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# 1 Windborne debris trajectories in tornado-like flow field 2 initiated from a low-rise building

3 Frederick Bourriez <sup>a</sup>, Mark Sterling <sup>a</sup>, Chris Baker <sup>a</sup>

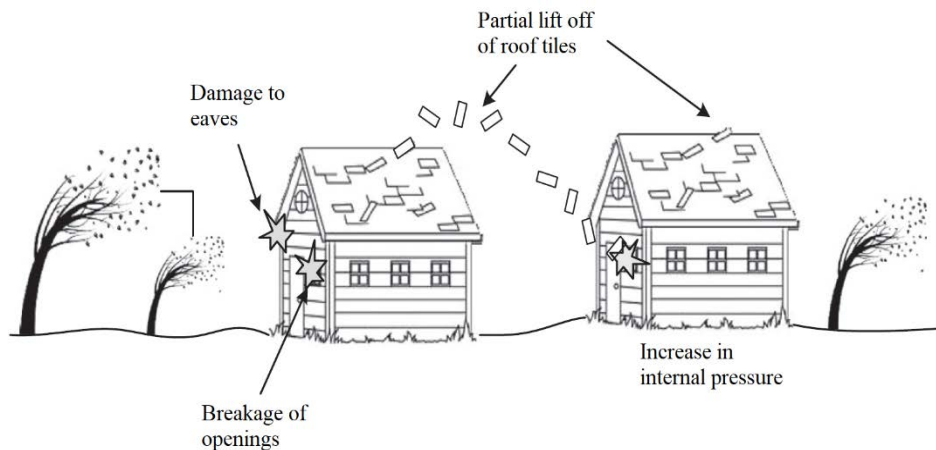
4  
5 <sup>a</sup> *University of Birmingham, School of Engineering, Edgbaston, Birmingham, U.K.*

6 ABSTRACT: This paper examines compact debris flight in tornado-like flow fields. The  
7 research focuses on physically simulating a specific tornado-like vortex and on investigating  
8 windborne debris flight with and without a low-rise building model. The low-rise building  
9 model, 1/8<sup>th</sup> scale with regard to the vortex core radius, was used to initiate the flight of  
10 Styrofoam spheres from its top. Debris motion was recorded using three high-speed cameras  
11 and data reduction was performed on open-source OpenPTV software. Flow field  
12 characterisation showed that including a building model does not considerably affect the  
13 averaged flow field but only the local instantaneous flow field. Debris flight analysis shows  
14 that the mean flight distance is not affected by the building model, but a change in the initial  
15 direction occurs. Comparison between local wind flow field and initial debris velocities shows  
16 good agreement, and therefore the variability in initial directions of the debris flight can be  
17 attributed to a wind-driven process. To compare with experimental data, experimental data  
18 were incorporated into debris flight equations to compute debris motion. Debris trajectories  
19 computed from experimental data show strong visual similarities with experimental trajectories  
20 and debris flight analysis presents good agreement with experimental data.

21 KEYWORDS: tornado-like vortex, debris, low-rise building, 3D-PTV.

## 22 1 INTRODUCTION

23  
24 Tornadoes are fast vertical swirling columns of air formed inside a storm and connecting the  
25 cloud base and the ground via a funnel cloud. Tornadoes account for the strongest and most  
26 devastating natural wind phenomenon with wind speeds up to 480 km/h. When fully developed,  
27 the width of a tornado can reach half a mile in diameter and travel long distances. Moreover,  
28 the debris cloud during a tornadic event, observable when dust and/or objects are lifted aloft  
29 and are swirling around the tornado, can account for twice the size of the tornado itself  
30 (Wurman *et al.*, 2013). In 2018, 628 tornadoes were reported in Europe. In the United  
31 Kingdom, about 30 tornadoes are reported every year, with 2.2 tornadoes per year per 10,000  
32 km sq. on average. This is more than in the U.S. with 1.3 tornadoes per year per 10,000km sq.  
33 and about 1,200 tornadoes occurring every year (Mulder and Schultz, 2015). Nevertheless,  
34 tornadoes in the U.S. are stronger than tornadoes usually occurring in Europe. As a result, over  
35 the last 40 years, about 60 fatalities were reported and billions of dollars of damages are caused  
36 every year on U.S. soil. Damage to property and infrastructure due to tornadoes can mainly be  
37 attributed to the high wind speeds, the low pressure inside the vortex core and windborne debris  
38 (Brooks and Doswell, 2001). Windborne debris embedded during strong wind events, such as  
39 downbursts, hurricanes or tornadoes, can lead to major wind disasters. When reaching high  
40 velocities, debris can significantly damage engineered and non-engineered buildings, resulting  
41 in the production of even more debris, known as debris chain (Figure 1).



**Figure 1** Windborne debris chain (from Kakimpa et al. (2012))

42 Due to their unpredictability and their danger, full-scale tornadoes are difficult to study,  
 43 therefore it is particularly challenging to obtain in-situ velocity and pressure data. Using  
 44 Doppler radar measurements, (Wurman and Alexander, 2005) collected wind field  
 45 measurements of the Spencer South Dakota tornado of May 30, 1998. However, due to the  
 46 earth's curvature and Doppler radar elevation angle, no data are usually obtained below 30m  
 47 AGL (Wurman *et al.*, 2013), which is usually more than standard low-rise building heights. To  
 48 overcome this limitation, Bluestein and Unruh (1989) used a Portable Doppler radar to  
 49 intercept and get as close as 10km distance from a tornado. As a result, between 1995 and 2008,  
 50 150 tornadoes have been measured using Doppler On Wheels (DOWs) (Alexander and  
 51 Wurman, 2008) with measurements within 20m Above Ground Level (AGL) (Kosiba and  
 52 Wurman, 2013).

53 Recourse to physical modelling of tornadoes has had a significant impact in our understanding  
 54 of the processes governing tornado-like flow fields. Ward (1972) pioneered the tornado  
 55 generator that could generate realistic tornado-like vortices and sub-vortices. Ward-type  
 56 generators use guide vanes to introduce angular momentum and an exhaust fan to generate an  
 57 updraft. Several Ward-type generators were then developed to further investigate tornado-like  
 58 vortices in small-scale (Church *et al.*, 1977), small/medium scale (Gillmeier *et al.*, 2018) and  
 59 large-scale generators (Tang *et al.*, 2018). Over the past decades, tornado generators with new  
 60 design have been developed: Mishra et al. (2008a; 2008b) developed a Ward-type generator  
 61 where the circulation is driven by 16 slotted jets; (Haan *et al.*, 2008) developed a new large-  
 62 scale translating generator with a central exhaust fan and an annular duct recirculating the flow  
 63 downward; Refan et al. (2014) used the new large-scale 3D wind testing facility (WindEEE  
 64 dome) to generate tornado-like vortices. Through the years and improvements, each generator  
 65 presented improving similarities with full-scale tornadoes. The pressure drop in the vortex core  
 66 matched with the 2004 Manchester (U.S.) tornado (Mishra *et al.*, 2008a) and velocity  
 67 components showed good agreements with full-scale measurements (Haan *et al.*, 2008).  
 68 Tornado-like vortex generators can also model one and two-cell vortices (Haan *et al.*, 2008;  
 69 Tang *et al.*, 2018), and improved the understanding of the vortex breakdown during the  
 70 transition between one and two-cell structure (Church and Snow, 1985). However, Baker and  
 71 Sterling (2019) recently questioned the efficacy of tornado simulators to reproduce full-scale  
 72 tornadoes. The authors looked at various dimensionless group to evaluate the performance of  
 73 several tornado simulators with regards to dynamic, kinematic and geometric similarities. The

74 main conclusion drawn from this analysis is that some are able to match geometric similarity  
75 for some tornadoes, some are able to match kinematic similarity but none of the current tornado  
76 simulators are capable to show all similarities with full-scale tornadoes and therefore have  
77 limited capabilities.

78 Different analytical vortex models have been developed to reproduce tornado-like flow field  
79 (Rankine, 1882; Burgers, 1948; Rott, 1958; Sullivan, 1959; Baker and Sterling, 2017).  
80 However, due to the assumptions made and the complexity of tornado flow fields, some vortex  
81 models tend to only reproduce the tangential velocity component correctly and remain  
82 incapable of realistically modelling the radial and axial components. Gillmeier et al. (2018)  
83 compared vortex models with physically modelled flow fields and showed that Baker and  
84 Sterling's model is currently the most realistic on those cited above.

85 Debris flight models have also been developed, although they were originally simplified and  
86 only considered the drag forces (McDonald *et al.*, 1974) or the drag and lift forces (Lee, 1974).  
87 Subsequently, effort has been made to understand the forces involved during the flight of  
88 debris. As a result, Twisdale et al. (1979) performed an analysis of tornado missile transport  
89 using a 'random orientation 6-degree of freedom' 3D model including drag, lift and side forces;  
90 Tachikawa (1983, 1988) looked at the flight of flat plates in a wind-tunnel, resulting in the  
91 definition of the Tachikawa number (Holmes *et al.*, 2006) and accounting for the ratio of  
92 aerodynamic to gravity forces; Holmes (2004) and Holmes et al. (2006) developed debris flight  
93 equations for spherical and plate-type debris, respectively; Kordi and Kopp (2011) investigated  
94 the flight of windborne plate debris from a building in a wind-tunnel; Baker (2007) developed  
95 debris flight equations with a dimensionless approach. In a further paper, Baker and Sterling  
96 (2017) developed the debris flight equations for tornado-like wind field applications.

97 Although great effort has been undertaken over the past few decades to model and understand  
98 the flow field of a tornado, the flight of debris in a tornado flow field is still poorly understood.  
99 It is therefore not surprising that only a little can be found in the literature focusing on  
100 windborne debris flight in a tornado-like vortices (Sassa *et al.*, 2009; Maruyama, 2011; Noda  
101 *et al.*, 2013; Baker and Sterling, 2017). The current work presents an introductory overview of  
102 the experimental investigation of windborne debris flight in tornado-like flow field using a  
103 tornado generator and the Particle Tracking Velocimetry (PTV) technique and a methodology  
104 to compute debris trajectories using experimental data.

## 105 2 DEBRIS FLIGHT INVESTIGATION

### 106 2.1 *Experimental methodology*

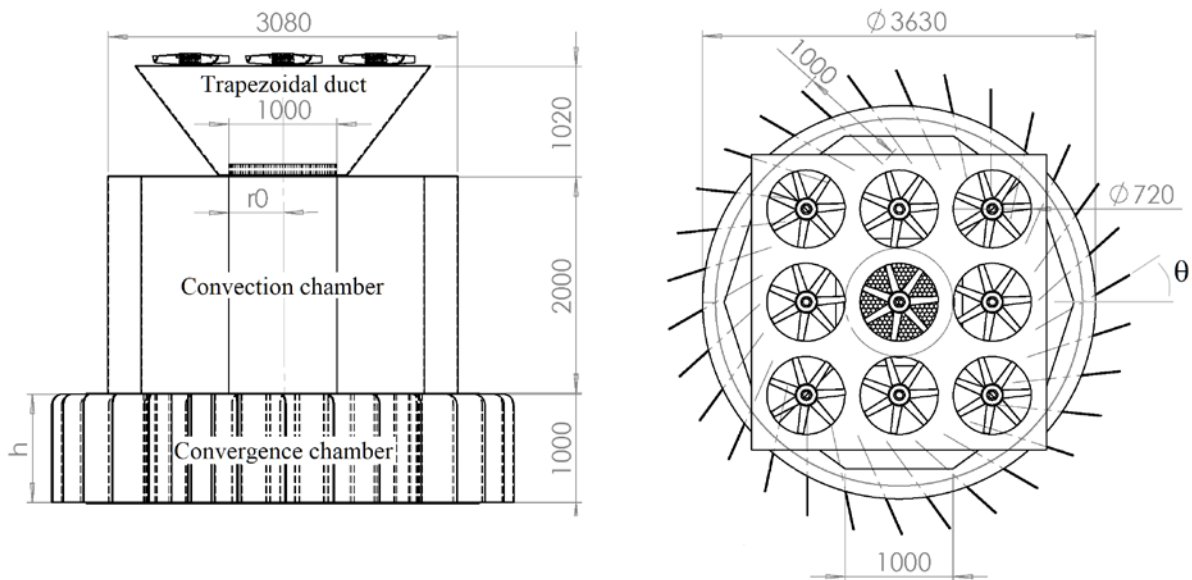
#### 107 2.1.1 *University of Birmingham Tornado-like Vortex Generator*

108 The University of Birmingham Tornado-like Vortex Generator (UoB-TVG) is a Ward-type  
109 generator used to investigate tornado-like vortices (Figure 2). A description of the generator  
110 can be found in Gillmeier et al. (2018) but for the sake of completeness a brief description is  
111 given below.

112 The generator consists of three different sections:

- 113 - a convergence chamber (1m height x 3.6m diameter) located at the bottom. 30 guide
- 114 vanes equally spaced around the convergence chamber introduce angular momentum,
- 115 and the vane angle ( $\theta$ ) can be set up from 0 degree, i.e. no swirling, to 70 degrees.
- 116 - a convection chamber (2m height x 3.1m diameter) sitting on top of the convergence
- 117 chamber.
- 118 - a trapezoidal duct with 9 identical fans sitting on top on the convection chamber and
- 119 generate the uplift through the updraft hole.

120 The updraft hole has a radius ( $r_0$ ) of 0.5m and a honeycomb is placed to reduce secondary  
 121 vortices from the fans interacting with the generated vortex (Ward, 1972). The fans generate a  
 122 mean vertical velocity ( $W$ ) of around 10 m/s at the top of the convection chamber, resulting in  
 123 a radial Reynolds number ( $Re_r = Wr_0/4\pi\nu$ ) of  $5.4 \times 10^4$ , with  $\nu$  the kinematic viscosity of air.  
 124 In the current work, the vane angle was set to 50 degrees for a resulting swirl ratio of 0.3. The  
 125 swirl ratio is defined as a measure of the circulation strength relative to the updraft ( $S = \tan\theta/2a$ ,  
 126 with  $a$  the aspect ratio defined by the ratio of the inflow height  $h$  and the updraft radius  $r_0$ ).



**Figure 2** Schematic and dimensions of the University of Birmingham Tornado-like Vortex Generator (UoB-TVG).

### 127 2.1.2 Wind velocity measurements

128 The velocity flow field was measured using a TFI 4-hole Cobra probe (TFI, 2011). The Cobra  
 129 can measure velocity data with magnitude higher than 2 m/s within a cone of influence of  $45^\circ$ .  
 130 The probe was mounted on a two-axis traverse system located off-centre inside the generator  
 131 and was held approximately 0.7m away from the traverse system to minimize disturbances.  
 132 Point measurements were undertaken radially every 0.02m from the centre of the generator up  
 133 to 0.5m and then every 0.04m up to 0.7m; and vertically every 0.02cm from 0.005m to 0.065m  
 134 and then every 0.04m up to 0.425m. The point measurements were taken with a precision of  
 135 less than 1mm. The velocity data were measured and averaged over a sampling duration of 82s  
 136 and uncertainties are taken to be  $\pm 0.5$ m/s for the tangential velocity and  $\pm 0.2$ m/s for the radial  
 137 and vertical velocity (Gillmeier *et al.*, 2018).

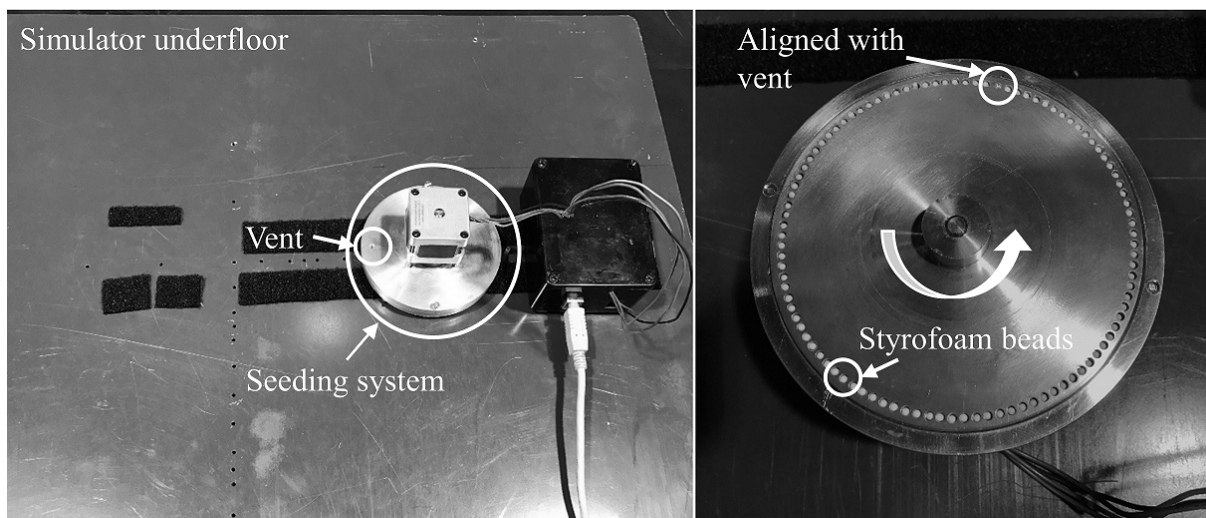
### 2.1.3 3D Particle Tracking Velocimetry

138  
139 A full description of the 3D-PTV technique and development can be found in the literature  
140 (Maas *et al.*, 1993; Malik *et al.*, 1993). In the current work, the 3D-PTV system consists of  
141 three digital high-speed cameras, Sony NEX-FS700RH, used to record the motion of the  
142 compact debris. Two cameras are positioned at the top of the convection chamber and one  
143 camera at the top of the convergence chamber. The cameras were set up to record videos at 480  
144 Hz at a resolution of  $1920 \times 1080$  pixels. Four 50W LED lights were placed inside the chambers  
145 to illuminate the volume of interest. The cameras were synchronised using an ALE718 Multi  
146 Camera LANC controller, developed by Applied Logic Engineering, Inc. The cameras  
147 calibration is performed using a two-dimensional calibration target placed inside the UoB-  
148 TVG. A three-dimensional target was 3D printed and used to estimate the spatial uncertainties.  
149 The resulting root mean square errors in the x, y and z directions are 1.05mm, 1.20mm and  
150 2.15 mm, respectively. The measurement volume in the x, y and z directions is approximately  
151 is  $2 \times 1.2 \times 0.3$  m, respectively. The digital images were pre-processed by substituting a mask  
152 from the images to filter background noises and to enhance the debris visibility. The data  
153 reduction was then performed using the open-source particle tracking software OpenPTV  
154 (OpenPTV, 2012). The output data were subsequently post-processed using MATLAB®.

### 2.1.4 Compact debris and cubic model

155  
156 Spherical Styrofoam beads were used as compact debris in the current work. A particle size  
157 and shape analysis using ImageJ software was undertaken to characterise the distributions for  
158 a large set of sieved beads (around 3000 samples). The bead diameters were found to be  
159 between 1.6mm and 2.3mm, with a mean diameter of  $1.94 \pm 0.1$ mm and a circularity around  
160 0.93.

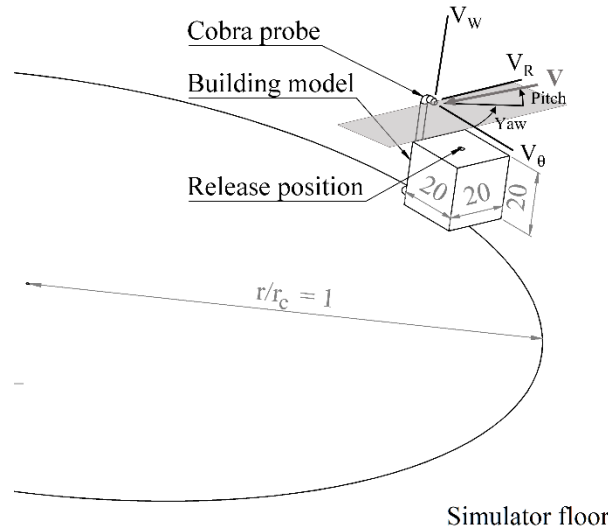
161 In order to inject the debris inside the simulator, a seeding system was designed and built  
162 (Figure 3). It consists of a disc with 100 equidistant holes of 0.002m diameter connected to a  
163 stepper motor. The disc is enclosed into a chamber and rotates at a constant speed using a  
164 motor. A vent passing through the seeding system and aligned with the holes on the wheel  
165 allowed to inject 2mm Styrofoam beads into the simulator. The pressure difference between  
166 inside and under the simulator naturally sucked the spheres through the vent. The seeding



**Figure 3** Seeding system used to inject debris into the simulator.

167 system was attached to the floor under the simulator and the vent was connected to the surface  
 168 pressure taps. This allowed, as the motor rotates, to inject each individual particle at a time.

169 A low-rise building model (20mm sides length x 20mm height, Figure 4) was 3D printed and  
 170 used to initiate debris flight from the building's roof. The building model has a geometric scale  
 171 of 1/8 with regards to the size of the experimental core radius. Wurman et al. (2013) measured  
 172 winds using a Doppler on Wheel radar during an F2 tornado and were able to determine the  
 173 size of the tornado and debris cloud around it. The tornado was found to have a core radius of  
 174 100m with a debris cloud twice as large as the tornado core itself. Therefore, the building model  
 175 would be equivalent to a building of around  $13 \times 13 \times 13$ m at full-scale. In the current work,  
 176 the building model is located at the core radius location and one hundred spheres (compact  
 177 debris) were injected into the flow field from the top of it.



**Figure 4** Building model and Cobra probe located at the core radius position.  $V_\theta$ ,  $V_R$  and  $V_W$  denote the wind tangential, radial, and vertical component, respectively, and  $\mathbf{V}$  the wind velocity magnitude.

## 178 2.2 Numerical methodology

179 The 3D motion of Styrofoam beads in a tornado-like vortex was computed numerically to  
 180 determine 3D trajectories, following a similar approach from Sassa et al. (2009). A compact  
 181 debris with only drag forces acting and no rotation is assumed. The accelerations of the compact  
 182 debris are defined as:

$$\frac{dV_{\theta,d}}{dt} = k \mathbf{V} (V_\theta - V_{\theta,d}) \quad (1)$$

$$\frac{dV_{R,d}}{dt} = k \mathbf{V} (V_R - V_{R,d}) \quad (2)$$

$$\frac{dV_{W,d}}{dt} = k \mathbf{V} (V_W - V_{W,d}) - g \quad (3)$$

183 where  $V_\theta$  and  $V_{\theta,d}$  are the local tangential wind and debris velocities respectively,  $k$  is a  
 184 buoyancy parameter (Eq. 5),  $t$  represents the time,  $V_R$  and  $V_{R,d}$  are the local radial wind and  
 185 debris velocities respectively,  $V_W$  and  $V_{W,d}$  are the local vertical wind and debris velocities  
 186 respectively, and  $\mathbf{V}$  is the vector of the relative velocity between the wind and debris defined  
 187 as:

$$V = \sqrt{(V_\theta - V_{\theta,d})^2 + (V_R - V_{R,d})^2 + (V_W - V_{W,d})^2} \quad (4)$$

188 The buoyancy parameter (Holmes, 2004) is defined as:

$$k = \frac{\rho_{air} C_D}{2 \rho_{debris} l} \quad (5)$$

189 where  $\rho_{air}$  is the air density,  $C_D$  the drag coefficient,  $\rho_{debris}$  the debris density and  $l$  the debris  
 190 characteristic length. The parameters were set to  $\rho_{air} = 1.2 \text{ kg/m}^3$ ,  $C_D = 0.5$ ,  $\rho_{debris} = 24 \text{ kg/m}^3$   
 191 and  $l = 2 d_{debris}/3$ , with  $d$  the debris diameter, giving a value of  $k = 9.5$ .

192 The velocity and position components, denoted hereafter with  $\bullet$ , were computed numerically  
 193 using a linear method as:

$$V_\bullet^{t+1} = V_\bullet^t + A_\bullet^t \Delta t \quad (6)$$

$$X_\bullet^{t+1} = X_\bullet^t + V_\bullet^t \Delta t \quad (7)$$

194 with  $\Delta t$  a timestep of 0.001s. A timestep of 0.001s offers a good ratio of stability/time efficiency  
 195 considering the maximal wind velocities of the flow field used in the current work.

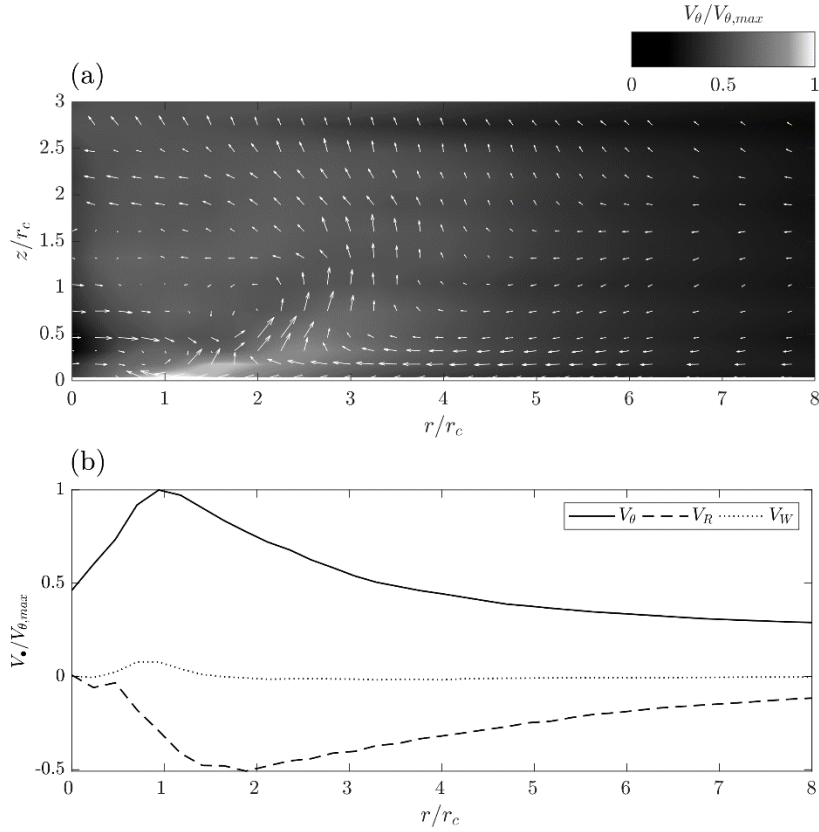
196 Experimental data were incorporated into the calculations, by incorporating the experimental  
 197 wind and tracked data into the calculated accelerations (Eq. 1-4). For  $t=0$ , the initial conditions  
 198 ( $V_\bullet^0$  and  $X_\bullet^0$ ) were set to the velocity and position data obtained from the particle tracking  
 199 experiment (Eq. 6-7). To determine at each timestep the wind velocities to be incorporated into  
 200 the calculations, bilinear interpolation was used to interpolate sub-grid velocities from the  
 201 positions  $X_\bullet^t$ . Since the measurement grid was coarse, this technique improves the accuracy and  
 202 robustness of the calculations by interpolating wind velocities rather than finding the closest  
 203 measured point. The modelling of the debris trajectory was stopped when the debris first  
 204 impacts the ground.

## 205 3 RESULTS

### 206 3.1 Tornado wind field

207 Figure 5 shows the 3D velocity flow field and radial profile of the velocity components at the  
 208 lowest height measured ( $z/r_c = 0.03$ ) in the UoB-TVG for a swirl ratio of 0.3. Heights and radial  
 209 positions are normalised by the core radius  $r_c$ , located at  $r/r_c = 1$ , and velocities are normalised  
 210 by the maximal tangential velocity  $V_{\theta,max}$ . The core radius was estimated as the radius where  
 211 the maximal tangential velocity occurs and was found to be equal to 0.144m. Figure 5 (a) shows  
 212 the complexity of a tornado-like flow field, with a strong inflow in the lower level towards the  
 213 vortex core region and a strong updraft at the corner region ( $r/r_c = 1$ ). The strong tangential  
 214 velocity region located at the core radius region, the core radius increase with height, as well  
 215 as the positive radial inflow inside the vortex core are consistent with previous work undertaken  
 216 by Haan et al. (2008), Gillmeier et al. (2018) and Tang et al. (2018). The recirculation located  
 217 at  $z/r_c = 1$  and  $r/r_c = 2$  seems to indicate the presence of a vortex breakdown, which would  
 218 indicate that for  $S=0.3$  the tornado flow field is not yet fully developed into a two-cell vortex.  
 219 Vortex wandering characterisation (not shown here) revealed that for the current swirl ratio,  
 220

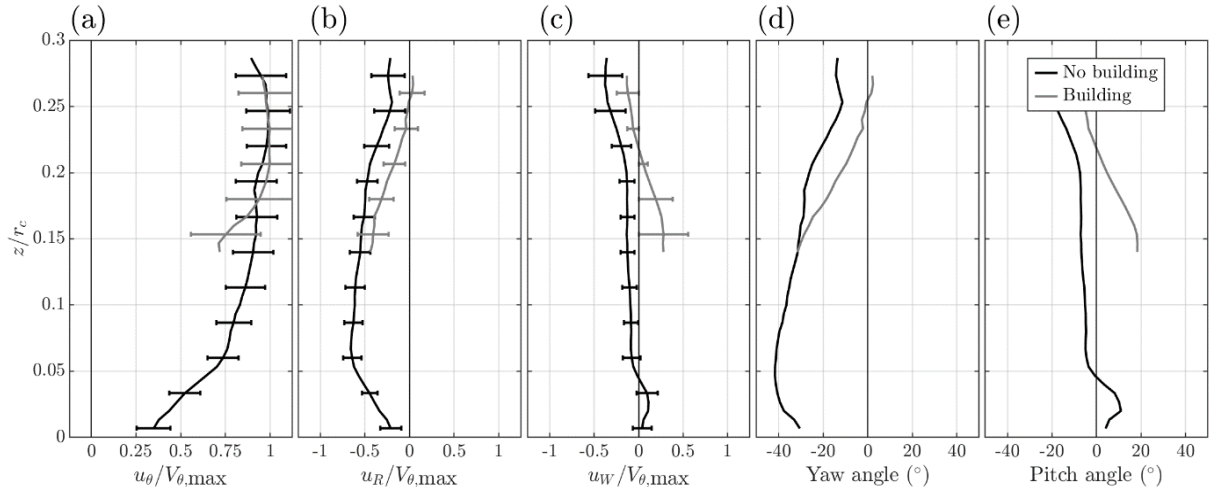




**Figure 5** (a) 3D velocity field for  $S=0.3$ . Contour denotes the tangential velocity and the vectors denote the velocity vector between radial and vertical velocity. (b) Velocity components for the lowest height (5mm), located in the boundary-layer.

221 the vortex typically wanders around the centre of the simulator up to  $0.15r_c$ . Overall, the vortex  
 222 structure shows a relatively good agreement with previous work from Gillmeier et al. (2018)  
 223 and the VorTECH generator for a similar swirl ratio (Tang *et al.*, 2018). Figure 5 (b) shows the  
 224 tangential, radial and vertical velocity profiles for the lowest height measured and accounts for  
 225 a tangential velocity increasing until the core radius  $r_c$ , a strong radial inflow up (negative radial  
 226 velocities) to half the maximum tangential velocity ( $V_R/V_{\theta,max} = 0.44$ ) at  $r/r_c = 2$  and a weak  
 227 updraft around the core radius location.

228 Figure 6 shows the vertical profiles of velocities closed to the ground at the core radius location  
 229 with and without building model included. To measure the vertical profile on top of the model,  
 230 the Cobra probe's head was mounted under the simulator at the core radius location. The  
 231 probe's head was placed perpendicular to core centre's direction and at the building model's  
 232 corner between the top and leeward side, as shown in Figure 4. With no building model, the  
 233 tangential velocity (Figure 6 (a)) at this location is similar in many respects to a typical  
 234 boundary-layer profile. The radial velocity (Figure 6 (b)) is predominantly negative over the  
 235 height and accounts for the radial inflow as found in Figure 5. The vertical velocity (Figure 6  
 236 (c)) shows a change of sign at around  $z/r_c = 0.05$  which is difficult to interpret since the velocity  
 237 remains relatively close to zero and the measurement is taken in a turbulent region. When the  
 238 building model was included, the profile was measured from the top of the leeward side, to be  
 239 as close as possible from the debris release position. It shows that the tangential velocity is  
 240 typically not affected with height by the building model (Figure 6 (a)), and that the radial and  
 241 vertical velocities are reduced in magnitude but still display a similar profile with height

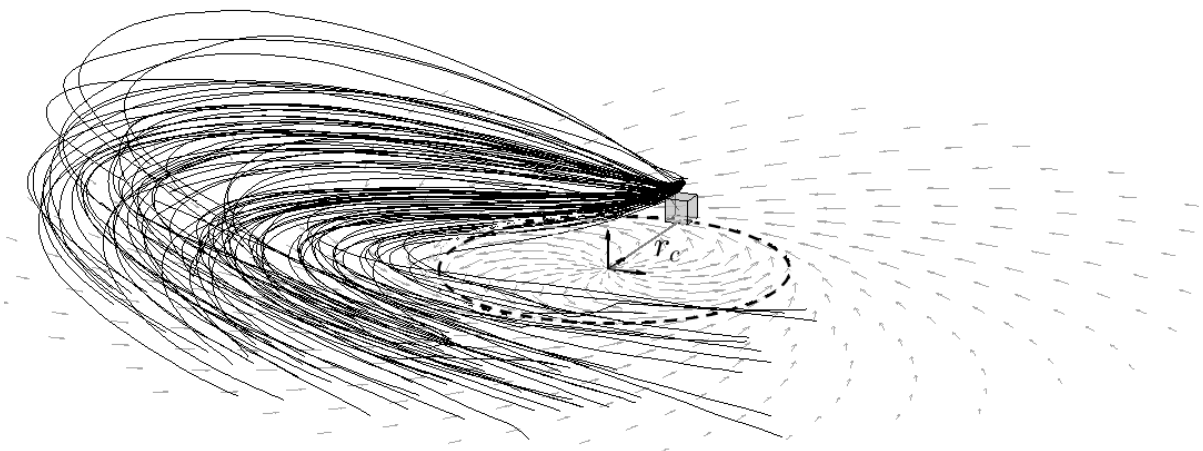


**Figure 6** Vertical velocity profiles of normalised – (a) tangential, (b) radial, (c) vertical velocity, respectively, (d) yaw angle and (e) pitch angle. The black line denotes when the building model is not included into the flow field, and the grey line when it is included. The horizontal capped lines denote the measurement uncertainties. The yaw angle denotes the horizontal angle between the tangential and radial velocities and the pitch angle the vertical angle with regard to the horizontal plane.

242 (Figure 6 (b and c)). The largest discrepancies observed around  $z/r_c = 0.15$  are associated with  
 243 a potential flow separation occurring on top the building model. Numerical simulations of a  
 244 low-rise building in tornado-like flow field with a similar swirl ratio ( $S=0.4$ ) showed that a  
 245 recirculation is generated on the lee-ward side of the building (Nasir, 2017). That recirculation  
 246 produces a downward flow in the wake of the building which most likely explain the debris  
 247 behaviour described in the following section. The yaw angle (Figure 6 (d)) refers to the  
 248 horizontal wind direction facing the probe while the pitch angle (Figure 6 (e)) refers to the  
 249 vertical wind direction.

### 250 3.2 Example of 3D trajectories from a low-rise building

251 Figure 7 illustrates an example of 3D debris trajectories initiated from the top of a low-rise  
 252 building located at the core radius location and obtained using 3D-PTV. It shows that the debris  
 253 are initiated with different horizontal and vertical directions resulting in a large spread of the  
 254 falling locations. Further interpretation and analysis are discussed in the following section.

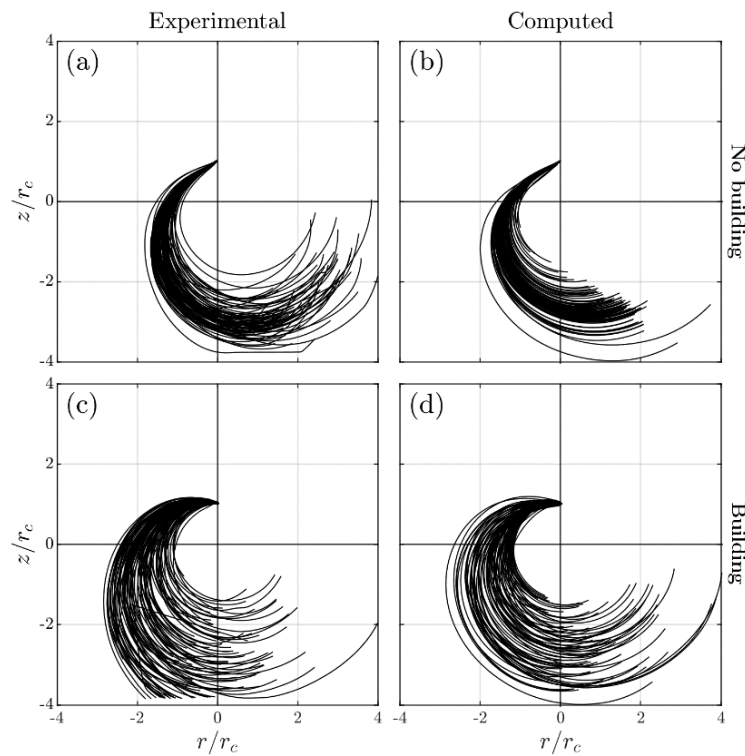


**Figure 7** 3D representation of the flight of debris from a low-rise building at the core radius ( $r_c$ ) location. The dash line denotes the size the core radius, the arrows denotes the idealised flow-field at building height.

255 3.3 Comparison of debris flight trajectories with no building and a building included

256 The comparison between different debris flight configurations is shown in Figure 8. The debris  
 257 trajectories obtained experimentally are shown in Figure 8 (a and c), while Figure 8 (b and d)  
 258 show the numerical reproduction of these trajectories. In both cases, visual similarities are  
 259 observed and are quantified in Figure 10. Figure 8 (a and c) illustrate the typical set of  
 260 experimental data obtained using 3D-PTV when injecting 100 particles from the core radius  
 261 position. Similar trajectory behaviours were found in the literature, although studied  
 262 numerically, when debris were released from the core radius location (Maruyama, 2011; Noda  
 263 *et al.*, 2013). In the present cases, the debris were injected from the floor simulating the  
 264 initiation of debris flight in a landscape environment (i.e. no building), and from the top of a  
 265 building. The tracking was stopped when the particles impacted the ground after being airborne  
 266 to reflect what could be observed during a full-scale event. However, it is worth noting that  
 267 due to the elasticity of the Styrofoam bead, its light mass and the smooth floor surface in the  
 268 tornado generator, the beads bounce when falling to the ground. For the sake of the current  
 269 analysis, such motion has been neglected and is not considered further.

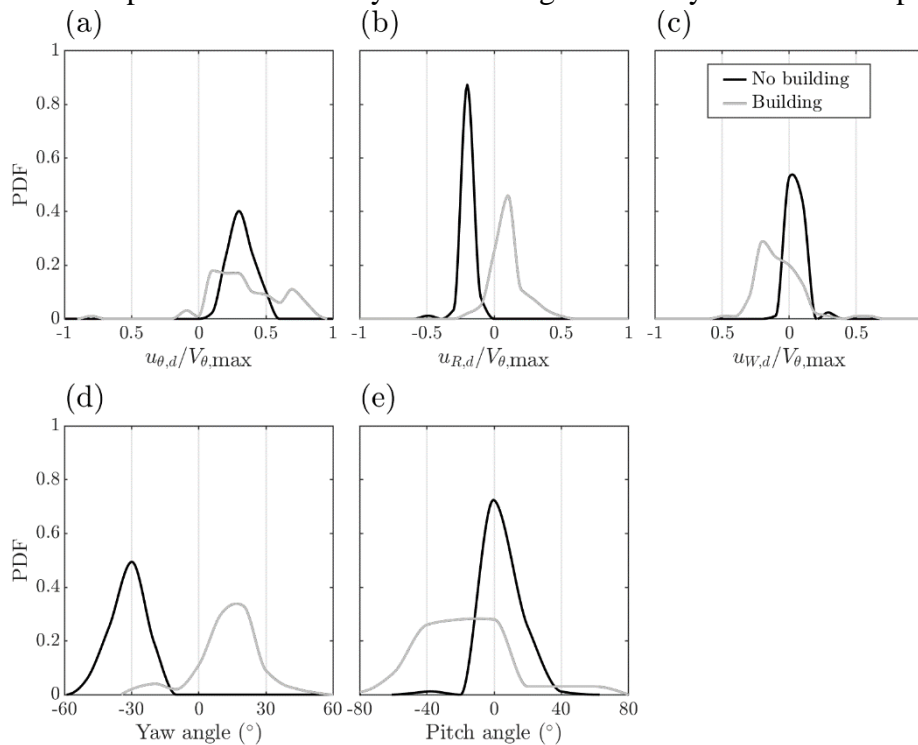
270 Figure 8 (a) shows that the particles are ejected outward from the vortex following a relatively  
 271 straight line for a short period of time. The particles are then swirling back again around the  
 272 vortex before falling to the ground. It also illustrates that a slight variation of the early stage  
 273 trajectories can result in a large variation of the impact location. The variability in the early  
 274 stage trajectories observed in the trajectories could be due to the varying size of the beads,  
 275 although it is most likely that the variation is due to the vortex wandering and/or to turbulent  
 276 fluctuations of the local wind field. Figure 8 (c) shows the debris trajectories when injected



**Figure 8** Top view of experimental and numerical debris trajectories. Left column denotes the experimental trajectories (a-c), right column the computed trajectories (b-d). The top row denotes the trajectories when the building model is not included (a-b), the bottom row the trajectories when the building model is included (c-d).

277 from the top of the cubic model. It illustrates a greater variation of the early stage trajectories  
 278 than without the model resulting in a larger region where the debris falls. The larger variation  
 279 can be attributed to the disturbance of the local flow field around the building model. This  
 280 assumption can be supported by the initial velocity distributions, i.e. the first debris velocities  
 281 calculated using 3D-PTV (Figure 9). It appears that debris are driven by positive tangential  
 282 velocities in both situations (Figure 9 (a)), however radial and vertical velocities have opposite  
 283 sign between the two. As a result, an overall difference of 40 degrees in direction can be noticed  
 284 when including the building into the flow field (Figure 9 (d)). Furthermore, Figure 9 (c) shows  
 285 that debris are mainly driven by both positive and negative vertical velocities when released  
 286 from the top of the building, which could be explained by a potential flow separation. When  
 287 debris are driven by negative vertical velocity, the debris falls quickly to the ground and is  
 288 travelling longer distances due to strong local tangential velocities, and therefore accounts for  
 289 the closest trajectories to the vortex core. On the other hand, the debris tends to be ejected  
 290 outward the core and to fall at even longer distances away from the core radius. This behaviour  
 291 is visually observed in Figure 7. Therefore, the risks with windborne debris initiated from a  
 292 low-rise building are not only associated with long-range falling debris but also with possible  
 293 short-range impacts at higher velocities.

294 Figure 8 (b and d) presents the computed debris trajectories using experimental wind and  
 295 tracking data. In both cases, the computed trajectories are in good agreement with the tracked  
 296 trajectories. Although the flow field is assumed axisymmetric and does not include any  
 297 turbulence, it shows that the present methodology can reproduce the trajectories variability.  
 298 Table 1 compares the components of wind and initial debris velocities and shows that when  
 299 the building is not included, the velocity components are matching wind and debris. This  
 300 validates the assumption that variability in debris flight is mainly a wind-driven process. When



**Figure 9** Distribution of debris initial tangential (a), radial (b), vertical (c) velocities, yaw (d) and pitch (e) angle, when the low-rise building model is and is not included in the tornado-like flow field.

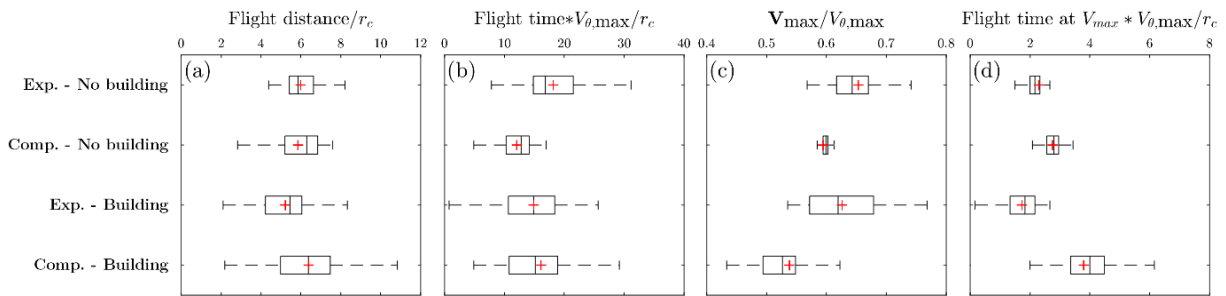
301 the building is included, large discrepancies are observed between wind and debris which could  
 302 be explained by the fact that velocities are not measured exactly at the same position.

**Table 1** Comparison between wind and initial debris velocities, pitch and yaw angles, when the building model was included and not included in the simulations.

		$u_{\theta} / V_{\theta, \max}$	$u_R / V_{\theta, \max}$	$u_W / V_{\theta, \max}$	Pitch ( $^{\circ}$ )	Yaw ( $^{\circ}$ )
No building	Wind	0.30	-0.20	0.04	25.8	4.10
	Debris	0.27	-0.16	0.04	30.7	7.10
Building	Wind	0.70	-0.44	0.27	-32.1	17.7
	Debris	0.38	0.11	-0.09	-16.7	-17.3

303

304 Figure 10 shows the debris flight analysis with normalised flight parameters determined from  
 305 the debris trajectories (Figure 8). The flight distance (Figure 10 (a)) accounts for the cumulative  
 306 distance travelled by a debris, the flight time (Figure 10 (b)) the duration while the debris is  
 307 airborne,  $V_{\max}$  (Figure 10 (c)) the maximal velocity reached by the debris, and flight time to  
 308  $V_{\max}$  (Figure 10 (d)) the duration taken by a debris to reach the maximal velocity. Overall, it  
 309 shows that including a building model does not affect the flight behaviour of the debris  
 310 significantly. This corroborates the assumption that the overall flow field is not affected by the  
 311 building model, as shown in Figure 6. The numerical simulations tend to reproduce some  
 312 behaviour, although more variability is generally observed. The flight distance shows good  
 313 agreement with experimental results and coincides with the debris trajectories overall  
 314 behaviour (Figure 8), as does the flight time. However, the numerical simulations show  
 315 limitations in modelling the maximum velocity  $V_{\max}$ , and therefore the flight time to reach  $V_{\max}$   
 316 (Figure 10 (c and d)). This could be explained by the flow field variability (due to wandering  
 317 and/or turbulence) that is diminished after averaging and could result in a stronger  
 318 instantaneous flow field at the time of the experiment.



**Figure 10** Debris flight analysis – the red cross denotes the mean value, the box interquartile range, the vertical line the median and the whiskers the range from minimal to maximal values.

319

#### 4 CONCLUDING REMARKS

In this work, physical and numerical simulations were used to investigate the flight of debris in tornado-like vortices initiated with and without a low-rise building model. A medium-range vortex was simulated in a tornado-like vortex generator with a swirl ratio of 0.3. Wind measurements displayed a complex flow field with a strong tangential velocity in the corner region of the vortex, and strong radial inflow outside the vortex core. Vertical profiles of winds show that overall, the time-averaged flow field is not significantly affected by included a low-rise building model but tends to reduce in magnitude of the radial and vertical velocities with height. The current study focused on releasing Styrofoam beads in a tornado-like flow field and investigates the effect of releasing the debris from the top of the building model. Debris trajectories with and without the building model included show differences in early stage trajectory behaviors. When the building model is not included, trajectories follow a narrower path in early-stage flight with respect to the mean local flow field direction. When the building model is included, the early-stage debris trajectories have a different main direction with more variability resulting in a wider falling region. However, the analysis of different characteristic flight parameters shows good agreement between the two situations. It confirms that in the present case the overall flow field in a tornado-like vortex is not affected by including a building model; however, the local flow field in the proximity of the model is affected. Numerical calculations were used to simulate the debris trajectories and flight behaviour. Using debris flight equations with wind and tracked experimental data, the computed trajectories show strong similarities with experimental trajectories. The debris flight analysis shows that the simulations could reproduce the overall trajectories but had limitations in simulating maximum velocities. The limitations could be attributed to the experimental wind data used to compute the trajectories and not necessarily to the methodology itself. Finally, this numerical methodology presents advantages to investigate debris flight in tornadoes:

- Gillmeier et al. (2018) showed that analytical vortex models can reproduce some of the behaviours observed in full scale and modelled tornadoes but fail to reproduce the complexity of the 3D flow field. Wind measurements help describing the complexity that models fail to reproduce and therefore integrate a more “realistic” flow field to the study of debris flight.
- Using bilinear interpolation to better approximate the local wind field increases the accuracy of the computed positions and velocities and does not require a fine measurement grid. It also significantly reduces the computation time and does not require high computational power.
- The methodology also allows any type of debris in tornado-like flow field to be studied. By adjusting the debris flight equations to either compact (Holmes, 2004) or plate (Holmes *et al.*, 2006), a large range of debris type, size and density (i.e. Tachikawa number) could be investigated.
- The methodology could also be applied to full-scale tornado wind data, however, Doppler radars often fail to measure wind below 30m AGL (Wurman *et al.*, 2013), which is where most of debris are becoming airborne and impacting structures. When data becomes available, comparison and validation could be undertaken at full-scale.

362 Although the methodology has many advantages, it also has limitations. The assumption of an  
363 axisymmetric flow field does not reflect the complexity of a tornado-like flow field. Volumetric  
364 wind measurements using neutral buoyancy tracers and laser illumination (Tomographic PIV,  
365 3D-PTV, etc.) could help retrieving instantaneous 3D flow field. Finally, the methodology only  
366 considers averaged flow fields. Turbulence could be incorporated into the simulation (Holmes,  
367 2004) to recreate wind field fluctuations but would require a better understanding of the  
368 turbulence in tornadoes/tornado-like vortices.

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