

# On the interstitial induced lattice inhomogeneities in nitrogen-expanded austenite

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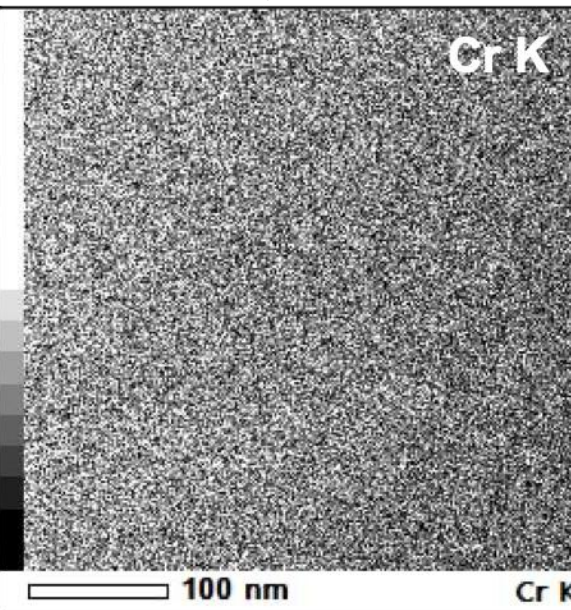
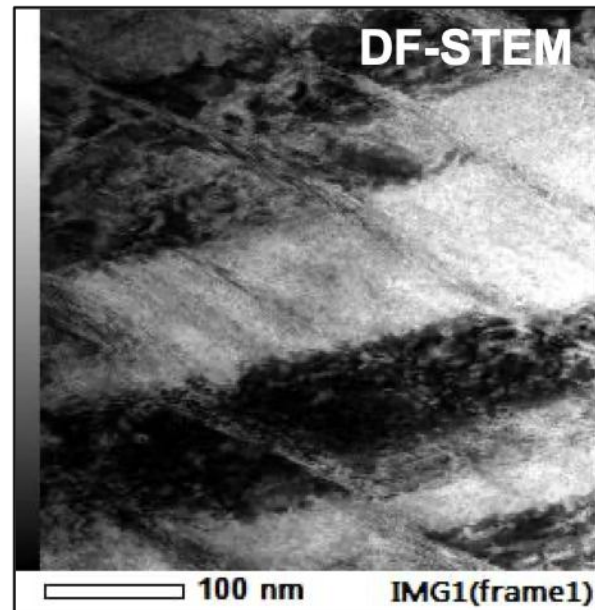
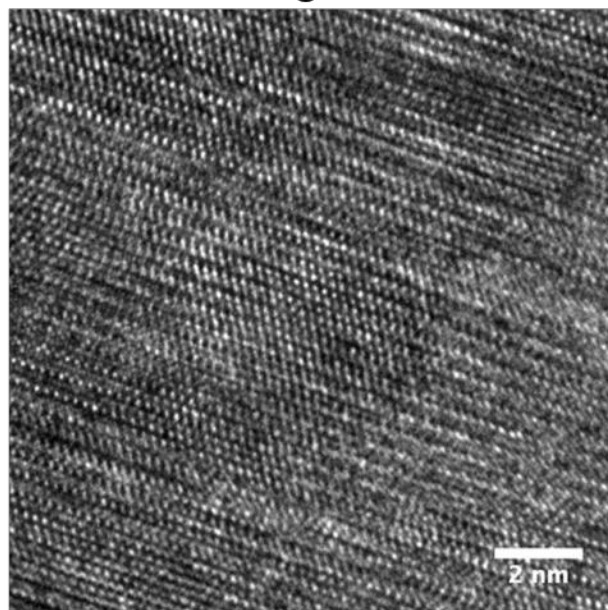
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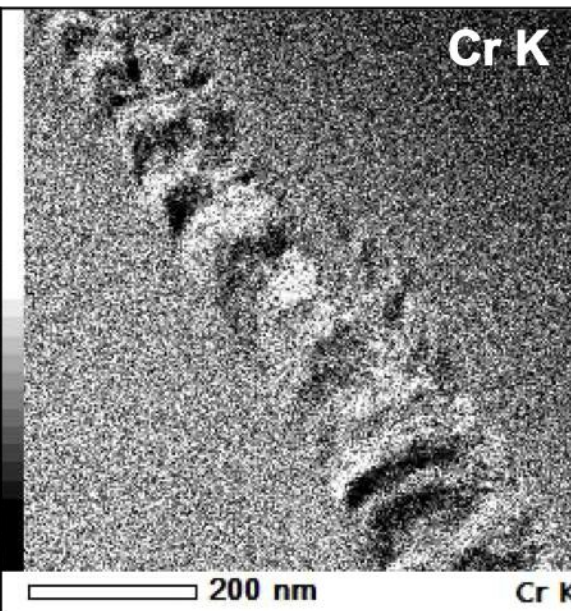
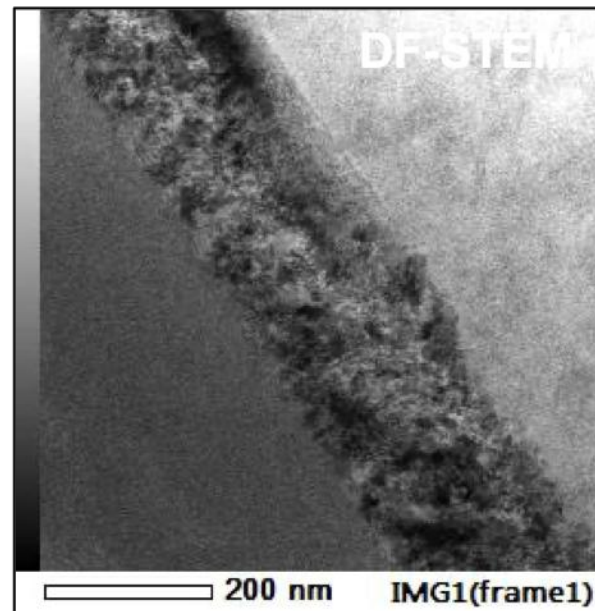
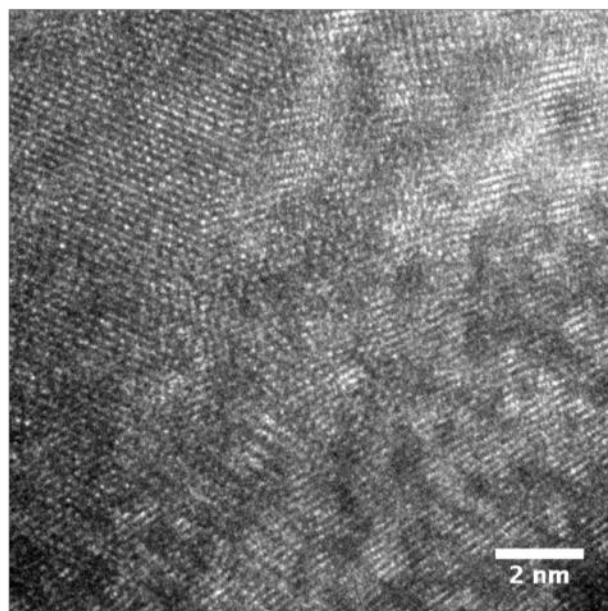
HRTEM of the  
inhomogeneities

STEM-EDX for  
inhomogeneities in parent  $\gamma_N$

HCP- $\epsilon_N$



Shear  
band



1 On the nitrogen interstitial induced lattice inhomogeneities in nitrogen-expanded austenite

2

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5

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10

11 Abstract

12 Lattice inhomogeneities, i.e. nitrogen interstitial-induced hexagonal-close-packed martensite  
13 (HCP- $\epsilon_N$ ) and shear bands, can form in face-centred cubic nitrogen-expanded austenite (FCC-  
14  $\gamma_N$ ) synthesised on Fe-Cr-Mn and Fe-Cr-Ni austenitic stainless steel (ASS) using triode-plasma  
15 nitriding (TPN). Homogenous elemental distribution between HCP- $\epsilon_N$  and FCC- $\gamma_N$  supports  
16 the displacive shear transformation mechanism proposed for a high-Mn, low stacking fault  
17 energy ASS under nitrogen interstitial-induced deformation. While being a product of  
18 transformation-induced plasticity effect, HCP- $\epsilon_N$  exhibits similar lattice expansion behaviour  
19 to the parent FCC- $\gamma_N$ . However, inhomogeneous elemental distributions in the shear bands  
20 formed in  $\gamma_N$  layers on a 400°C TPN-treated high-Ni ASS indicate local migration of  
21 substitutional elements.

22

23

24 Keyword: nitriding; expanded austenite; shear banding; martensitic phase transformation;  
25 high-resolution electron microscopy (HREM)

26 Low temperature thermochemical diffusion treatments have been developed as an important  
27 and effective surface engineering method to improve the wear resistance of austenitic stainless  
28 steels (ASSs), generating hard, yet corrosion resistant surface layers supersaturated with  
29 interstitial nitrogen and/or carbon, named ‘expanded austenite’ or ‘S-phase’ [1-5]. A  
30 metastable, expanded face-centred cubic (FCC) phase appears to be the most plausible and  
31 widely applicable interpretation so far for this type of interstitially-supersaturated material  
32 formed on ASSs under low temperature diffusion treatment, with a widely used designation  
33 ‘ $\gamma_N$ ’ and ‘ $\gamma_C$ ’ (when supersaturated with nitrogen and carbon, respectively) [6-10]. After  
34 prolonged treatment at ‘low’ temperatures (these being typically below  $\sim 450^\circ\text{C}$  for nitrogen  
35 and  $\sim 550^\circ\text{C}$  for carbon), an excessive concentration of interstitial atoms is diffused into the  
36 substrate near-surface (e.g. up to  $\sim 38$  at.% N in AISI 316 ASS [8]) which, although many times  
37 the expected equilibrium solubility limit in ASSs, does not lead to direct phase transformation  
38 of the original FCC structure.

39

40 Following the work of Ichii et al. in 1986 [11], the anomalous XRD peak shifting and  
41 anisotropic lattice expansion of  $\gamma_N$  has been known and investigated for many years. The lattice  
42 expansion anomalies observed have been attributed to stacking fault (SF) generation and to the  
43 hkl-dependent (anisotropic) ‘deformation’ behaviour of FCC material in elastic and/or plastic  
44 deformation regimes (i.e. accommodation of supersaturated concentrations of nitrogen and/or  
45 carbon elasto-plastically) [6, 7, 10]. A more recent study [12] reveals anisotropic lattice rotation  
46 for AISI 316 ASS after low temperature nitriding using electron backscatter diffraction (EBSD),  
47 where a reasonably good correlation was found between nitrogen composition-induced  
48 anisotropic lattice rotation and out-of-plane tensile deformations, that could be simulated using  
49 the Taylor-Bishop-Hill model [13].

50

51 Under transmission electron microscopy (TEM),  $\gamma_N$  has often been reported to contain a large  
52 number of lattice defects, e.g. dislocations, nanotwins and stacking faults (SFs) [14-19]. These  
53 crystallographic defects can interrupt the local crystallinity of  $\gamma_N$ , but do not necessarily alter  
54 the underlying crystal structure. Recently it was proposed that, similar to FCC materials under  
55 mechanical deformation, the observed interstitial supersaturation-induced deformation  
56 mechanisms (and their prevalence) could depend on the materials stacking fault energy (SFE)  
57 [20]. Hexagonal-close-packed (HCP)  $\epsilon_N$  [19, 20] and shear bands [20] could form in  $\gamma_N$  layers  
58 when a large amount of such defects accumulate locally, causing the underlying

59 crystallographic structure to be severely interrupted (and/or altered) over comparatively large  
60 lengthscales. The “mono-phased” metastable nitrogen diffusion layer could contain  
61 inhomogeneous regions (formed predominately via chemical composition-mediated  
62 deformation phenomena, occurring during/after low temperature thermochemical diffusion  
63 treatments) in place of, or prior to, the known nitride phase precipitation, that occurs at more  
64 elevated treatment temperatures and/or prolonged treatment times. A good understanding of  
65 these ‘deformation’-mediated inhomogeneities (although not likely to be found in ‘Cr-free’  $\gamma_N$   
66 [21] at comparatively low interstitial supersaturation levels) would provide valuable insights  
67 into the structure (and process-structure-property relationships) of metastable, interstitially  
68 highly-supersaturated, anisotropically-expanded austenite(s).

69

70 The nominal compositions of Staballoy AG17® and RA 330® are Fe-17Cr-20Mn-0.5N and  
71 Fe-19Cr-35Ni, respectively. In this work, both alloys were treated under triode-plasma  
72 nitriding (TPN) at 400°C for 20hrs (-200V substrate bias, 0.4 Pa chamber pressure and 7:3  
73 N<sub>2</sub>:Ar gas volumetric ratio); detailed property characterisation and nanostructural studies are  
74 reported elsewhere [20]. For the results presented here, thin foils were prepared (in cross-  
75 section) from the treated surfaces by focused ion beam (FIB) milling, using an FEI Quanta 200  
76 3D instrument (with Gallium ion beam). In this investigation, high-resolution TEM (HRTEM)  
77 images were obtained from the prepared foils using a cold field emission gun (C-FEG) JEOL  
78 R005 double aberration corrected TEM operating at 300 kV. Scanning TEM (STEM) imaging  
79 and energy-dispersive X-ray (EDX) elemental mapping at high point-to-point resolution was  
80 carried out using a C-FEG JEOL F200 microscope at 200 kV (with twin, solid state, ultra-  
81 sensitive silicon drift X-ray detectors). EDX data was analysed using Analysis Station 4  
82 software provided by JEOL. Glancing angle X-ray diffraction (GAXRD) was carried out at 2°  
83 incident angle using a PANalytical Xpert3 diffractometer (with monochromated CuK $\alpha_1$   
84 0.15406 nm), at 45 kV and 40 mA.

85

86 In the  $\gamma_N$ -AG17 layer synthesised on alloy AG17 after TPN at 400°C and 20hrs, the stacking  
87 sequences change from ABCA... in FCC- $\gamma_N$  (**Fig. 1b**) to ABAB... in HCP- $\epsilon_N$  (**Fig. 1c**). The  
88 FCC- $\gamma_N$  and HCP- $\epsilon_N$  interfaces follow a Shoji-Nishiyama orientation relationship, i.e.  
89  $\gamma_N(111)//\epsilon_N(0002)$ . The inverse FFT images (**Fig. 1d, e**) signify that the close-packed planes of  
90 neither  $\gamma_N$  nor  $\epsilon_N$  appear as continuous straight lines – that may be attributed to the defects  
91 generated during the insertion of interstitial nitrogen. The interplanar spacings of  $\gamma_N(111)$  and

92  $\epsilon_N(002)$  are approximately the same under HRTEM (**Fig. 1b, c**) at  $\sim 0.219$  nm. One could  
93 estimate that  $a_{\gamma_N} = d_{\gamma_N(hkl)} \times \sqrt{h^2 + k^2 + l^2} = 0.379$  nm,  $a_{\epsilon_N} = \frac{\sqrt{2}}{2} a_{\gamma_N} = 0.268$  nm  
94 and  $c_{\epsilon_N} = 2 \times d_{\gamma_N(111)} = 0.438$  nm ( $a_{\gamma_N}$  is the lattice parameter for  $\gamma_N$ ;  $a_{\epsilon_N}$  and  $c_{\epsilon_N}$  are the  
95 lattice parameters for  $\epsilon_N$ , where  $c_{\epsilon_N}/a_{\epsilon_N} \sim 1.63$ ).

96

97 STEM-EDX analysis suggests a homogenous elemental distribution in the TPN 400°C/20hrs  
98 treated AG17 sample (**Fig. 2a-f**), that is self-consistent to the previously proposed displacive  
99 martensitic shear transformation mechanism for the formation of HCP- $\epsilon_N$  in expanded austenite  
100 [20]. Evidenced by the homogenous N distribution seen in **Fig. 2f**,  $\epsilon_N$  regions have a nitrogen  
101 content equivalent to the  $\gamma_N$  matrix at the same depth in the diffusion layer, and may also  
102 possess a decreasing nitrogen concentration-depth profile in the treatment layer, similar to  $\gamma_N$ .  
103 The formation of HCP- $\epsilon_N$  would start from (and require a minimum composition-induced  
104 lattice deformation at) some nitrogen concentration threshold. Thus,  $\epsilon_N$  should have a  
105 composition window that overlaps with that of  $\gamma_N$  at higher nitrogen concentration levels.

106

107 More importantly, both  $\gamma_N$  and  $\epsilon_N$  peaks are evident under GAXRD (**Fig. 3**). Taking the  $\gamma_N(111)$   
108 peaks at  $\sim 40.93^\circ$  and  $\sim 39.52^\circ$  in **Fig. 3**, one could calculate that  $a_{\gamma_N} = \sim 0.382$  nm and  $a_{\gamma_N} =$   
109  $\sim 0.395$  nm after 4hrs and 20hrs of TPN, respectively. Taking the S-N orientation relationship  
110 and  $d_{\gamma_N(111)} = d_{\epsilon_N(002)}$ , the lattice parameters for  $\epsilon_N$  can be estimated as  $a_{\epsilon_N} = \sim 0.270$  nm  
111 and  $c_{\epsilon_N} = \sim 0.441$  nm after 4hrs of TPN, and  $a_{\epsilon_N} = \sim 0.279$  nm and  $c_{\epsilon_N} = \sim 0.456$  nm after  
112 20hrs of TPN. The estimated peak positions of  $\epsilon_N$ , strikingly, match well with those observed  
113 (hence, indexed) in **Fig. 3**. Noticeably, the X-ray attenuation depth (estimated using AbsorbDX  
114 software) on untreated alloy AG17 is  $\sim 0.4$   $\mu\text{m}$  under GAXRD at  $2^\circ$  glancing angle and  $\sim 2$ -3  
115  $\mu\text{m}$  under  $\theta$ -2 $\theta$  geometry. The observation of  $\epsilon_N$  peaks under GAXRD in this study (but not  
116 under  $\theta$ -2 $\theta$  XRD, as in Ref. [20]) is believed to be due to the large volume fraction of  $\epsilon_N$  in the  
117 uppermost regions of the diffusion layer.

118

119 In good agreement with the increasing volume fraction of  $\epsilon_N$  with treatment time at 400°C [20],  
120 the relative intensity of the  $\epsilon_N(101)$  reflection increases from 4hrs to 20hrs (**Fig. 3**). Most  
121 importantly, comparing with  $\gamma_N$ , HCP- $\epsilon_N$  also shows increasing peak shifts to lower  $2\theta$  angles  
122 with treatment time, indicating that  $\epsilon_N$  undergoes similar lattice expansion under interstitial  
123 nitrogen absorption. Based on this peak-shift behaviour – and the dilated crystal structure that  
124 causes it – HCP- $\epsilon_N$  phase may tentatively be named as “expanded martensite”. However,

125 whereas  $\gamma_N$  forms from unexpanded  $\gamma$  via N absorption,  $\epsilon_N$  likely forms directly from  $\gamma_N$  (and  
126 not from unexpanded  $\epsilon$ ) through a composition-induced martensitic shear transformation above  
127 some critical N-concentration. Based on its formation mechanism,  $\epsilon_N$  could therefore be  
128 categorised as a special deformation-induced martensite (DIM) with expanded lattice  
129 parameters formed under composition-induced deformation, in contrast to ‘conventional’  
130 DIMs formed under mechanical deformation. Since the expanded lattice of  $\gamma_N$  would “shrink”  
131 under low-temperature annealing,  $\epsilon_N$  very likely exhibits similar reversibility in  
132 expansion/contraction. The nitrogen composition-induced martensitic transformation from  $\gamma_N$   
133 to  $\epsilon_N$  could also be reversible upon the loss of interstitial nitrogen. In either case, the level of  
134 compressive residual stress (established under supersaturation and loss of interstitial nitrogen)  
135 would play a determining role. These arguments imply the need for further studies on the  
136 reversibility of plasticity mechanisms in  $\gamma_N$ ; e.g. via low-temperature annealing of TPN-treated  
137 layers (carried out in combination with residual stress and chemical composition analysis).

138

139 Noticeably, owing to stress relaxation during TEM sample thinning, the lattice parameters  
140 estimated based on HRTEM are systematically smaller than those estimated based on XRD  
141 measurement. The lattice parameters of  $\epsilon_N$ , if estimated using the  $\gamma_N(200)$  peak positions  
142 instead of  $\gamma_N(111)$ , give  $\epsilon_N$  peak positions at  $2\theta$  angles slightly lower than those observed from  
143 GAXRD. Although the true lattice parameters of  $\gamma_N$  are difficult to determine with precision, a  
144 plausible HCP- $\epsilon_N$  structure is proposed based on the  $\gamma_N(111)$  XRD peak position, that gives  
145 satisfactory explanation to GAXRD profiles. However, the true structure of  $\epsilon_N$  may deviate  
146 from the estimated one, considering the defects in both  $\epsilon_N$  and  $\gamma_N$  implied by **Fig. 1(d, e)**, and  
147 the high residual stress developed in the diffusion layer.

148

149 Owing to twinning-induced plasticity (TWIP) and/or transformation-induced plasticity (TRIP)  
150 effects, low-SFE alloys are known for their good uniform plastic ductility with continuous  
151 working hardening and delayed necking upon deformation, while  $\gamma_N$ -316 layers synthesised at  
152 high nitriding potential have been seen to show a tendency to crack [22]. With the ability to  
153 undergo martensitic shear transformation and local strain hardening, the composition induced  
154 plastic deformation could be distributed more uniformly within the surface treatment layer on  
155 alloy AG17, i.e. that TPN-treated low-SFE substrates may accommodate a higher amount of  
156 (composition induced) strain than high-SFE ones, without “failure”. Compared to the localised  
157 shear bands and micro-cracks seen in  $\gamma_N$ -330 [20], the large amount of  $\epsilon_N$  phase in  $\gamma_N$ -AG17

158 appears to “share” the composition-induced deformation. Additionally, substantial HCP- $\epsilon_N$  is  
159 found in  $\gamma_N$ -AG17 [20] (and most likely occurs also in  $\gamma_N$  synthesised on other low-SFE  
160 substrates), with such treatment layers possessing a duplex  $\gamma_N/\epsilon_N$  microstructure. The nitrogen  
161 interstitial-induced TRIP effect (and the unique duplex microstructure formed) could have a  
162 complex and profound influence on material mechanical properties.

163

164 Low-temperature thermochemical diffusion treatments (such as carburising and nitriding) have  
165 long been studied (and used) as hard protective surface layers for metallic components from  
166 the aspect of Surface Engineering. Nevertheless, Ren Zheng et al. [23, 24] recently  
167 demonstrated the unique intrinsic properties of  $\gamma_C$ -316. As the size of the “substrate” reduces,  
168 the diffusion-treated layers play an increasingly important role in the performance of the bulk  
169 component. The insertion of interstitials (such as carbon and nitrogen) into the parent lattice  
170 via low-temperature thermochemical diffusion treatments could be exploited to tailor material  
171 strength-ductility and/or other properties for small or thin-section components (such as medical  
172 stents). However, although contributing to improved bulk yield strengths, the  $\gamma_C$ -316 layers  
173 synthesised on both sides on thin AISI 316 foils (of thicknesses ranging from  $\sim 27 \mu\text{m}$  to  $\sim 100$   
174  $\mu\text{m}$ ) after carburising exhibited reduced ductility, showing brittle cleavage fracturing under  
175 tensile testing [23]. AISI 316 foil ( $\sim 20 \mu\text{m}$  thick) was also reported as fracturing into powder  
176 during low-temperature nitriding [7]. Brittle fracturing limits the use of these interstitial-  
177 supersaturated materials (i.e. both  $\gamma_N$ -316 and  $\gamma_C$ -316), typically when the “substrates” are  
178 small. In this regards, the crack-free  $\gamma_N$ -AG17 layer synthesised (as compared to the crack-  
179 containing  $\gamma_N$ -330) [20] and the nitrogen composition-induced TRIP effect provide new  
180 insights on potentially breaking the strength-ductility trade-off, towards exploiting these  
181 interstitially-supersaturated materials in bulk.

182

183 In contrast to  $\epsilon_N$ , shear bands in  $\gamma_N$  (with localised strain and high ‘stored’ strain energy under  
184 the nitrogen composition introduced deformation) appear to be problematic in terms of the  
185 micro-cracking observed along them (as in [20]). **Fig. 4** shows the distorted and discontinuous  
186 lattice fringes inside this shear band. FFT image suggests a strained and somewhat  
187 polycrystalline sub-structure. The elemental distribution, while appearing rather homogenous  
188 for parent  $\gamma_N$ -330, is inhomogeneous within the shear band (**Fig. 2g-l**), indicating redistribution  
189 of both substitutional and interstitial atoms. While the relatively low treatment temperature at  
190  $400^\circ\text{C}$  for this sample limits diffusion of substitutional elements [5], the local migration of such



191 elements within shear bands might be facilitated by the high defect density and/or the intensive  
192 shear stress. Additionally, Cr and N atoms tend to agglomerate in accordance to their high  
193 chemical affinity, and an associated rejection of Ni and Fe atoms. In EDX line profile (**Fig.**  
194 **2h**), Cr content tends to increase with N content, but decreases with increasing Ni and Fe  
195 content. The bright areas in **Fig. 2j** and **Fig. 2l** (i.e. of high Cr and N content) appear dark in  
196 **Fig. 2i** and **Fig. 2k** (i.e. low Fe and Ni content, respectively). This element segregation in shear  
197 bands in the  $\gamma_{\text{N-330}}$  layer raises concerns about material corrosion performance, as the local  
198 low-Cr volumes could be more vulnerable to galvanic attack and may result in pitting or (when  
199 combined with the micro-crack network) crevice corrosion.

200

201 Nevertheless, there is currently no direct evidence of the formation of a stoichiometric CrN  
202 nitride phase within shear bands in  $\gamma_{\text{N-330}}$  layers after 400°C nitriding (see [20] and GAXRD  
203 profiles in **Appendix**). The 3rd EDX spot (along the scan direction in **Fig. 2g**) shows the lowest  
204 (but still significant) Ni content (**Fig. 2h**). Given the high Fe and Ni content still in the high Cr  
205 regions (**Fig. 2h**) and considering the prohibited substitutional migration at 400°C, it is  
206 anticipated that the observed elemental re-distribution is moving towards thermodynamic  
207 equilibrium, but still in a paraequilibrium state [25], without stoichiometric CrN compound  
208 phase formation. The paraequilibrium phase(s) could have slightly different lattice parameter(s)  
209 to the parent  $\gamma_{\text{N-330}}$ . Additionally, the minor alloying elements in the substrate (i.e. 1.4 wt.%  
210 Mn, 1.2 wt.% Si, 0.16 wt.% Al, 0.14 wt.% Ti, 0.11 wt.% Cu and 0.05 wt.% C) could have  
211 migrated and played important roles in the formation of these intermediate phase(s). Rejection  
212 of Mn from high-Ni regions and Si redistribution was observed within a shear band after more  
213 extensive STEM-EDX analysis (see **Appendix**). Probably owing to their small amounts, the  
214 re-distribution of Al, Ti, Cu and C (although likely to have occurred) cannot clearly be seen.

215

216 In this paper, we investigated the structure and local chemical composition of the HCP- $\epsilon_{\text{N}}$   
217 regions and shear bands in the metastable nitrogen diffusion layers on two special ASSs using  
218 advanced TEM analysis. The evident  $\epsilon_{\text{N}}/\gamma_{\text{N}}$  duplex structure under HRTEM and the  
219 homogenous elemental distribution under STEM-EDX in the  $\gamma_{\text{N-AG17}}$  TPN layer is consistent  
220 to the previously proposed nitrogen-composition induced displacive shear transformation  
221 mechanism. HCP- $\epsilon_{\text{N}}$  ‘expands’ in accompaniment with  $\gamma_{\text{N}}$  and might be named as ‘nitrogen-  
222 expanded martensite’. On the other hand, the inhomogeneous elemental distribution in shear  
223 bands in the  $\gamma_{\text{N-330}}$  TPN layer indicates segregation of substitutional elements – even at a low

224 nitriding temperature of 400°C. Although there is no direct evidence of stoichiometric CrN in  
225  $\gamma_N$ -330, volumes of different chemical composition found along shear bands suggest the  
226 formation of intermediate metastable phase(s). It is evident that local volumes of lattice  
227 inhomogeneity may form via displacive plasticity mechanisms in  $\gamma_N$ . Both nitrogen  
228 composition-induced  $\epsilon_N$  formation and shear banding could disrupt the expanded FCC structure  
229 and would play a significant role in resulting material properties and performance. Nitrogen-  
230 composition induced TRIP effects might be exploited to achieve improved strength-ductility  
231 combinations, and permit the use of interstitially-supersaturated ASS materials in bulk form.

232

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237

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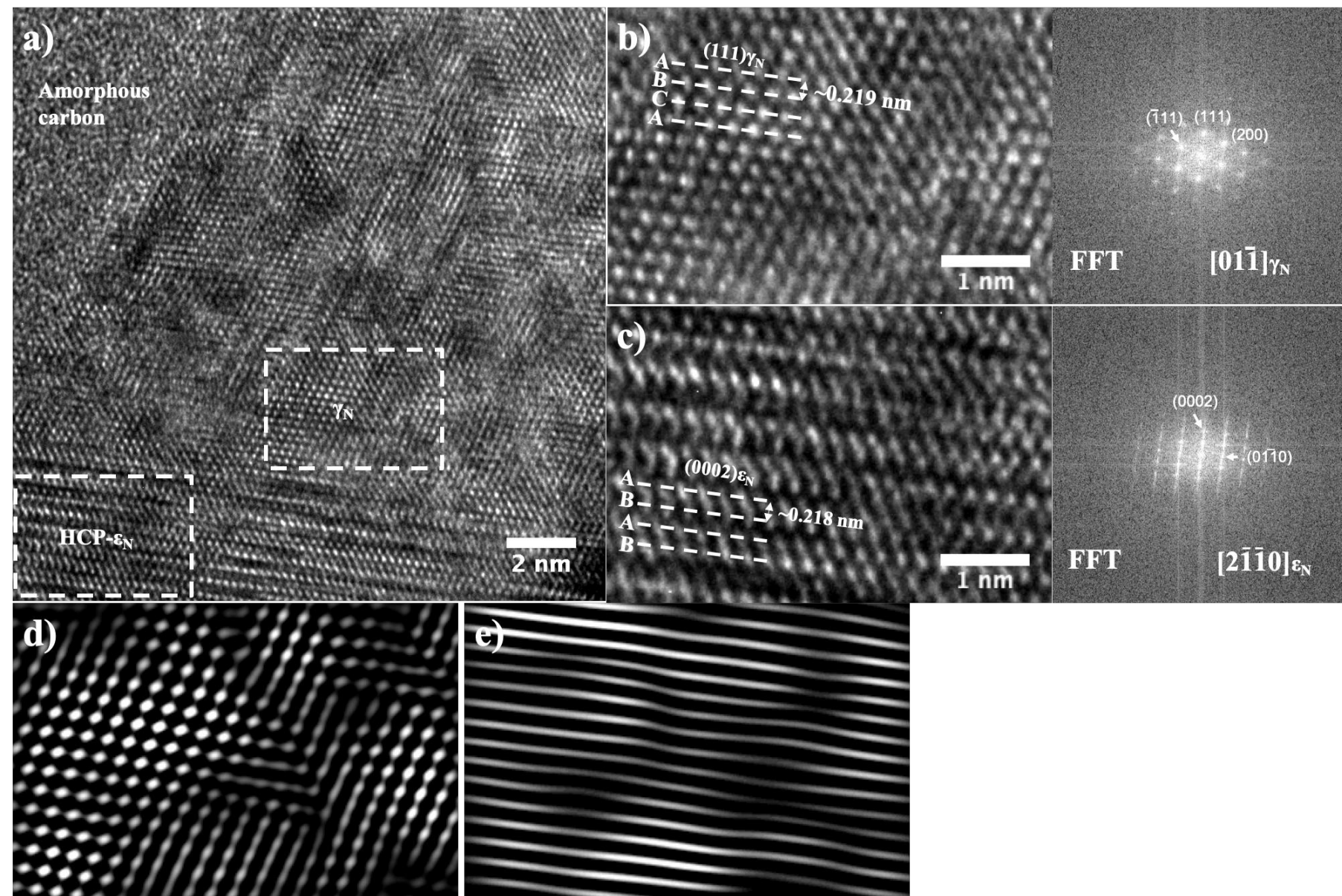
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**Figure 1** a) HRTEM image at the surface of  $\gamma_N$ -AG17 layer; HRTEM images and respective fast Fourier transform (FFT) for b) FCC- $\gamma_N$  and c) HCP- $\epsilon_N$  as indicated in dashed lines in **Fig. 1a**; d) Inverse FFT image for **Fig. 1b**, signifying the 111 and  $\bar{1}\bar{1}\bar{1}$  planes; e) Inverse FFT image for **Fig. 1c**, signifying the 0002 planes. The amorphous carbon layer (at top left corner in **Fig. 1a**) was deposited on the sample surface during the FIB preparation process to protect the underlying material surface from ion beam damage.



**Figure 2** DF-STEM images for HCP- $\epsilon_N$  regions in  $\gamma_N$ -AG17 (**Fig. 2a**) and a typical  $\sim 200$  nm wide shear band in  $\gamma_N$ -330 (**Fig. 2b**); corresponding STEM-EDX maps for  $\epsilon_N$  in  $\gamma_N$ -AG17 (**Fig. 2c-f**) and shear band in  $\gamma_N$ -330 (**Fig. 2i-l**); EDX line scans across  $\epsilon_N$  (**Fig. 2b**) and shear band (**Fig. 2h**). EDX analysis was performed with respect to the main alloy elements (i.e. Fe, Cr, Mn, N for nitrated AG17 and Fe, Cr, Ni, N for nitrated 330). The EDX spots on the line scan are also presented in elemental maps in **Fig. 2i-l**.

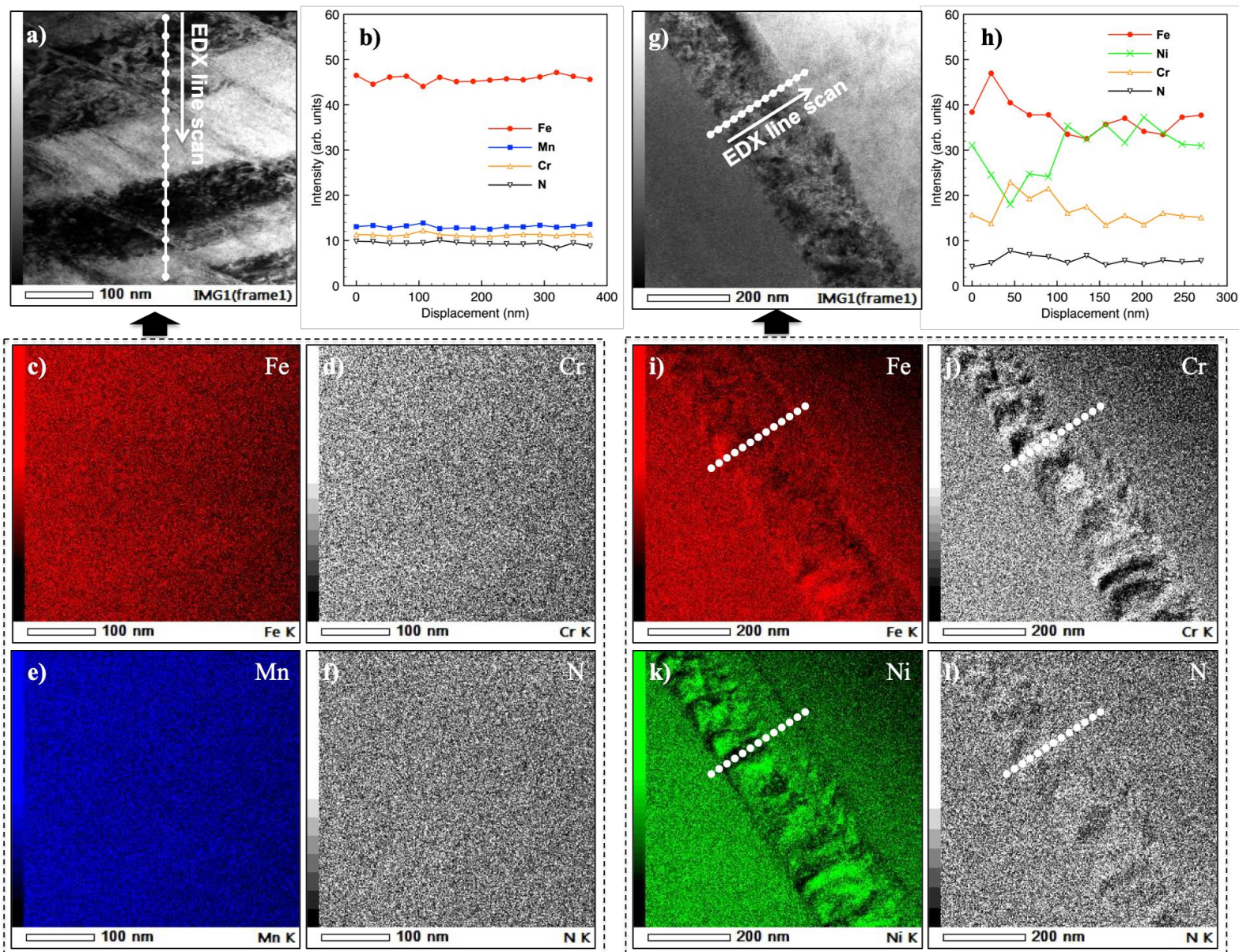


Figure 3 XAXRD profiles of alloy AG17 before and after nitriding treatment at 400° C for 4hrs and 20hrs

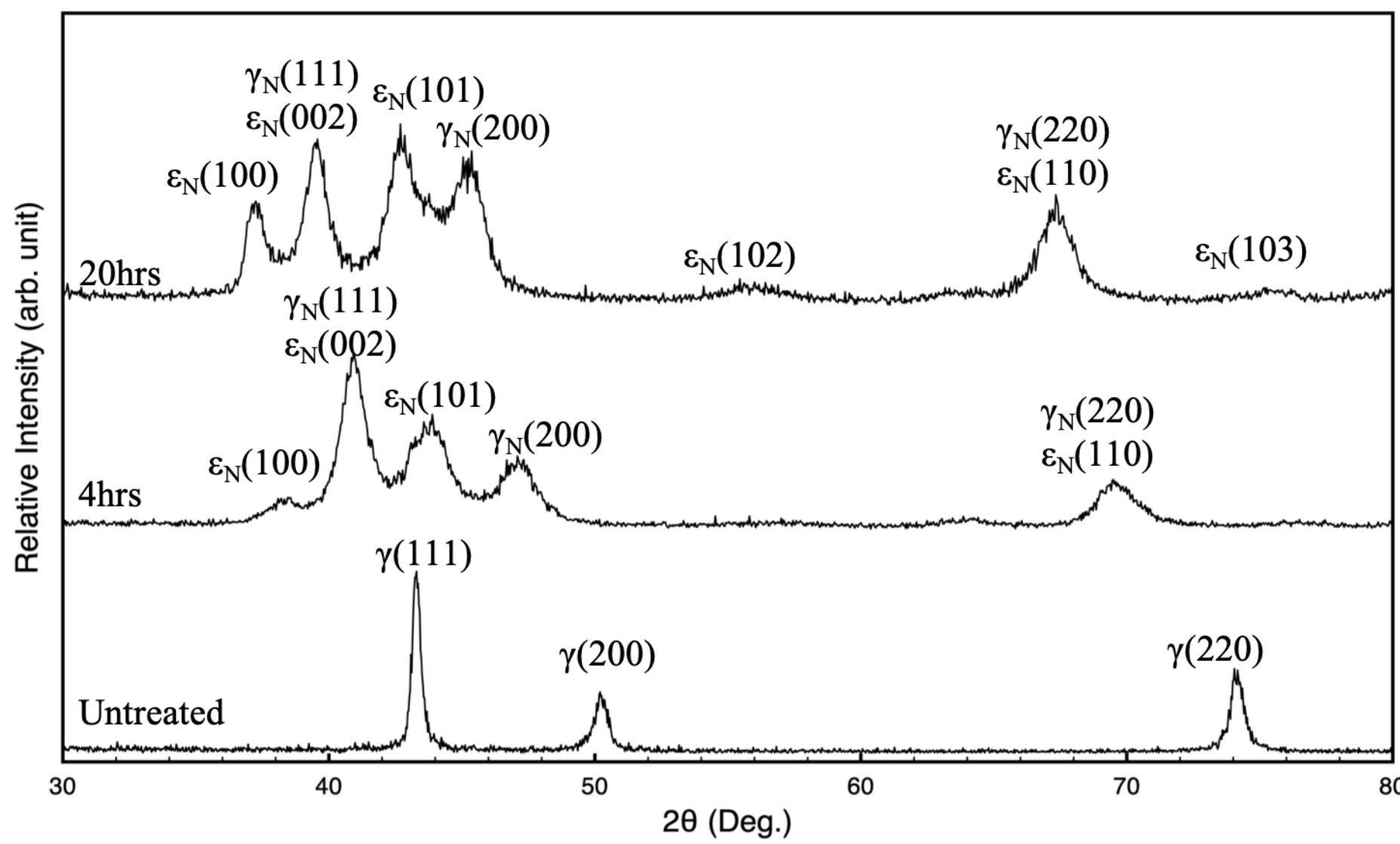
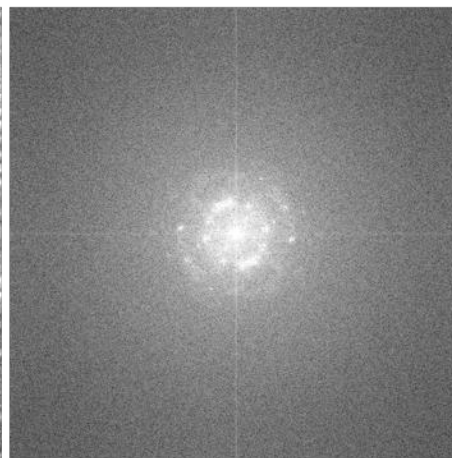
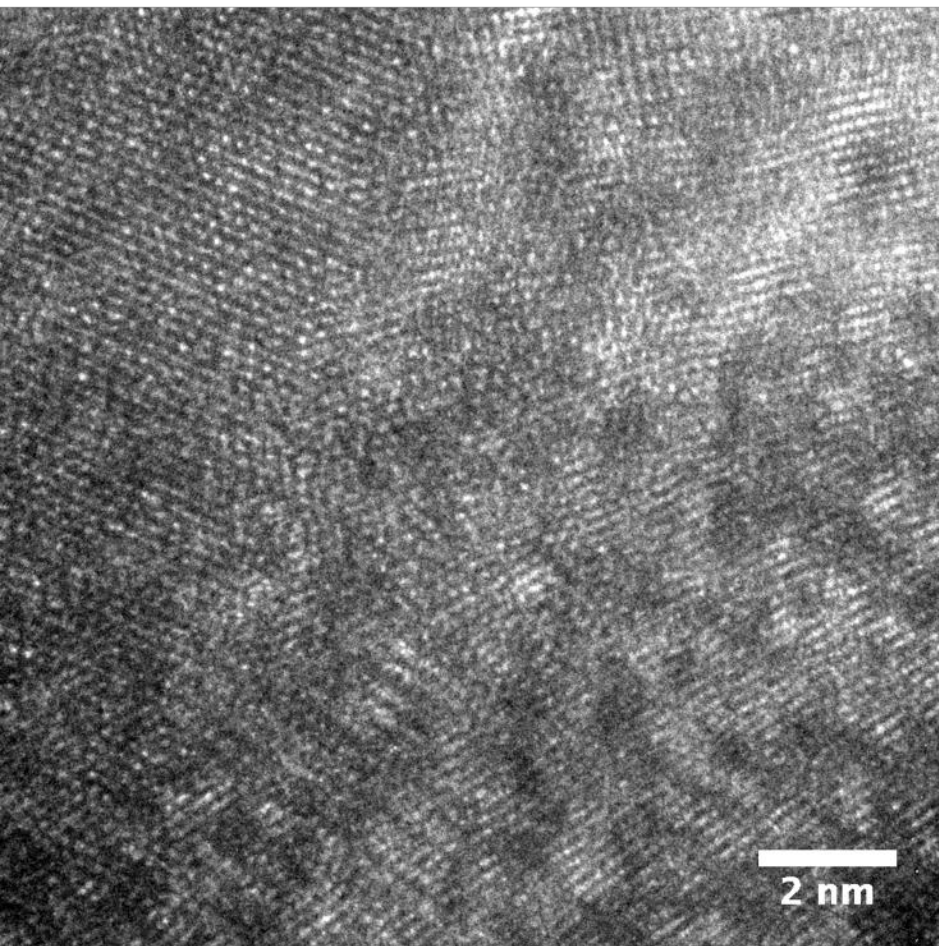
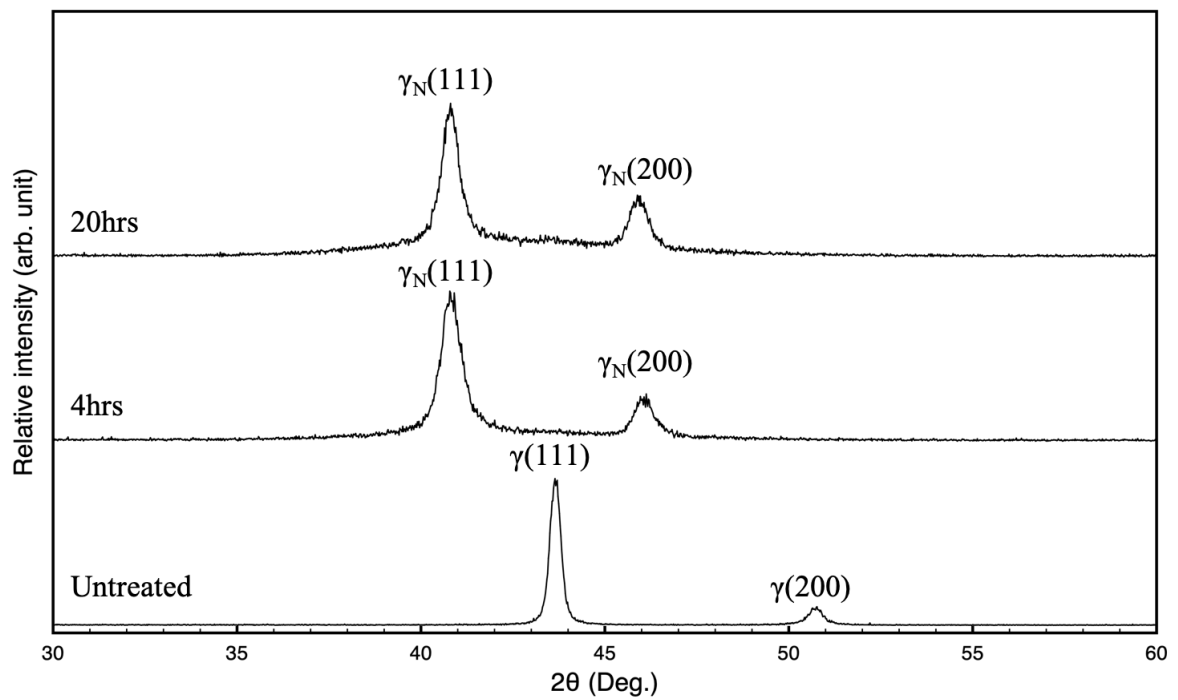


Figure 4 HRTEM images inside a shear band in  $\gamma_N$ -330 layer and corresponding FFT

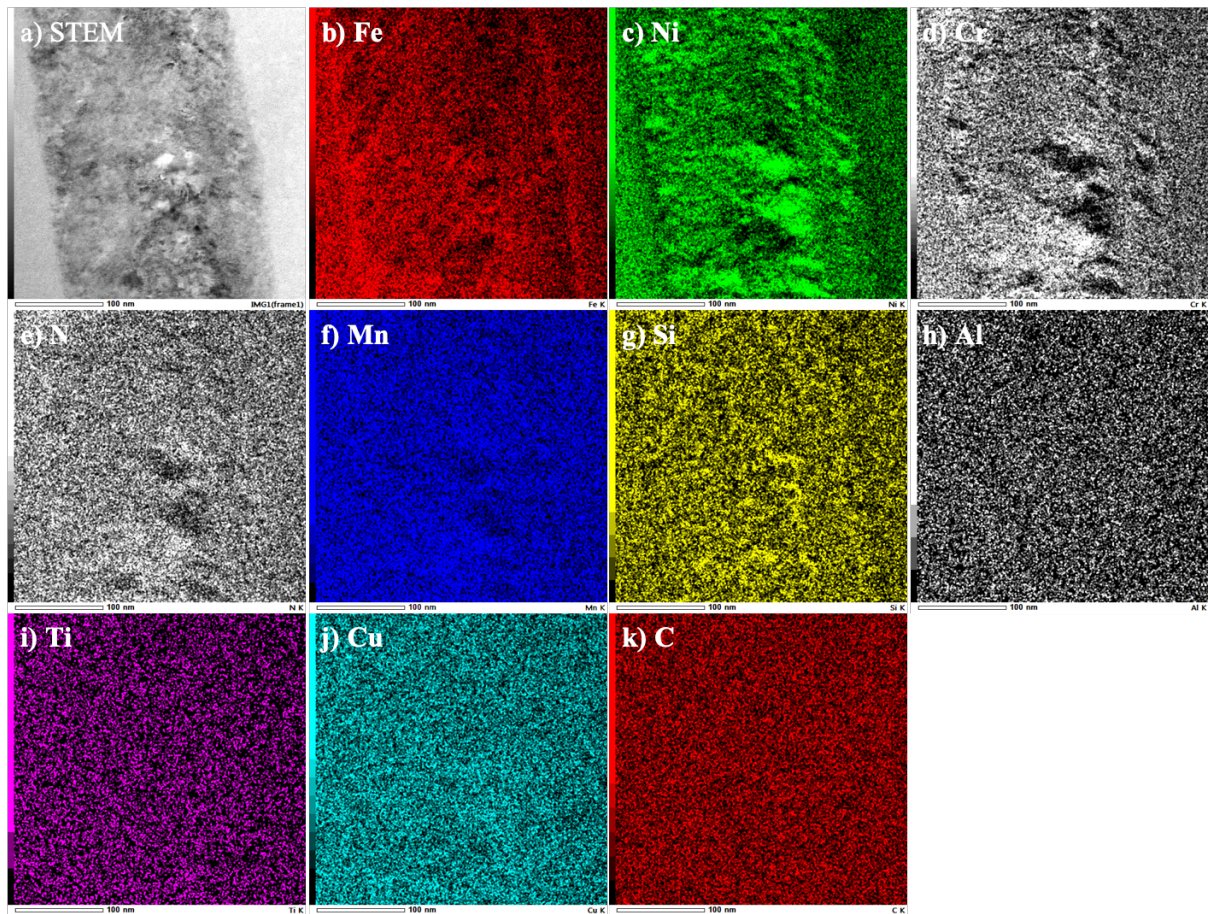


## Appendix



**Fig. 1** GAXRD profiles of alloy 330 before and after nitriding treatment at 400°C for 4hrs and 20hrs. Only  $\gamma_N$  peaks were observed after TPN treatments. A saturated expansion was reached after 4hrs of TPN at 400°C.





**Fig. 2** STEM-EDX mapping for all elements in  $\gamma_N$ -330. a) BF-STEM image for the field of interest; STEM-EDX mapping for elements: b) Fe, c) Ni, d) Cr, e) N, f) Mn, g) Si, h) Al, i) Ti, j) Cu, and k) C. STEM-EDX conditions: C-FEG 200 kV, 10  $\mu\text{m}$  C2 aperture, and Spot size 6.