2O_3 . Higher oxides may also be present but were not spectroscopically identified.

$$Cl_2 + hv(UV) \rightarrow 2Cl$$
 (R4)

$$CI + O_3 \rightarrow CIO + O_2$$
 (R5)

$$CIO + CIO + M \rightarrow CIOOCI + M \tag{R6}$$

The cold trap was constructed from quartz, and its initial trapping temperature was set at ~150 K. This allowed for the capture of ClOOCI, Cl₂, and higher ClOx species while allowing O₃, which possesses a higher vapor pressure, to escape. Correspondingly, there was no spectroscopic evidence of codeposition of O₃ within the trap. Once sufficient ClOOCI has been collected, the trap is warmed to release the trapped gas into the absorption cell. The trap is slowly warmed, with a typical ramp rate of ~ 1 K min⁻¹, to an upper trap temperature which led to the release of ClOOCI and Cl₂ but no other spectral impurities. The OCIO and Cl₂O₃ species leave the trap at higher temperatures and thus were not observed under the experimental conditions.

The vapor pressure of Cl₂ is greater than ClOOCI, and therefore, Cl₂ leaves the cold trap at lower temperatures than ClOOCI. This physical difference between Cl₂ and ClOOCI was exploited to increase the concentration ratio of the desorbed ClOOCI:Cl₂ entering the absorption cell which is kept at a constant temperature of 197 K. Conditions were optimized to produce a desorption period which contained high ClOOCI concentrations (up to $\sim 4 \times 10^{12}$ cm⁻³) combined with low Cl₂ concentrations. Complete separation of the ClOOCI and Cl₂ is not achieved, and hence, all raw spectra contain some influence from the Cl₂ contaminant.

Both spectroscopic channels operate simultaneously and sample the same gas mixture. The combination of the visible and UV channels allows for the subtraction of the Cl₂ absorbance from both channels. In the UV channel, both Cl₂ and ClOOCI have similar shaped broad structureless spectra in the wavelength range ~300-360 nm (see Figure S3 of the supporting information) making it impossible to subtract the Cl₂ signal unambiguously without further information. This extra information is acquired from the visible spectrum, where the clear structure of Cl₂ (see Figure 1) allows for the determination of the Cl₂ concentration. The values of the absorption cross sections of the visible Cl₂ spectrum are temperature dependent; therefore, a pure reference spectrum is required and is provided by Young et al. [2011]. Knowledge of the Cl₂ concentration allows for the subtraction of the Cl₂ absorbance within both the visible and UV spectra resulting in the pure ClOOCI spectra.

Within the absorption cell, the UV absorbance (AUV) of the gas mixture is measured using single-pass spectroscopy and follows the Beer-Lambert relationship (for i different absorbers) given in equation (E2), where I and I₀ are the intensities of light transmitted through the cell with and without the absorbers present, respectively. The I is the path length, $\sigma(\lambda)$ is the wavelength-dependent absorption cross section, and c is the concentration of absorber.

$$A_{\text{UV}} = \ln \frac{I_0}{I} = \sum_i \sigma_i \ I \ c_i$$
 (E2)

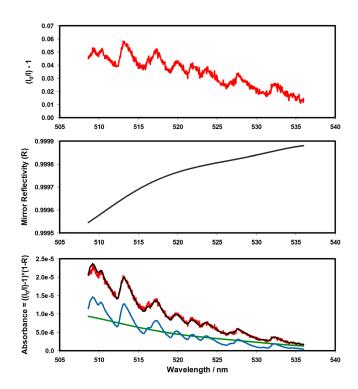


Figure 1. Cavity-enhanced absorption spectra of ClOOCI and Cl₂. (top) The laboratory-measured spectrum of a $CIOOCI + CI_2$ mixture of path integral concentrations of CIOOCI and $CI_2 = 4.1 \times 10^{16}$ cm⁻² and 1.6×10^{17} cm⁻², respectively. For a given wavelength: $(I_0/I)-1=\sum_i \sigma_i I c_i/(1-R)$. (middle) The BBCEAS mirror reflectivity obtained from a spectrum of pure Cl₂ of known concentration. (bottom) The absorbance spectra of the ClOOCI + Cl₂ mixture and individual ClOOCI and Cl₂ components: red line: laboratory-measured spectrum of ClOOCl + Cl₂ mixture, black line: fitted spectrum of CIOOCI + CI₂ mixture, blue line: fitted CI₂ spectrum, and green line: fitted CIOOCI spectrum.

The absorbance of the gas mixture in the visible channel has a much lower absorbance and hence is measured with the ultra sensitive BroadBand Cavity-Enhanced Absorption Spectroscopy (BBCEAS) technique. Figure 1 illustrates how the absorbance, due only to ClOOCI, is obtained from the visible spectrum. The wavelength-dependent absorbance of the ClOOCI and Cl₂ mixture can be calculated from the intensity of light transmitted through the BBCEAS channel with (I) and without (I_0) the absorbers present and the mirror reflectivity (R) via equation (E3) which is valid under conditions of low absorbance [Fiedler et al., 2003]. The BBCEAS mirror reflectivity, and hence

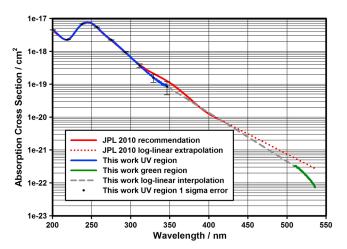


Figure 2. Visible and UV absorption cross sections of ClOOCI. The interpolation between the UV and green measurements from this study is log linear with boundary conditions provided by the longest wavelength UV measurement and shortest wavelength green measurement.

the effective path length for each experiment, was determined by recording a spectrum of pure Cl₂ of known concentration. We previously measured the temperature-dependent high-resolution absorption cross sections of Cl₂, as documented by Young et al. [2011], and these values are used to determine the effective mirror reflectivity for these experiments. The path integral concentration of Cl₂ (Cl₂ concentration multiplied by path length) is obtained from the UV measurement of Cl₂.

$$A_{vis} = \left(\frac{I_0}{I} - 1\right)(1 - R)$$
 (E3)

Once the absorbance spectrum of the Cl₂ and ClOOCI mixture is known, the relative contribution from Cl₂ can be

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Wavelength (nm)	Absorption Cross Section (cm ²)	1σ Standard Deviation (cm ²)
200	4.47E-18	2.39E-20
205	3.54E-18	2.08E-20
210	2.76E-18	1.86E-20
215	2.26E-18	1.66E-20
220	2.22E-18	1.46E-20
225	2.81E-18	1.30E-20
230	4.08E-18	1.43E-20
235	5.71E-18	1.96E-20
240	7.08E-18	2.58E-20
245	7.58E-18	2.97E-20
250	7.02E-18	3.00E-20
255	5.79E-18	2.86E-20
260	4.48E-18	2.79E-20
265	3.43E-18	2.88E-20
270	2.70E-18	3.10E-20
275	2.17E-18	3.39E-20
280	1.75E-18	3.69E-20
285	1.38E-18	3.96E-20
290	1.09E-18	4.08E-20
295	8.37E-19	4.22E-20
300	6.50E-19	4.32E-20
305	5.06E-19	4.42E-20
310	3.92E-19	4.49E-20
315	3.05E-19	4.51E-20
320	2.38E-19	4.44E-20
325	1.84E-19	4.34E-20
330	1.48E-19	4.22E-20
335	1.25E-19	4.06E-20
340	9.98E-20	3.85E-20
345	9.12E-20	3.65E-20
510	3.21E-22	1.80E-23
515	2.77E-22	1.64E-23
520	2.12E-22	1.50E-23
525	1.61E-22	1.35E-23
530	1.19E-22	1.17E-23
535	8.20E-23	9.15E-24

^aHigh-resolution data interpolated onto a 5 nm grid using air wavelengths. The standard deviation does not take into account the uncertainty on the normalized peak cross section.

easily subtracted because of its highly structured spectrum in the wavelength region 509-535 nm. At longer wavelengths, the amount of structure diminishes; hence, the retrieval error increases, and as such, we only report the spectrum for the highly structured region. This retrieval of the Cl₂ absorbance provides the Cl₂ concentration, since the absorption cross sections and path length are known, which can then be used to subtract the correct amount of Cl₂ from the visible spectrum. The residual spectrum, that is left once the Cl₂ contribution has been subtracted from the composite (ClOOCl + Cl₂) spectrum, consistently possesses the same shape that is well fitted to a polynomial of order 4. The chronology of the appearance and the changing magnitude of the residual spectrum are synchronous with the UV spectrum of CIOOCI. Therefore, we confidently attribute the residual spectrum to be the visible absorbance spectrum of ClOOCI. The absolute absorption cross sections are obtained by normalizing the absorbance spectrum with the same factor used for the UV spectrum. This approach is valid because the same gas is being sampled in both the UV and visible channels.

The gases used in this study were used without further purification: 5% Cl₂/N₂ (BOC), O₂ (BOC), and N₂ (BOC). The O_3 was generated from the photolysis of synthetic air using a mercury lamp (Penray lamp).

3. Results and Discussion

The UV and visible absorption cross sections of pure ClOOCI are shown in Figure 2 and are compared to the most recent 2010 JPL recommendation [Sander et al., 2011]. Table 1 provides the averages and the 1σ

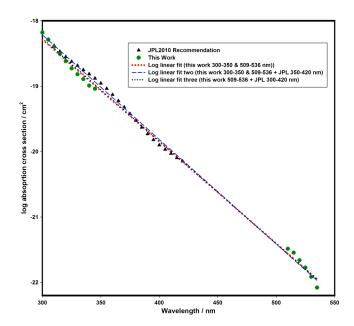


Figure 3. ClOOCI absorption cross sections in the photolytically important wavelength region (λ > 300 nm). The JPL recommendation is based upon the work of *Papanastasiou et al.* [2009]. Three different interpolations are used to fill the data gap between the laboratory measurements: interpolation 1 (red dotted line) uses only the UV and visible data from this study leaving a data gap between 350 and 509 nm. Interpolation 2 uses the UV and visible data from this study in addition to the *Papanastasiou et al.* data in the range 350–420 nm leaving a data gap between 420 and 509 nm. Interpolation 3 uses the UV data of *Papanastasiou et al.* and the visible data from this study leaving a data gap between 420 and 509 nm.

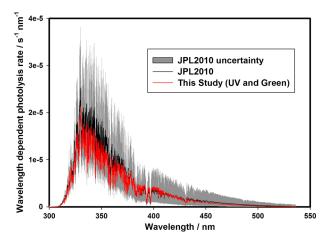


Figure 4. Wavelength-dependent ClOOCI photolysis rates in the polar stratosphere. The photolysis rates are calculated using equation (E1). The quantum yield (Φ) for reaction (R2) is set equal to 0.8 as recommended by the JPL kinetics evaluation [Sander et al., 2011]. High-resolution actinic flux (F) data are obtained from the output of the MYSTIC model used with the following model parameters: altitude = 20 km, no aerosol or clouds, ozone concentration = 350 Dobson unit, surface albedo = 0.3, and solar zenith angle = 85°[Mayer, 2009]. MYSTIC is operated as one of several solvers of the libRadtran radiative transfer package [Mayer and Kylling, 2005]. The spectra were calculated in a fully spherical geometry [Emde and Mayer, 2007] using the absorption line importance sampling method [Emde et al., 2011]. The JPL 2010 recommendation [Sander et al., 2011] for the absorption cross sections is based upon the laboratory work of Papanastasiou et al. [2009] and uses a log linear extrapolation at wavelengths longer than 420 nm. The uncertainty on the JPL recommendation represents the consensus of the panel members on JPL kinetics evaluation committee [Sander et al., 2011]. The absorption cross sections from this work are the interpolated data from the combined UV and green data sets (blue, pink, and green lines in Figure 2).



standard deviations for both the UV and visible ClOOCI absorption cross sections interpolated to a resolution of 5 nm. Higher-resolution (0.5 nm) spectral data are provided in the supporting information.

The measured UV ClOOCI spectrum is slightly weaker, at wavelengths long of 300 nm, than the JPLrecommended Papanastasiou et al. [2009] data but falls within its stated uncertainty. This study provides the first measurement of the visible CIOOCI spectrum. The spectrum is found to be broad and unstructured; it rapidly reduces in intensity with increasing wavelength. A log linear interpolation between the UV and visible data sets is shown in Figure 2. This combination of the UV and visible data suggests that the absorption cross section of CIOOCI follows an approximately log linear relationship with wavelength at wavelengths longer than 300 nm. This result is consistent with the absence of additional nonnegligible electronic transitions occurring at wavelengths longer than 300 nm. However, it should be noted that this study cannot rule out additional absorption features appearing in the ~159 nm data gap between 350 and 509 nm.

The good agreement, within errors, between this study and the study of Papanastasiou et al. [2009], which extends the long UV measurement to 420 nm, allows for the combination of data sets to reduce the laboratory data gap to 89 nm (420-509 nm), which further constrains the ClOOCI photolysis rate. Figure 3 highlights the measured and interpolated absorption spectra of ClOOCI in the photolytically important wavelength region. It can be seen that the visible measurement, provided by this study, effectively anchors the data and removes much of the uncertainty within the long-wavelength tail of ClOOCI assuming a log linear relationship holds over the whole wavelength range. This is consistent with the RECONCILE field measurements which suggested that it is unlikely that any photolytically significant absorption will exist at wavelengths longer than 420 nm [Suminska-Ebersoldt et al., 2012]. However, even a small additional peak ($>1 \times 10^{-20}$ cm²) in the unmeasured 420–509 nm region could significantly increase the photolysis rate due to the high actinic flux present at these wavelengths.

Figure 4 compares the wavelength-dependent photolysis rates, relevant for the polar arctic stratospheric spring, derived from the interpolated absorption cross section data measured in this study and the recommended JPL evaluation values including the recommended extrapolation at wavelengths longer than 420 nm. The calculated photolysis rate for CIOOCI integrated over all relevant UV and visible wavelengths detailed in Figure 3 is found to be 7.3×10^{-4} s⁻¹, approximately 22% lower than the current JPL recommendation, but it is well within the stated uncertainty of the JPL recommendation [Sander et al., 2011]. The absorption cross sections measured in the green region of the spectrum are found to be too small to be photolytically important. However, as previously stated, their major value is to anchor the UV data measurements and hence provide a powerful constraint upon the overall photolysis rate of CIOOCI.

4. Conclusions

The photolysis of chlorine peroxide (CIOOCI) is a key chemical step in the depletion of polar stratospheric ozone. This study provides an optically pure CIOOCI spectrum in the photolytically important UV wavelength region (200–350 nm) that is in good agreement with other recently published spectra. Second, the study provides the first ever measurement of the ClOOCI spectrum in the visible region (509–536 nm) conclusively demonstrating the photolysis of ClOOCI to be negligible in this spectral region. The visible measurement can be used to anchor other previously recorded UV spectra and hence reduce the uncertainty in CIOOCI photolysis rate calculations. The calculated photolysis rate from this study indicates that current models of stratospheric ozone depletion are correct. Furthermore, the new visible spectrum measurement greatly improves the confidence in this result.

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